LOG-CONCAVE SAMPLING ON COMPACT SUPPORTS: A VERSATILE PROXIMAL FRAMEWORK

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ABSTRACT

In this paper, we investigate the theoretical aspects of sampling from strongly log-concave distributions defined on convex and compact supports. We propose a general proximal framework that involves projecting onto the constrained set, which is highly flexible and supports various projection options. Specifically, we consider the cases of Euclidean and Gauge projections, with the latter having the advantage of being performed efficiently using a membership oracle. This framework can be seamlessly integrated with multiple sampling methods. Our analysis focuses on Langevin-type sampling algorithms within the context of constrained sampling. We provide nonasymptotic upper bounds on the W_1 and W_2 errors, offering a detailed comparison of the performance of these methods in constrained sampling.

1 INTRODUCTION

Sampling from probability distributions plays a critical role in various fields of science and engineering, especially when dealing with convex and compact sets (Andrieu et al., 2003; Gelman et al., 1995; Stuart, 2010). In this context, the problem involves sampling from a probability measure ν on such sets, characterized by its density function

$$\nu(x) = \frac{e^{-U(x)}}{\int_{\mathbb{R}^p} e^{-U(s)} \mathrm{d}s} \,,$$

Here, $U(x) = f(x) + \ell_{\mathcal{K}}(x)$, where f(x) represents a potential function and $\ell_{\mathcal{K}}(x)$ is an indicator function ensuring x lies within the convex and compact set $\mathcal{K} \subset \mathbb{R}^p$. Specifically, $\ell_{\mathcal{K}}(x)$ takes the form

$$\ell_{\mathcal{K}} := \begin{cases} +\infty & \text{if } x \notin \mathcal{K} \\ 0 & \text{if } x \in \mathcal{K} \end{cases}$$

Solving this constrained sampling problem is challenging and has garnered considerable interest 037 across various fields, including computer science and statistics. In the realm of computer science, a line of research initiated by Dyer et al. (1991) explored polynomial-time algorithms for uniformly sampling convex bodies. This has been followed by seminal studies on the convergence proper-040 ties of the Ball Walk and the Hit-and-Run algorithm toward uniform density on a convex body 041 or, more broadly, to log-concave densities (Kannan et al., 1997; Kook et al., 2024; Lovász, 1999; 042 Lovász & Simonovits, 1993; Lovász & Vempala, 2007; Smith, 1984). Other Markov Chain Monte 043 Carlo (MCMC) methods, such as Gibbs sampling (Gelfand et al., 1992) and Hamiltonian Monte 044 Carlo (Brubaker et al., 2012; Gürbüzbalaban et al., 2022; Kook et al., 2022), have also been adapted 045 and enhanced to sample from distributions defined on convex and compact sets.

In recent years, leveraging optimization techniques to facilitate sampling has become a prevalent approach. By formulating the sampling challenge as an optimization problem, methods like projected stochastic gradient descent (Bubeck et al., 2015; 2018; Lamperski, 2021; Lehec, 2023), proximal approaches (Brosse et al., 2017; Durmus et al., 2018; Salim & Richtárik, 2020), particle-based algorithms Li et al. (2022), and mirror descent (Ahn & Chewi, 2021; Chewi et al., 2020; Hsieh et al., 2018; Zhang et al., 2020) have proven effective in navigating the target distribution to generate samples. Further innovations have emerged from the intersection of deep learning and neural networks, leading to the development of novel sampling techniques via generative adversarial networks (Goodfellow et al., 2014) and variational autoencoders (Kingma & Welling, 2013). These advanced

methodologies offer promising pathways for sampling from intricate and high-dimensional distributions, particularly those defined on convex and compact sets (Ortiz-Haro et al., 2022). We will present a detailed discussion comparing our work with previous mentioned studies in Appendix A.

In this work, we tackle the challenge posed by a non-smooth target density ν by employing a 058 proximal method. This method leverages a regularization technique that involves projecting onto set \mathcal{K} , effectively transforming the constrained sampling problem into an unconstrained one. Our 060 framework is versatile, accommodating various projection options, such as the Euclidean projec-061 tion—which corresponds to the Moreau envelope of the indicator function $\ell_{\mathcal{K}}$ (Rockafellar & Wets, 062 2009; Durmus et al., 2018; Brosse et al., 2017; Pereyra, 2016)—and the Gauge projection (Lu et al., 063 2022; Mhammedi, 2022), which can be efficiently performed using a membership oracle. Specif-064 ically, we introduce a smooth and strongly convex surrogate distribution that closely mimics the target density. By fine-tuning the regularization parameter, we can significantly reduce approxima-065 tion errors. The advantageous properties of our surrogate distribution enable the effective application 066 of various sampling techniques. In this work, we specifically explore the vanilla Langevin Monte 067 Carlo (LMC) (Roberts & Tweedie, 1996; Dalalyan, 2017; Durmus & Moulines, 2017; Erdogdu & 068 Hosseinzadeh, 2021; Mousavi-Hosseini et al., 2023; Raginsky et al., 2017; Erdogdu et al., 2018; 069 Mou et al., 2022; Erdogdu et al., 2022), kinetic Langevin Monte Carlo (KLMC) (Cheng et al., 2018; 070 Dalalyan & Riou-Durand, 2020; Shen & Lee, 2019; Ma et al., 2021; Zhang et al., 2023), and the 071 parallelized versions of the midpoint randomization method for these algorithms (Shen & Lee, 2019; Yu & Dalalyan, 2024; He et al., 2020; Yu et al., 2023) (referred to as pRLMC and pRKLMC, re-073 spectively), within the context of constrained sampling. We derive the convergence rates for these 074 algorithms in Wasserstein-1 and Wasserstein-2 distances, and provide a detailed comparison of their 075 performance. To this end, we make the following contributions.

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- We establish both general upper and lower bounds for the distance between the smooth approximation of the target density and its original form in Wasserstein-q distance for any q ≥ 1. These bounds are detailed in Proposition 2.1 and Proposition 2.2, respectively.
- In Section 2, we demonstrate that our proposed framework can seamlessly incorporate various projection options, including both Euclidean and Gauge projections.
- In Section 3, we incorporate several MCMC sampling methods, such as (kinetic) Langevin Monte Carlo and the parallelized midpoint randomization method for these algorithms, into our general framework. We present a detailed convergence analysis of these methods in both Wasserstein-1 and Wasserstein-2 distances.

In summary, we develop a comprehensive framework specifically designed for sampling from con-087 vex and compact sets, utilizing a regularization technique that involves projecting onto these constrained sets. Our study notably presents an improved error bound for LMC in constrained sampling settings, as well as the first convergence analysis for KLMC, pRLMC, and pRKLMC algorithms in this scenario. We emphasize that the convergence analysis for these three algorithms relies heavily 091 on the smooth properties of the target density—conditions that are not met in the constrained set-092 ting we examine. Our proposed framework addresses these challenges and is exceptionally flexible 093 and adaptable, accommodating various projection operators and sampling methods. Additionally, it enables clear comparisons of the behaviors of different sampling methods in constrained sampling 094 scenarios. Overall, this new framework provides valuable theoretical insights into the dynamics of 095 constrained sampling. 096

