# 404 Appendix

### 405 **Computing environment**

MNIST benchmark runs were executed on a stand-alone x86 PC using 8 CPU cores, 8 GB RAM,
and no GPUs, The system installs Windows 10, Jupyter notebook 6.4.3, python 3.8.5, and PyTorch
1.9.1. We used MNIST data in PyTorch divided by a batch size of 128.

CIFAR10&100 benchmark runs were executed on an x86 cluster using 8 cores in a single node,
 typically 4-5 GB RAM allocated, and no GPUs, The system installs Linux version 8.5. python
 3.6.9, and PyTorch 1.2.0. We used CIFAR10&100 data in PyTorch divided by a batch size of 256<sup>2</sup>.

## 412 **Present limitations**

<sup>413</sup> Currently, the neural network representation is not perfectly asynchronous because of Eq. 32 to <sup>414</sup> pipeline the coarse-grained dynamics. However, this limitation may make sense considering that <sup>415</sup> biological brains also use slow brain waves to lively regulate their operations. The strategy can <sup>416</sup> reduce  $w_{ii}$  update frequency without much affecting online tracking performance.

Another limitation is that synaptic networks consisting of the axon, synapse, and dendrite are as-417 418 sumed linear, and nonlinear operations are presently dumped into neurons in conjunction with  $\sigma_1$ and  $\sigma_2$ . For example, batch normalization is a nonlinear operation that may be implemented by 419 more naturally adjusting the amplitude and delay distributions of neural signals. Computing the loss 420 function (in this research, the cross-entropy loss was used) is another case in which the computation 421 is now performed outside operator-discretized networks. Instead of making the synaptic network 422 nonlinear in a Cartesian-product state space, it may stay linear in a tensor product state space, which 423 is not entirely impossible as stated later. 424

Furthermore, this representation expects the data to take an event-driven (e.g., time-stamped) format, rather than synchronous streams like video data. In latter cases, some sort of front-end to convert frame-synchronous data may help.

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If the input data streams are appropriately arranged, the application of the idea to a variety of tempo-429 ral ANNs beyond tSNNs should be of interest as a follow-on investigation. The operator-discretized 430 representation will make sense for encoding information into temporal sequences and processing 431 them in time as seen in spike trains in biological neural systems. The fine-grained spike dynamics 432 can be nicely decoupled from the coarse-grained behavioral one. We may regard what is happen-433 ing in the fine-grained part as a sort of vector to time conversion, which is somewhat opposite to 434 435 time2vec [A1]. This approach will work better when the fine-grained dynamics is dominated by predefined temporal correlations, rather than blindly learned from data. 436

The operator representation may be combined with other well-established machine learning tech-437 niques, such as kernel methods [A2], as a means to constitute appropriate basis sets in large spa-438 tiotemporal dimensions. The nonstationary operators can add unique value to such methods by con-439 sistently handling temporal dynamics using amplitude and/or temporal coding. Adjoint operators 440 will naturally incorporate backward dynamics for learning. The use of log probability is popular, 441 such as in log-likelihood or Viterbi algorithm, to better deal with product events. Relating the log 442 of the probability to the Euclidean norm of the signal amplitude as in Eq. 9 for p = 2 is consistent 443 with, for example, what has been done in Viterbi algorithm under Gaussian noise [A3]. 444

Our emulation strategy can become a searchlight to explore future neuromorphic HW. The bidirec-445 tional and elastic nature of our operators may help to natively investigate other physically-oriented 446 (e.g., mechanical) models, such as equilibrium propagation [A4]. The faster turnaround of the pro-447 posed emulation methodology will facilitate detailed comparison across multiple design choices, 448 for example, digital and analog spiking [A5, A6] against reduced precision approaches [A7] using 449 modern AI workloads. We believe that temporal coding is essential for neuromorphic HW to be 450 truly as efficient as the biological brain. Encoding information into the pulse width may rather be 451 considered as non-return-zero rSNNs without much information temporarily. 452

<sup>&</sup>lt;sup>2</sup>Similar results with imagenet in multinode distributed data-parallel to be published upon approval

#### 453 **Biological implications**

Though the present research is on artificial neural networks, originally initiated by treating neural signals as slowly-traveling elastic waves [31, 32], there are growing pieces of works on the importance of wave dynamics in biological neural networks as well [33, 34]. Thus, it may also be worth investigating more biological neuron models, such as Izhikevich neurons [A8] within our representation. Specific functional forms for  $\sigma_1$  and  $\sigma_2$  will affect amplifying and/or damping of collective wave dynamics, which may better elucidate what is happening in biological systems.

Since the geometrical size of spike signals is less than typical axon-synapse-dendrite network sizes, 460 it makes sense, also from a biological standpoint, to explicitly deal with the traveling wave dynamics 461 of spike signals. The width in size of a spike signal is  $\sim 1$  mm for a velocity  $\sim$  of 1 m/s and a width in 462 time of  $\sim 1$  ms. Thus, the axon-synapse-dendrite networks need to be treated as distributed entities, 463 rather than lumped. The biological implications of present LC TL models for axons and dendrites 464 should be argued further in comparison to the conventional RC cable models. LC TLs are superior 465 in better transmitting signal energy and information without dissipation. However, since the origin 466 of a reasonable amount of L is still controversial, the use of L in the biology literature is limited [A9] 467 to our knowledge. 468

Here, we speculate that significant *L* could be caused by more than 4 orders of magnitude heavier masses of ions than that of electrons  $(3.81754 \times 10^{-23} \text{g for Na}^+ \text{ and } 9.1093837 \times 10^{-28} \text{g for e}^-)$ .