Notation. Denote the *p*-dimensional Euclidean space by \mathbb{R}^p . The letter θ denotes the deterministic vector and its calligraphic counterpart ϑ denotes the random vector. We use \mathbf{I}_p and $\mathbf{0}_p$ to denote, respectively, the $p \times p$ identity and zero matrices. Define the relations $\mathbf{A} \preccurlyeq \mathbf{B}$ and $\mathbf{B} \succeq \mathbf{A}$ for two symmetric $p \times p$ matrices \mathbf{A} and \mathbf{B} to mean that $\mathbf{B} - \mathbf{A}$ is positive semi-definite. The gradient and the Hessian of a function $f : \mathbb{R}^p \rightarrow \mathbb{R}$ are denoted by ∇f and $\nabla^2 f$, respectively. Given any pair of measures μ and ν defined on $(\mathbb{R}^p, \mathcal{B}(\mathbb{R}^p))$, the Wasserstein-q distance between μ and ν is defined as

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$$\mathsf{W}_{q}(\mu,\nu) = \left(\inf_{\varrho\in\Gamma(\mu,\nu)}\int_{\mathbb{R}^{p}\times\mathbb{R}^{p}}\|\boldsymbol{\theta}-\boldsymbol{\theta}'\|_{2}^{q}\,\mathrm{d}\varrho(\boldsymbol{\theta},\boldsymbol{\theta}')\right)^{1/q}, \quad q \ge 1$$

106 where the infimum is taken over all joint distributions ρ that have μ and ν as marginals. For a set 107 $\mathcal{K} \subset \mathbb{R}^p$, we use \mathcal{K}^c to denote its complement. The ceiling function maps $x \in \mathbb{R}$ to the smallest 108 integer greater than or equal to x, denoted by $\lceil x \rceil$.

¹⁰⁸ 2 Smooth Approximation for the Target Density

The absence of smoothness in the target distribution ν presents significant challenges because sampling algorithms often depend on the smoothness properties of the target distribution to effectively explore the space and produce representative samples. Motivated by this, we consider the approximation for $\ell_{\mathcal{K}}$ of the form

$$\ell^{\lambda}_{\mathcal{K}}(oldsymbol{ heta}) := rac{1}{2\lambda^2} \|oldsymbol{ heta} - \mathtt{P}_{\mathcal{K}}(oldsymbol{ heta})\|_2^2$$

where $\lambda > 0$ is the tuning parameter and $P_{\mathcal{K}} : \mathbb{R}^p \to \mathcal{K}$ denotes a projection operator which projects the vector $x \in \mathbb{R}^p$ onto the set \mathcal{K} . Define $U^{\lambda}(\theta) := f(\theta) + \ell_{\mathcal{K}}^{\lambda}(\theta)$, and we the define the corresponding surrogate target density ν^{λ}

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 $\nu^{\lambda}(\boldsymbol{\theta}) = \frac{e^{-U^{\lambda}(\boldsymbol{\theta})}}{\int_{\mathbb{R}^{p}} e^{-U^{\lambda}(\boldsymbol{\theta}')} \mathrm{d}\boldsymbol{\theta}'} \,. \tag{2.1}$

121 Throughout the paper, we define $\mu_k(\nu) = \int \|\boldsymbol{\theta}\|_2^k \nu(d\boldsymbol{\theta})$ as the *k*-th moment of the distribution ν 122 $(k \ge 1)$, and we assume the convex and compact set \mathcal{K} satisfy the following assumption.

Assumption 2.1. Given a positive constant $r \in (0, \infty)$, we assume the Euclidean ball centered at the origin with radius r, denoted by $\mathcal{B}_2(r)$, is contained in \mathcal{K}

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$$\mathcal{B}_2(r) = \{ \boldsymbol{\theta} \in \mathbb{R}^p : \|\boldsymbol{\theta}\|_2 \leqslant r \} \subset \mathcal{K}.$$

This assumption has been commonly made in the work of constrained sampling (Lamperski, 2021; Bubeck et al., 2018; Brosse et al., 2017; Gürbüzbalaban et al., 2022). Moreover, we assume that the potential function f satisfies the following assumption.

Assumption 2.2. The function $f : \mathbb{R}^p \to \mathbb{R}$ is continuously differentiable, and its Hessian matrix $\nabla^2 f$ satisfies

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 $m\mathbf{I}_p \preccurlyeq \nabla^2 f(\boldsymbol{\theta}) \preccurlyeq M\mathbf{I}_p, \qquad \forall \boldsymbol{\theta} \in \mathbb{R}^p.$

133 *Moreover, we assume the function* f *is lower bounded and* $\mathbf{0} = \arg \min_{\boldsymbol{\theta} \in \mathcal{K}} f(\boldsymbol{\theta})$.

The assumption that f attains the minimum at the origin simplifies the presentation of our results. All our results will still hold even if this condition is not met. Given the minimizer of the function f over \mathcal{K} , denoted by θ_* , we simply need to shift all the coordinates and consider the constrained set $\mathcal{K} + \theta_*$. Additionally, we note that the approximation U^{λ} inherits the strong convexity of f. Moreover, we require U^{λ} to be integrable, continuously differentiable, and smooth, even if U is not. **Condition 2.1.** The function U^{λ} is m-strongly convex, continuously differentiable, and M^{λ} -smooth. Moreover, it holds that $\int_{\mathbb{R}^p} e^{-U^{\lambda}(\theta)} d\theta < \infty$.

Below, we present two examples of the approximation $\ell_{\mathcal{K}}^{\lambda}$, each utilizing a different, commonly used projection operator. We will demonstrate that these examples effectively approximate the indicator function $\ell_{\mathcal{K}}$ and the resulting potential functions fulfill Condition 2.1.

Example 2.1. [Moreau envelope] One example of a projection operator is the Euclidean operator, which results in the Moreau envelope of the indicator function $\ell_{\mathcal{K}}$

$$\ell_{\mathcal{K}}^{E,\lambda}(\boldsymbol{\theta}) = \inf_{\boldsymbol{\theta}' \in \mathbb{R}^p} \left(\ell_{\mathcal{K}}(\boldsymbol{\theta}') + \frac{1}{2\lambda^2} \|\boldsymbol{\theta} - \boldsymbol{\theta}'\|_2^2 \right) = \frac{1}{2\lambda^2} \|\boldsymbol{\theta} - \mathsf{P}_{\mathcal{K}}^E(\boldsymbol{\theta})\|_2^2,$$

where $\mathbb{P}^{E}_{\mathcal{K}}$ denotes the Euclidean projection on \mathcal{K} .

¹⁵¹ Define the corresponding surrogate potential $U^{E,\lambda}(\boldsymbol{\theta}) := f(\boldsymbol{\theta}) + \ell_{\mathcal{K}}^{E,\lambda}(\boldsymbol{\theta})$. This surrogate potential $U^{E,\lambda}$ satisfies the following property.

Lemma 2.1. Assume Assumptions 2.1 and 2.2 hold, the potential $U^{E,\lambda}$ satisfies Condition 2.1 with $M^{\lambda} = 1/\lambda^2 + M$.

Example 2.2. [Gauge projection] Another example of the projection operator is the Gauge projection $\mathbb{P}^{G}_{\mathcal{K}} : \mathbb{R}^{p} \to \mathbb{R}$, which is defined as

$$\mathsf{P}^G_{\mathcal{K}}(\boldsymbol{\theta}) := rac{\boldsymbol{\theta}}{g_{\mathcal{K}}(\boldsymbol{\theta})}$$

with a variation of Gauge function (also known as the Minkowski function) $g_{\mathcal{K}} : \mathbb{R}^p \to \mathbb{R}$ of set \mathcal{K} $q_{\mathcal{K}}(\boldsymbol{\theta}) := \inf\{t \ge 1 : \boldsymbol{\theta} \in t\mathcal{K}\}.$ (2.2)

This function encapsulates the scaling required to project θ onto the set \mathcal{K} .

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Set $\ell_{\mathcal{K}}^{G,\lambda}(\boldsymbol{\theta}) = \frac{1}{2\lambda^2} \|\boldsymbol{\theta} - \mathbf{P}_{\mathcal{K}}^G(\boldsymbol{\theta})\|_2^2$. The corresponding surrogate function is defined via $U^{G,\lambda}(\boldsymbol{\theta}) := f(\boldsymbol{\theta}) + \ell_{\mathcal{K}}^{G,\lambda}(\boldsymbol{\theta})$. This surrogate potential $U^{G,\lambda}$ satisfies the following property.

Lemma 2.2. Assume Assumptions 2.1 and 2.2 hold, the surrogate potential $U^{G,\lambda}$ satisfies Condition 2.1 with $M^{\lambda} = 1/\lambda^2 + M$.

In the following, we aim to measure the discrepancy between the surrogate distribution and the target distribution. To this end, we establish an upper bound in W_q distance between the target distribution ν and the approximated distribution ν^{λ} defined in equation 2.1. Our aim is to demonstrate that as λ approaches zero, the distance between these two distributions converges to zero.

Proposition 2.1. Under Assumption 2.1, assume that the potential function f is convex with 0 as its minimizer. Then, for any $q \ge 1$ and any $\lambda \in (0, \frac{r}{2(p+q)})$, it holds that

$$\mathsf{W}_{q}^{q}(\nu,\nu^{\lambda}) \leqslant 3\mu_{q}(\nu)\lambda\left(2p+q\right)/r\,,$$

where $\mu_q(\nu) = \int \|\boldsymbol{\theta}\|_2^q \nu(\mathrm{d}\boldsymbol{\theta}), q \ge 1.$

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In this proposition, we provide a general and precise quantification of the distance between the ap-177 proximate distribution ν^{λ} to the target distribution ν in W_q distance for any $q \ge 1$. When q = 1, 178 our result aligns with the bound established for $W_1(\nu,\nu^{\lambda})$ as stated in Proposition 5 of Brosse et al. 179 (2017). Notably, our Assumption 2.1 is less restrictive than the assumption outlined in Proposition 5 180 of that work, as we do not require the domain to be contained within a ball, as specified in Assump-181 tion H2 of that paper. In Proposition 2.1, the size of \mathcal{K} is captured in the bounds through the radius 182 r and the q-th moment of the target distribution μ_q . In this context, μ_q plays a role analogous to the 183 radius of the ball containing \mathcal{K} , denoted as R in Brosse et al. (2017), while offering a more precise description of the domain's size. 185

We also derive the lower bound for the distance between the approximate distribution ν^{λ} to the target distribution ν in W_q distance, specifically for the scenario where $\mathcal{K} = \mathcal{B}_2(r)$ and the projections discussed in Examples 2.1 and 2.2. In this case, the Gauge projection and Euclidean projection coincide, yielding the following results.

Proposition 2.2. Let $\mathcal{K} = \mathcal{B}_2(r)$ and assume the potential function f is convex and M-smooth, with **0** as its minimizer and $f(\mathbf{0}) = 0$. Set $\ell_{\mathcal{K}} = \ell_{\mathcal{K}}^{E,\lambda}$. Under Assumptions 2.1, for any $q \ge 1$ and any $\lambda \in (0, \frac{r}{2(p+q)} \land \frac{1}{\sqrt{M}})$, it holds that

$$\mathsf{W}_{q}^{q}(\nu,\nu^{\lambda}) \ge \mu_{q}(\nu) \min\left\{ \left| \left(\sqrt{\frac{\pi}{8}} e^{-\frac{3}{4}Mr^{2}} \frac{\lambda(p+q)}{r+3\lambda p} + \frac{r}{r+3\lambda p} \right)^{1/q} - 1 \right|^{q}, \left| \left(\frac{3\lambda(p+q)}{r} + 1 \right)^{1/q} - 1 \right|^{q} \right\}.$$

This result provides a quantitative assessment of the tightness of the upper bound established in Proposition 2.1. Before we discuss the results of the lower bound, we present the following corollary. **Corollary 2.1.** Under the assumptions stated in Proposition 2.2, when $\lambda = o(r/p)$ and q = 1, it holds that

$$\mathsf{W}_1(\nu,\nu^{\lambda}) \geqslant C_{M,r}\mu_1(\nu)\frac{p+1}{r}\lambda\,,$$

where $C_{M,r} > 0$ is a constant that depends on M and r exponentially.

Based on the results outlined in the preceding corollary, when q = 1 and with a sufficiently small λ , the lower bound aligns with the rate estimates from Proposition 2.1 up to a constant factor, thereby confirming the optimality of the upper bound for the case of q = 1. However, when q > 1 and λ is small, the lower bound follows the order $((1 + \lambda)^{1/q} - 1)^q$, which increases at a rate slower than λ . This indicates the suboptimality¹ of the rate stated in Proposition 2.1, highlighting a difference in behavior for higher values of q.

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3 PROXIMAL LANGEVIN ALGORITHMS

A central theme of this work is the approximation of the non-smooth potential ν using a carefully crafted smooth surrogate density ν^{λ} . We note that the convergence analysis for the log-concave

¹We conjecture that this suboptimality arises from the general technique used in our proof. With specific assumptions about the density, a refined lower bound for the Wasserstein distance between ν and ν^{λ} .

sampling algorithms typically relies heavily on the strongly log-concave and smooth properties of the target density. However, these conditions are not met in the constrained setting considered in this work. Owing to the desirable properties of the surrogate density ν^{λ} , we now can utilize various sampling methods that are effective for distributions with smooth and strongly convex densities.

220 In this section, we demonstrate how the proposed scheme can be adapted to various sampling meth-221 ods that collectively aim to sample from the density ν^{λ} . We particularly focus on Langevin Monte 222 Carlo (LMC) and its variant, Kinetic Langevin Monte Carlo (KLMC). LMC employs stochastic 223 differential equations to sample effectively from complex distributions, making it a powerful tool 224 in high-dimensional spaces. KLMC enhances efficiency by incorporating kinetic energy dynamics. 225 Additionally, we explore the parallelized randomized midpoint method for these two algorithms, 226 which improves sampling speed by leveraging parallel computations. Due to space constraints, we provide only a brief overview of each algorithm for the convenience of readers in the following sub-227 section. We refer interested readers to the works of Dalalyan (2017); Durmus & Moulines (2017); 228 Shen & Lee (2019); Yu & Dalalyan (2024); He et al. (2020); Yu et al. (2023) for further details on 229 these algorithms. 230

231 More specifically, we combine the discretization error with the approximation error analyzed in Section 2 to evaluate errors in Wasserstein-1 and Wasserstein-2 distances between the sample dis-232 tribution and the target density ν in the context of constrained sampling. Table 1 below offers a 233 comparison of the results for the four sampling methods mentioned above, using the specific pro-234 jection operators discussed in Example 2.1 and Example 2.2^2 . For simplicity, we omit the constants 235 and logarithmic terms that appear in the bounds. The table illustrates that when the number of par-236 allel steps R > 1, the randomized midpoint method markedly enhances the performance of both 237 the vanilla Langevin Monte Carlo and kinetic Langevin Monte Carlo algorithms. When the number 238 of parallel steps R = 1, the pRLMC and pRKLMC algorithms correspond exactly to the random-239 ized midpoint method applied to Langevin Monte Carlo (RLMC) and kinetic Langevin Monte Carlo 240 (RKLMC), respectively, as detailed in Yu & Dalalyan (2024). The convergence rate for RLMC is 241 comparable to that of LMC when assessing the error in W_2 distance. Further details are provided 242 following Corollary 3.2. 243

	KLMC	pRLMC	pRKLMC
$W_1 \parallel \widetilde{\mathcal{O}}(\varepsilon^{-4})$	$\widetilde{\mathcal{O}}(\varepsilon^{-4})$	$\widetilde{\mathcal{O}}(R^{-1/3}\varepsilon^{-10/3})$	$\widetilde{\mathcal{O}}(R^{-1/3}\varepsilon^{-8/3})$
$\mathbb{W}_2 \parallel \widetilde{\mathcal{O}}(\varepsilon^{-6})$	$\widetilde{\mathcal{O}}(\varepsilon^{-7})$	$\widetilde{\mathcal{O}}(R^{-1/3}\varepsilon^{-6})$	$\widetilde{\mathcal{O}}(R^{-1/3}\varepsilon^{-5})$

Table 1: The number of iterations required by {L,KL,pRL,pRKL}MC algorithms to achieve an error bounded by $\varepsilon \sqrt{p/m}$ in W₁ distance and W₂ distance. $R \ge 1$ denotes the number of parallel steps.