471 This is because the kinetic inductance  $L_K$  due to the elastic inertia without scattering given by the

following equation [A10] is also more than 4 orders of magnitude higher:

$$L = L_{EM} + L_K, \quad L_K = \frac{m}{nq^2} \frac{l}{A} \tag{A1}$$

where *m* is mass, *n* density, *q* charge, *l* length, *A* area. For the electric TLs, it is well known that, though microscopic electron motions are diffusive, the coherent electrical signal waves driven by the macroscopic charge density offset is not much affected by them. It should be carefully investigated further whether the same situation holds in more electro-mechanical biological environments with much more complex ion dynamics, and therefore, whether the heavier ion masses can indeed make  $L_K$  a dominant component as Eq. A1 indicates.

Ingress and egress operators can naturally represent orthodromic and antidromic spike transmissions [A11]. Thus, our operator formalism may help to systematically model bidirectional spike
 transmissions in biological systems.

#### 482 Perspective on future AI and QC

It is one of the primary agendas in future computing how AI and QC would evolve in parallel with conventional computing systems. The present approach will shed new light on it as "AI  $\cup$ QC" arguments, alternative to historical "AI  $\cap$  QC" [A12], and facilitate us to unlock unknown mechanisms of the brain. This is because the present idea seems to suggest that classical wave dynamics alone can achieve some limited functionalities of QC by taking advantage of Euclideannorm computing features that have been considered unique to QC [29].

Figure A1 (a) illustrates delay and sum beamforming with delay precoding. It examplifies how the 489 Euclidean norm can enhance the contrast between desired and undesired signals for better perfor-490 mance under a given energy budget (known as the beamforming gain in wireless literature [30]). 491 Freespace should be replaced with a waveguide network for neural signals and the geometrical size 492 will be orders of magnitude reduced as the signal velocity is reduced from the speed of light (3.0 493  $\times 10^8$  m / s) to  $\sim 1$  m/s [31,32]. Well-arranged attenuated superpositions of neural signals from pre-494 ceding neurons can increase desired signal amplitudes with constructive interference, and decrease 495 undesired ones with destructive interference. We still need to carefully work on how to exploit this 496 feature in future computing, but let us discuss some interesting possibilities below. 497

Grover algorithm is well known as a QC algorithm with quadratic speed-up by using amplitude amplification. Here, we would like to argue that the same algorithm is also possible with cubits as shown in Fig.A1(b). By using *n* normalized full cubits instead of  $m = \log_2 n$  qubits, the state  $||\rho\rangle\rangle$ is given as

$$||\rho\rangle\rangle = ||\rho_1\rangle\rangle \oplus \dots \oplus ||\rho_n\rangle\rangle = ||\rho'\rangle\rangle \oplus ||\rho''\rangle\rangle.$$
(A2)



Figure A1: (a) Delay and sum beamforming in wireless communications (freespace is to be replaced with a waveguide network for neural signals); (b) Grover algorithm with cubits; (c) A tensor product state with two cubits. All take advantage of classical wave physics with the Euclidean norm.

502 This corresponds to the following m qubit state:

$$|\rho\rangle = \rho_1 |0_0\rangle \otimes |0_1\rangle \otimes \dots \otimes |0_m\rangle \oplus \rho_2 |1_0\rangle \otimes |0_1\rangle \otimes \dots \otimes |0_m\rangle \oplus \dots \rho_n |1_0\rangle \otimes |1_1\rangle \otimes \dots \otimes |1_m\rangle.$$
(A3)

So, *n* cubits and *m* qubits span the same state space of  $n = 2^m$  dimension. Then the following exactly the same two unitary operations, starting with  $||\rho\rangle\rangle = ||\rho^{init}\rangle\rangle$ , are repeated  $O(n^{1/2})$  times on the cubit state:

$$U_{\rho'} = 2 \left| \left| \rho \right\rangle \right\rangle \left\langle \left\langle \rho \right| \right| - I,\tag{A4}$$

506 and

$$U_{\rho''} = \begin{cases} -||\rho_i\rangle\rangle & \text{if } ||\rho_i\rangle\rangle = ||\rho''\rangle\rangle, \\ ||\rho_i\rangle\rangle & \text{otherwise.} \end{cases}$$
(A5)

Though an exponentially larger number of cubits is required (this may make sense considering  $\sim 10^{11}$  neurons in our brain), the complications associated with encoding *n* data into *m* qubits can be avoided.

In contrast, the exponential speedup in specific QC algorithms, such as quantum Fourier transform and factoring, seems difficult since they take full advantage of tensor product state spaces. Though tensor product state of cubits, such as

$$|\rho_i\rangle\rangle\otimes||\rho_j\rangle\rangle$$
 (A6)

can be defined and constructed with signal multipliers as exemplified in Fig. A1(c), the states can
 entangle only locally and the dimension is limited by the bandwidth.

There are other distinct differences to be mentioned. The interference discussed here of classical waves can occur for signals from different sources. This is a noticeable difference since the qubits interfere only from the same sources. In addition, the coherence time of classical waves can be quite long even at room temperature, as is observed in sound waves, radio waves, ocean waves, and so on [31–34].

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