3.1 LANGEVIN MONTE CARLO (LMC)

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Let ϑ_0 be a random vector drawn from a distribution ν^{λ} on \mathbb{R}^p and let $\mathbf{W} = (\mathbf{W}_t : t \ge 0)$ be a *p*-dimensional Brownian motion that is independent of ϑ_0 . To sample from the approximation distribution ν^{λ} , we consider the vanilla Langevin diffusion, which is a strong solution to the stochastic differential equation

$$d\boldsymbol{L}_t^{\mathsf{LD}} = -\nabla U^{\lambda}(\boldsymbol{L}_t^{\mathsf{LD}}) dt + \sqrt{2} d\boldsymbol{W}_t, \qquad t \ge 0, \qquad \boldsymbol{L}_0^{\mathsf{LD}} = \boldsymbol{\vartheta}_0.$$
(3.1)

This equation has a unique strong solution, which is a continuous-time Markov process, termed Langevin diffusion. Under the further assumptions on the potential U^{λ} , such as strong convexity, the Langevin diffusion is ergodic, geometrically mixing and has ν^{λ} as its unique invariant distribution (Bhattacharya, 1978). Moreover, we can sample from the distribution defined by ν^{λ} by using a suitable discretization of the Langevin diffusion. LMC algorithm is based on this idea, combining the considerations with the Euler discretization. Specifically, for small values of $h \ge 0$ and

 ²In this work, we mainly focus on the Wasserstein metric. To our knowledge, the only comparable results are those provided by Brosse et al. (2017) concerning the application of Langevin Monte Carlo. We offer a detailed discussion of this comparison following Corollary 3.1.

270 $\Delta_h W_t = W_{t+h} - W_t$, the following approximation holds

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 $\boldsymbol{L}_{t+h}^{\text{LD}} = \boldsymbol{L}_{t}^{\text{LD}} - \int_{0}^{h} \nabla U^{\lambda}(\boldsymbol{L}_{t+s}^{\text{LD}}) \, \mathrm{d}s + \sqrt{2} \, \Delta_{h} \boldsymbol{W}_{t} \approx \boldsymbol{L}_{t}^{\text{LD}} - h \nabla U^{\lambda}(\boldsymbol{L}_{t}^{\text{LD}}) + \sqrt{2} \, \Delta_{h} \boldsymbol{W}_{t}.$

 $\boldsymbol{\vartheta}_{k+1}^{\mathsf{LMC}} = \boldsymbol{\vartheta}_{k}^{\mathsf{LMC}} - h\nabla U^{\lambda}(\boldsymbol{\vartheta}_{k}^{\mathsf{LMC}}) + \sqrt{2}\left(\boldsymbol{W}_{(k+1)h} - \boldsymbol{W}_{kh}\right).$

By repeatedly applying this approximation with a small step-size h, we can construct a Markov chain $(\vartheta_k^{\text{LMC}} : k \in \mathbb{N})$ that converges to the target distribution ν^{λ} as h goes to zero. More precisely, $\vartheta_k^{\text{LMC}} \approx \boldsymbol{L}_{kh}^{\text{LD}}$, for $k \in \mathbb{N}$, is given by

We now provide explicit upper bounds for the error of the LMC algorithm in terms of W_1 and W_2 distances within the context of constrained sampling.

Theorem 3.1. Under Assumptions 2.1 and 2.2, we further assume that the potential U^{λ} satisfies Condition 2.1. Let the step size $h \in (0, 1/M^{\lambda})$ and the tuning parameter $\lambda \in (0, \frac{r}{2(p+2)})$. Then, for every $n \ge 1$, the distribution ν_n^{LMC} of ϑ_n^{LMC} satisfies

$$\mathsf{W}_2(\nu_n^{\mathsf{LMC}},\nu) \leqslant e^{-\frac{mnh}{2}} \mathsf{W}_2(\nu_0^{\mathsf{LMC}},\nu) + 2\sqrt{\frac{3\mu_2(\nu)\lambda(2p+2)}{r}} + \sqrt{\frac{2phM^{\lambda}}{m}}$$

Moreover, when the initial distribution ν_0^{LMC} is set to be the Dirac measure at the minimizer of f, it holds that

$$\mathsf{W}_1(\nu_n^{\mathsf{LMC}},\nu)\leqslant e^{-\frac{mnh}{2}}\sqrt{\frac{p}{m}}+\frac{3\mu_1(\nu)\lambda(2p+1)}{r}+\sqrt{\frac{2phM^\lambda}{m}}$$

The term M^{λ} introduces an additional factor of λ into the convergence rate. To clarify this dependence on λ , we specify in the following corollary that $M^{\lambda} = M + 1/\lambda^2$. This specification corresponds to the specific projection operators used in Example 2.1 and Example 2.2. We then optimize λ to obtain the results presented below.

Corollary 3.1. Let $\varepsilon \in (0,1)$ be a small number, and $M^{\lambda} = M + \frac{1}{\lambda^2}$. (a) Set $\lambda = \sqrt{2/3} (ph/m)^{1/4} \sqrt{r/(\mu_1(\nu)(2p+1))}$, choose h > 0 and $n \in \mathbb{N}$ so that

$$h=2^{-10}3^{-2}(p/m)r^2\big(\mu_1(\nu)(2p+1)\big)^{-2}\varepsilon^4\quad \text{and}\quad n\geqslant \tfrac{2}{mh}\log(2/\varepsilon)\,,$$

then we have $W_1(\nu_n^{\mathsf{LMC}},\nu) \leq \varepsilon \sqrt{p/m}$.

(b) Set
$$\lambda = 2^{-2/3} 3^{-1/3} (rh)^{1/3} (\mu_2(\nu)m)^{-1/3}$$
, choose $h > 0$ and $n \in \mathbb{N}$ so that
 $h = 2^{-20} 3^{-2} r^2 (\mu_2(\nu)m)^{-2} \varepsilon^6$ and $n \ge \frac{2}{mh} \log(2/\varepsilon)$,

then we have $W_2(\nu_n^{LMC}, \nu) \leq \varepsilon \sqrt{p/m}$.

According to this corollary, when evaluating the error in W₁ distance, the required sample size nto achieve a prespecified error level is of order $\tilde{\mathcal{O}}(\varepsilon^{-4})$. This represents an improvement over the rate of $\tilde{\mathcal{O}}(\varepsilon^{-6})$ obtained in Brosse et al. (2017). Additionally, our general framework provides the convergence rate in W₂ for the proximal LMC algorithm, thus complementing the study of the proximal LMC presented in Brosse et al. (2017).

3.2 PARALLELIZED RANDOMIZED MIDPOINT DISCRETIZATION OF VANILLA LANGEVIN DIFFUSION (PRLMC)

315 As an alternative to the Euler discretization method discussed in Section 3.1 for the stochastic dif-316 ferential equation 3.1, we consider the randomized midpoint method, initially introduced in Shen & 317 Lee (2019), as a different discretization framework. This method enables the discretization and sim-318 ulation of the Langevin diffusion 3.1, ultimately converging to the distribution ν^{λ} . Building upon 319 the foundations laid by Shen & Lee (2019); Yu et al. (2023), Yu & Dalalyan (2024) developed a 320 parallel computing scheme for this algorithm, significantly enhancing the efficiency of the sampling 321 process. The formal definition of pRLMC is defined in Algorithm 1. For simplicity, the superscript pRLMC is omitted therein. We now introduce the primary findings in the work, which establish 322 explicit upper bounds for the error associated with the pRLMC algorithm, measured in W_1 and W_2 323 distances, specifically within the framework of constrained sampling.

Algorithm 1 Parallelized RLMC (pRLMC) **Input:** number of parallel steps R, number of sequential iterations Q, step size h, number of iterations K, initial point $\boldsymbol{\vartheta}_0$. **Output**: iterate ϑ_{K+1} 1: for k = 1 to K do Draw U_{kr} uniformly from $\left[\frac{r-1}{R}, \frac{r}{R}\right], r = 1, \dots, R$. Generate $\boldsymbol{\xi}_{kr} = \boldsymbol{W}_{(k+U_{kr})h} - \boldsymbol{W}_{kh}, r = 1, \dots, R$. 2: 3: Set $\boldsymbol{\vartheta}_k^{(0,r)} = \boldsymbol{\vartheta}_k, r = 1, \dots, R.$ for q = 1 to Q - 1 do 4: 5: $\begin{aligned} q &= 1 \quad \text{to} \quad Q - 1 \text{ and} \\ \text{for } r &= 1 \quad \text{to} \quad R \text{ in parallel do} \\ a_{kj} &= \min\{\frac{1}{R}, U_{kr} - \frac{j-1}{R}\}, j = 1, \dots, r \\ \boldsymbol{\vartheta}_{k}^{(q,r)} &= \boldsymbol{\vartheta}_{k} - h \sum_{j=1}^{r} a_{kj} \nabla U^{\lambda} \big(\boldsymbol{\vartheta}_{k}^{(q-1,j)}\big) + \sqrt{2} \boldsymbol{\xi}_{kr}. \end{aligned}$ 6: 7: In parallel 8: 9: 10: end for $\boldsymbol{\vartheta}_{k+1} = \boldsymbol{\vartheta}_k - \frac{h}{R} \sum_{r=1}^R \nabla U^{\lambda} \left(\boldsymbol{\vartheta}_k^{(Q-1,r)} \right) + \sqrt{2} (\boldsymbol{W}_{(k+1)h} - \boldsymbol{W}_{kh}).$ 11: 12: end for

Theorem 3.2. Under Assumptions 2.1 and 2.2, we further assume that the potential U^{λ} satisfies Condition 2.1. Let the the tuning parameter $\lambda \in (0, \frac{r}{2(p+2)})$. Choose the step size h such that $M^{\lambda}h \leq 1/10$. Then, for every $n \geq 1$, the distribution ν_n^{PRLMC} of $\mathfrak{O}_n^{\mathsf{PRLMC}}$ satisfies

$$\begin{split} \mathsf{W}_{2}(\nu_{n}^{\mathsf{pRLMC}},\nu) &\leqslant \left(1 + \sqrt{\kappa M^{\lambda} h} \big(0.82(M^{\lambda} h)^{Q-1} + 0.94M^{\lambda} h/R \big) \Big) e^{-mnh/2} \mathsf{W}_{2}(\nu_{0}^{\mathsf{pRLMC}},\nu) \\ &+ \sqrt{\kappa M^{\lambda} h} \big(3.98(M^{\lambda} h)^{Q-1} + 6.91M^{\lambda} h/\sqrt{R} \big) \sqrt{p/m} \\ &\left(2 + \sqrt{\kappa M^{\lambda} h} \big(0.82(M^{\lambda} h)^{Q-1} + 0.94M^{\lambda} h/R \big) \big) \sqrt{3\mu_{2}(\nu)\lambda(2p+2)/r} \,, \end{split}$$

where $\kappa = M^{\lambda}/m$. Moreover, when the initial distribution ν_0^{pRLMC} is set to be the Dirac measure at the minimizer of f, it holds that

$$\begin{split} \mathsf{W}_1(\nu_n^{\mathsf{pRLMC}},\nu) &\leqslant \left(1 + \sqrt{\kappa M^{\lambda} h} \big(0.82(M^{\lambda} h)^{Q-1} + 0.94M^{\lambda} h/R\big) \Big) e^{-mnh/2} \sqrt{p/m} \\ &+ \sqrt{\kappa M^{\lambda} h} \big(3.98(M^{\lambda} h)^{Q-1} + 6.91M^{\lambda} h/\sqrt{R}\big) \sqrt{p/m} \\ &+ 3\mu_1(\nu)\lambda(2p+1)/r. \end{split}$$

To our knowledge, this represents the first convergence rate for the pRLMC algorithm within the context of constrained sampling. When R = 1, these results recover the convergence rate for the randomized midpoint method applied to Langevin Monte Carlo algorithm (RLMC) (He et al., 2020; Yu et al., 2023; Yu & Dalalyan, 2024) in constrained sampling. Below, we define $M^{\lambda} = M + 1/\lambda^2$, a setting that corresponds to the projection choices outlined in Examples 2.1 and 2.2. We then optimize λ to derive the following corollary.

Corollary 3.2. Given the number of parallel steps R, we set $Q = \lceil \log R \rceil + 1$. Let $\varepsilon \in (0, 1)$ be a small number, and $M^{\lambda} = M + \frac{1}{\lambda^2}$.

(a) Set
$$\lambda = (\frac{2^{4} \cdot 7}{3})^{1/5} (\frac{p}{R})^{1/10} (\frac{r}{m\mu_1(\nu)(2p+1)})^{1/5} h^{3/10}$$
, choose $h > 0$ and $n \in \mathbb{N}$ so that

$$h = 15.6^{-10/3} p^{2/5} m^{-1} \left(\frac{r}{\mu_1(\nu)(2p+1)}\right)^{8/3} R^{1/3} \varepsilon^{10/3} \quad and \quad n \ge \frac{2}{mh} \log\left(\frac{2.22}{\varepsilon}\right),$$

then we have $W_1(\nu_n^{\mathsf{pRLMC}}, \nu) \leq \varepsilon \sqrt{p/m}$.

(b)
$$\lambda = 320^{2/9} 3^{-1/3} \left(\frac{rp}{m^2 \mu_2(\nu)(2p+2)R}\right)^{1/9} h^{1/3}$$
, choose $h > 0$ and $n \in \mathbb{N}$ so that

$$h = 13^{-6} m^{-7/3} \left(\frac{rp}{\mu_2(\nu)(2p+2)} \right)^{8/3} R^{1/3} \varepsilon^6 \quad and \quad n \ge \frac{2}{mh} \log\left(\frac{2.22}{\varepsilon}\right),$$

then we have $W_2(\nu_n^{\text{pRLMC}},\nu) \leq \varepsilon \sqrt{p/m}$.

Comparing this convergence rate with that obtained for proximal LMC in Corollary 3.1, we observe superior performance of pRLMC over LMC when R > 1. When assessing the error in W₂ distance for the case of R = 1, the sample complexity $\tilde{\mathcal{O}}(\epsilon^{-6})$ aligns with that of LMC. This outcome results from our specific choice of λ during the optimization of the upper bound concerning λ , placing it within a region where the performance of LMC is comparable to that of RLMC.

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3.3 KINETIC LANGEVIN MONTE CARLO (KLMC)

In this part, we explore the application of kinetic Langevin diffusion in constrained sampling. Recall that the kinetic Langevin process L^{KLD} is a solution to a second-order stochastic differential equation that can be informally written as

$$\frac{1}{\gamma} \ddot{\boldsymbol{L}}_{t}^{\mathsf{KLD}} + \dot{\boldsymbol{L}}_{t}^{\mathsf{KLD}} = -\nabla U^{\lambda} (\boldsymbol{L}_{t}^{\mathsf{KLD}}) + \sqrt{2} \, \dot{\boldsymbol{W}}_{t}, \qquad (3.2)$$

with initial conditions $L_0^{\text{KLD}} = \vartheta_0$ and $\dot{L}_0^{\text{KLD}} = \mathbf{v}_0$. In equation 3.2, $\gamma > 0$, W is a standard *p*dimensional Brownian motion and dots are used to designate derivatives with respect to time $t \ge 0$. This can be formalized using Itô's calculus and introducing the velocity field \mathbf{V}^{KLD} so that the joint process ($L^{\text{KLD}}, \mathbf{V}^{\text{KLD}}$) satisfies

$$d\boldsymbol{L}_{t}^{\mathsf{KLD}} = \boldsymbol{V}_{t}^{\mathsf{KLD}} dt; \quad \frac{1}{\gamma} d\boldsymbol{V}_{t}^{\mathsf{KLD}} = -\left(\boldsymbol{V}_{t}^{\mathsf{KLD}} + \nabla U^{\lambda}(\boldsymbol{L}_{t}^{\mathsf{KLD}})\right) dt + \sqrt{2} \, \mathrm{d}\boldsymbol{W}_{t}. \tag{3.3}$$

Similar to the vanilla Langevin diffusion 3.1, the kinetic Langevin diffusion (L^{KLD}, V^{KLD}) is a Markov process that exhibits ergodic properties when the potential U^{λ} is strongly convex (see Eberle et al. (2019) and references therein). The invariant density of this process is given by

$$p_*(\boldsymbol{\theta}, \mathbf{v}) \propto \exp\{-U^{\lambda}(\boldsymbol{\theta}) - \frac{1}{2\gamma} \|\mathbf{v}\|^2\}, \quad \text{for all} \quad \boldsymbol{\theta}, \mathbf{v} \in \mathbb{R}^p.$$

Note that the marginal of p_* corresponds to θ coincides with the target density ν^{λ} . The kinetic Langevin Monte Carlo (KLMC) algorithm is a discretized version of KLD 3.3, where the term $\nabla U^{\lambda}(L_t)$ is replaced by $\nabla U^{\lambda}(L_{kh})$ on each interval [kh, (k+1)h). The resulting error bounds are given in the following theorem.

Theorem 3.3. Under Assumptions 2.1 and 2.2, we further assume that the potential U^{λ} satisfies Condition 2.1. Let the tuning parameter $\lambda \in (0, \frac{r}{2(p+2)})$. Choose γ and h so that $\gamma \ge 5M^{\lambda}$ and $\sqrt{\kappa} \gamma h \le 0.1$, where $\kappa = M^{\lambda}/m$. Assume that $\vartheta_0^{\text{KLMC}}$ is independent of $\mathbf{v}_0^{\text{KLMC}}$ and that $\mathbf{v}_0^{\text{KLMC}} \sim \mathcal{N}_p(0, \gamma \mathbf{I}_p)$. Then, for any $n \ge 1$, the distribution ν_n^{KLMC} of $\vartheta_n^{\text{KLMC}}$ satisfies

$$W_2(\nu_n^{\text{KLMC}},\nu) \leq 2\varrho^n W_2(\nu_0^{\text{KLMC}},\nu) + 3\sqrt{\frac{3\mu_2(\nu)\lambda(2p+2)}{r}} + 0.05\sqrt{\frac{\varrho^n \mathbb{E}[U^{\lambda}(\vartheta_0^{\text{KLMC}}) - f(\mathbf{0})]}{m}} + 0.9\gamma h\sqrt{\frac{\kappa p}{m}}$$
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where $\rho = \rho^{-mh}$. Moreover, let the starting points $\vartheta^{\text{KLMC}} = \mathbf{0}$, it holds that

where $\rho = e^{-mh}$. Moreover, let the staring points $\vartheta_0^{\text{KLMC}} = \mathbf{0}$, it holds that

$$\mathsf{W}_1(\nu_n^{\mathsf{KLMC}},\nu) \leqslant 2\varrho^n \sqrt{\tfrac{p}{m}} + \tfrac{3\mu_1(\nu)\lambda(2p+1)}{r} + 0.9\gamma h \sqrt{\tfrac{\kappa p}{m}}$$

To our knowledge, these results are the first reported convergence rate for the KLMC algorithm within the context of constrained sampling. Below, we set $M^{\lambda} = M + 1/\lambda^2$, aligning with the projection methods detailed in Example 2.1 and Example 2.2. We proceed by optimizing λ to establish the subsequent corollary.

417 **Corollary 3.3.** Let $\varepsilon \in (0, 1)$ be a small number. Set $\gamma = 5M^{\lambda}$, $M^{\lambda} = M + \frac{1}{\lambda^2}$. 418 419 (a) Set $\lambda = 3^{3/4}2^{1/8}r^{1/4}(\mu_1(\nu)(2p+1))^{-1/4}p^{1/8}m^{-1/4}h^{1/4}$, choose h > 0 and $n \in \mathbb{N}$ so that 420 $h = 2^{-7/2}r^3(\mu_1(\nu)(2p+1))^{-3}p^{3/2}m^{-1}\varepsilon^4$ and $n \ge \frac{1}{mh}\log(4/\varepsilon)$, 421 then we have $W_1(\nu_n^{\text{KLMC}}, \nu) \le \varepsilon\sqrt{p/m}$.

423 (b) Set
$$\lambda = 2^{1/7} 3^{3/7} (r/\mu_2(\nu))^{1/7} (h/m)^{2/7}$$
, choose $h > 0$ and $n \in \mathbb{N}$ so that
424 $h = 32.26^{-7} m^{-5/2} r^3 \mu_2(\nu)^{-3} \varepsilon^7$ and $n \ge \frac{1}{mh} \log(4/\varepsilon)$,
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425 426 then we have $W_2(\nu_n^{\mathsf{KLMC}}, \nu) \leq \varepsilon \sqrt{p/m}$.

We note that our error bounds depend in part on the synchronous coupling between the KLMC
and the KLD 3.3. However, for the vanilla Langevin algorithm, Durmus et al. (2019) have shown
that the dependency of the error bound on κ can be improved by employing alternative couplings.
We propose that similar improvements might be achievable for the KLMC algorithm in constrained
sampling scenarios using non-synchronous coupling. For further insights, interested readers are directed to Yu et al. (2023).

432 433 434 3.4 PARALLELIZED RANDOMIZED MIDPOINT DISCRETIZATION OF KINETIC LANGEVIN DIFFUSION (PRKLMC)

The randomized midpoint method, introduced and studied in Shen & Lee (2019), aims at providing
a discretization of the kinetic Langevin process 3.2 that reduces the bias of sampling as compared
to more conventional discretizations. The parallel computing of this algorithm is outlined in Yu &
Dalalyan (2024). This algorithm is referred to as pRKLMC, and for the convenience of the readers,
we restate it in Algorithm 2. To ease the notation, we omit the superscript pRKLMC therein. In

440 Algorithm 2 Parallelized RKLMC (pRKLMC)

441 **Input**: number of parallel steps R, number of sequential iterations Q, step size h, friction coefficient γ , number 442 of iterations K, initial points \mathbf{v}_0 and $\boldsymbol{\vartheta}_0$. 443 **Output**: iterates ϑ_{K+1} and \mathbf{v}_{K+1} 444 1: for k = 1 to K do 445 Draw U_{k1}, \ldots, U_{kR} uniformly from $\left[0, \frac{1}{R}\right], \ldots, \left[\frac{R-1}{R}, 1\right]$, respectively. 2: Generate $\bar{W}_s = W_{kh+s} - W_{kh}$ Generate $\bar{\psi}_s = W_{kh+s} - W_{kh}$ Generate $\xi_{kr} = \int_0^{U_{kr}h} (1 - e^{-\gamma(U_{kr}h-s)}) d\bar{W}_s, r = 1, \dots, R.$ 446 3: 447 4: Set $\boldsymbol{\vartheta}_k^{(0,r)} = \boldsymbol{\vartheta}_k, r = 1, \dots, R.$ for q = 1 to Q - 1 do 448 5: 449 6: for r = 1 to R in parallel do 7: 450 $a_{kr} = \frac{1 - e^{-\gamma h U_{kr}}}{\gamma}.$ $u_{kr} = \frac{\gamma}{p_{kj}} d_{kr} = \int_{\frac{(j-1)h}{R}}^{h\min(\frac{j}{R}, U_{kr})} (1 - e^{-\gamma(U_{kr}h - s)}) ds, j = 1, \dots, r.$ $\vartheta_k^{(q,r)} = \vartheta_k + a_{kr} \mathbf{v}_k - \sum_{j=1}^r b_{kj} \nabla U^{\lambda} (\vartheta_k^{(q-1,j)}) + \sqrt{2} \boldsymbol{\xi}_{kr}.$ In parallel of for 8: 451 452 9: 453 10: 454 end for 11: 455 12: end for $\boldsymbol{\vartheta}_{k+1} = \boldsymbol{\vartheta}_k + \frac{1 - e^{-\gamma h}}{\gamma} \mathbf{v}_k - \sum_{r=1}^R \frac{h}{R} (1 - e^{-\gamma h (1 - U_{kr})}) \nabla U^{\lambda} (\boldsymbol{\vartheta}_k^{(Q-1,r)}) + \sqrt{2} \int_0^h (1 - e^{-\gamma h (1 - U_{kr})}) \nabla U^{\lambda} (\boldsymbol{\vartheta}_k^{(Q-1,r)}) + \sqrt{2} \int_0^h (1 - e^{-\gamma h (1 - U_{kr})}) \nabla U^{\lambda} (\boldsymbol{\vartheta}_k^{(Q-1,r)}) + \sqrt{2} \int_0^h (1 - e^{-\gamma h (1 - U_{kr})}) \nabla U^{\lambda} (\boldsymbol{\vartheta}_k^{(Q-1,r)}) + \sqrt{2} \int_0^h (1 - e^{-\gamma h (1 - U_{kr})}) \nabla U^{\lambda} (\boldsymbol{\vartheta}_k^{(Q-1,r)}) + \sqrt{2} \int_0^h (1 - e^{-\gamma h (1 - U_{kr})}) \nabla U^{\lambda} (\boldsymbol{\vartheta}_k^{(Q-1,r)}) + \sqrt{2} \int_0^h (1 - e^{-\gamma h (1 - U_{kr})}) \nabla U^{\lambda} (\boldsymbol{\vartheta}_k^{(Q-1,r)}) + \sqrt{2} \int_0^h (1 - e^{-\gamma h (1 - U_{kr})}) \nabla U^{\lambda} (\boldsymbol{\vartheta}_k^{(Q-1,r)}) + \sqrt{2} \int_0^h (1 - e^{-\gamma h (1 - U_{kr})}) \nabla U^{\lambda} (\boldsymbol{\vartheta}_k^{(Q-1,r)}) + \sqrt{2} \int_0^h (1 - e^{-\gamma h (1 - U_{kr})}) \nabla U^{\lambda} (\boldsymbol{\vartheta}_k^{(Q-1,r)}) + \sqrt{2} \int_0^h (1 - e^{-\gamma h (1 - U_{kr})}) \nabla U^{\lambda} (\boldsymbol{\vartheta}_k^{(Q-1,r)}) + \sqrt{2} \int_0^h (1 - e^{-\gamma h (1 - U_{kr})}) \nabla U^{\lambda} (\boldsymbol{\vartheta}_k^{(Q-1,r)}) + \sqrt{2} \int_0^h (1 - e^{-\gamma h (1 - U_{kr})}) \nabla U^{\lambda} (\boldsymbol{\vartheta}_k^{(Q-1,r)}) + \sqrt{2} \int_0^h (1 - e^{-\gamma h (1 - U_{kr})}) \nabla U^{\lambda} (\boldsymbol{\vartheta}_k^{(Q-1,r)}) + \sqrt{2} \int_0^h (1 - e^{-\gamma h (1 - U_{kr})}) \nabla U^{\lambda} (\boldsymbol{\vartheta}_k^{(Q-1,r)}) + \sqrt{2} \int_0^h (1 - e^{-\gamma h (1 - U_{kr})}) \nabla U^{\lambda} (\boldsymbol{\vartheta}_k^{(Q-1,r)}) + \sqrt{2} \int_0^h (1 - e^{-\gamma h (1 - U_{kr})}) \nabla U^{\lambda} (\boldsymbol{\vartheta}_k^{(Q-1,r)}) + \sqrt{2} \int_0^h (1 - e^{-\gamma h (1 - U_{kr})}) \nabla U^{\lambda} (\boldsymbol{\vartheta}_k^{(Q-1,r)}) + \sqrt{2} \int_0^h (1 - e^{-\gamma h (1 - U_{kr})}) \nabla U^{\lambda} (\boldsymbol{\vartheta}_k^{(Q-1,r)}) + \sqrt{2} \int_0^h (1 - e^{-\gamma h (1 - U_{kr})}) \nabla U^{\lambda} (\boldsymbol{\vartheta}_k^{(Q-1,r)}) + \sqrt{2} \int_0^h (1 - e^{-\gamma h (1 - U_{kr})}) \nabla U^{\lambda} (\boldsymbol{\vartheta}_k^{(Q-1,r)}) + \sqrt{2} \int_0^h (1 - e^{-\gamma h (1 - U_{kr})}) \nabla U^{\lambda} (\boldsymbol{\vartheta}_k^{(Q-1,r)}) + \sqrt{2} \int_0^h (1 - e^{-\gamma h (1 - U_{kr})}) \nabla U^{\lambda} (\boldsymbol{\vartheta}_k^{(Q-1,r)}) + \sqrt{2} \int_0^h (1 - e^{-\gamma h (1 - U_{kr})}) \nabla U^{\lambda} (\boldsymbol{\vartheta}_k^{(Q-1,r)}) + \sqrt{2} \int_0^h (1 - e^{-\gamma h (1 - U_{kr})}) \nabla U^{\lambda} (\boldsymbol{\vartheta}_k^{(Q-1,r)}) + \sqrt{2} \int_0^h (1 - e^{-\gamma h (1 - U_{kr})}) \nabla U^{\lambda} (\boldsymbol{\vartheta}_k^{(Q-1,r)}) + \sqrt{2} \int_0^h (1 - e^{-\gamma h (1 - U_{kr})}) \nabla U^{\lambda} (\boldsymbol{\vartheta}_k^{(Q-1,r)}) + \sqrt{2} \int_0^h (1 - e^{-\gamma h (1 - U_{kr})}) \nabla U^{\lambda} (\boldsymbol{\vartheta}_k^{(Q-1,r)}) + \sqrt{2} \int_0^h (1 - e^{-\gamma h (1 - U_{kr})}) + \sqrt{2} \int_0^h (1 - e^{-\gamma h (1 - U_{kr})}) \nabla U$ 456 13: 457 $e^{-\gamma(h-s)}d\bar{W}_s.$ 458 $\mathbf{v}_{k+1} = e^{-\gamma h} \mathbf{v}_k - \gamma \sum_{r=1}^{R} \frac{h}{R} e^{-\gamma h(1-U_{kr})} \nabla U^{\lambda} (\boldsymbol{\vartheta}_k^{(Q-1,r)}) + \sqrt{2} \gamma \int_0^h e^{-\gamma (h-s)} d\bar{\boldsymbol{W}}_s.$ 14: 459 15: end for 460

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the theorem below, we quantify the error bounds for this algorithm when applied to constrained sampling.

Theorem 3.4. Under Assumptions 2.1 and 2.2, we further assume that the potential U^{λ} satisfies Condition 2.1. Let the tuning parameter $\lambda \in (0, \frac{r}{2(p+2)})$. Choose γ and h so that $\gamma \ge 5M^{\lambda}$ and $\gamma h \le 0.1 \kappa^{-1/6}$, where $\kappa = M^{\lambda}/m$. Assume that $\vartheta_0^{\mathsf{pRKLMC}}$ is independent of $\mathbf{v}_0^{\mathsf{pRKLMC}}$ and that $\mathbf{v}_0^{\mathsf{pRKLMC}} \sim \mathcal{N}_p(0, \gamma \mathbf{I}_p)$. Then, for any $n \ge 1$, the distribution ν_n^{pRKLMC} of $\vartheta_n^{\mathsf{pRKLMC}}$ satisfies

$$\begin{split} \mathsf{W}_{2}(\nu_{n}^{\mathsf{pRKLMC}},\nu) \leqslant 1.8\varrho^{n} \mathsf{W}_{2}(\nu_{0}^{\mathsf{pRKLMC}},\nu) + 2.8\sqrt{\frac{3\mu_{2}(\nu)\lambda(2p+2)}{r}} + 0.28\sqrt{\frac{\varrho^{n}\mathbb{E}[U^{\lambda}(\vartheta_{0}^{\mathsf{pRKLMC}}) - f(\mathbf{0})]}{m}} \\ + 44.78\sqrt{\frac{(\gamma M^{\lambda}h^{2})^{3}}{R^{2}}} + (\gamma M^{\lambda}h^{2})^{2Q-1}}\sqrt{\frac{\kappa p}{m}}, \end{split}$$

where $\rho = e^{-mh}$. Moreover, let the staring points $\vartheta_0^{\mathsf{pRKLMC}} = \mathbf{0}$, it holds that

$$\mathsf{W}_1(\nu_n^{\mathsf{pRKLMC}},\nu) \leqslant 1.8\varrho^n\sqrt{\frac{p}{m}} + \frac{3\mu_1(\nu)\lambda(2p+1)}{r} + 44.78\sqrt{\frac{(\gamma M^\lambda h^2)^3}{R^2}} + (\gamma M^\lambda h^2)^{2Q-1}\sqrt{\frac{\kappa p}{m}} + \frac{3\mu_1(\nu)\lambda(2p+1)}{r} +$$

To our knowledge, this represents the first convergence analysis for the pRKLMC algorithm within the context of constrained sampling. When R = 1, we recover the convergence rate for RKLMC Shen & Lee (2019); Yu et al. (2023); Yu & Dalalyan (2024) when applied in the context of constrained sampling. We define $M^{\lambda} = M + 1/\lambda^2$, in line with the projection options described in Examples 2.1 and 2.2. We then select the optimal λ to derive the following corollary.

Corollary 3.4. Given the number of parallel steps R, set $Q = \lceil \log R \rceil + 2$. Let $\varepsilon \in (0, 1)$ be a small number, $\gamma = 5M^{\lambda}$ with $M^{\lambda} = M + 1/\lambda^2$.

(a) Set
$$\lambda = (\frac{7 \cdot 8011}{3})^{1/8} (\frac{rh^3}{R\mu_1(\nu)(2p+1)})^{1/8} (\frac{p}{m^2})^{1/16}$$
, choose $h > 0$ and $n \in \mathbb{N}$ so that

$$h = 23.6^{-8/3} p^{7/6} m^{-1} r^{7/3} \left(\mu_1(\nu)(2p+1) \right)^{-7/3} R^{1/3} \varepsilon^{8/3} \quad and \quad n \ge \frac{1}{mh} \log \left(\frac{3.6}{\varepsilon} \right),$$

486 then we have $W_1(\nu_n^{\mathsf{pRKLMC}},\nu) \leq \varepsilon \sqrt{p/m}$. 487

(b) Set
$$\lambda = (\frac{14\cdot8011}{5})^{2/15} (\frac{rp}{m^2\mu_2(\nu)(2p+2)})^{1/15} (\frac{h^3}{R})^{2/15}$$
, choose $h > 0$ and $n \in \mathbb{N}$ so that

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$$(5)$$
 $(m^2\mu_2(\nu)(2p+2))$ (R) (R) $(2p+2)$

$$h = 20.9^{-5} p^{7/3} m^{-13/6} r^{7/3} \left(\mu_2(\nu)(2p+2)\right)^{-1/3} R^{1/3} \varepsilon^5 \quad and \quad n \ge \frac{1}{mh} \log\left(\frac{3.6}{\varepsilon}\right)$$

then we have $W_2(\nu_n^{\mathsf{pRKLMC}},\nu) \leq \varepsilon \sqrt{p/m}$.

The corollary demonstrates that pRKLMC achieves the best performance in terms of sample complexity compared to the other three algorithms considered in this work.

4 DISCUSSION

Strongly convexity In this work, we focus on a strongly convex potential function f coupled 499 with specific choices of projection operators, resulting in a strongly log-concave surrogate density 500 ν^{λ} . However, it is important to note that the approximation error bound between the target density 501 ν and the surrogate density ν^{λ} , as stated in Proposition 2.1, depends solely on the convexity of f. 502 By relaxing the strong convexity requirement of f and considering alternative projection operators, one can generate a surrogate density ν^{λ} that is less restrictive than strongly log-concave. This 504 modification enables the extension of these findings to a broader spectrum of sampling methods. 505

Other metrics The Wasserstein distance we employ is a natural metric for measuring sampling 506 errors due to its relevance to optimal transport theory. However, recent advancements in gradient-507 based sampling have investigated other metrics, including total variation distance, KL divergence, 508 and χ^2 divergence. An intriguing avenue for future research would be to establish error guarantees 509 for constrained sampling concerning these alternative metrics. 510

Limitation The primary focus of this work is to provide theoretical insights into the analysis of 511 constrained sampling. We hope that these insights pave the way for empirical evaluations of the 512 performance of Langevin-type algorithms in various settings, as well as for implementations in real 513 data applications, which offer promising directions for future research. 514

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702 **RELATION TO PREVIOUS WORKS ON CONSTRAINED SAMPLING** А

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704 Sampling from constrained log-concave distributions is a fundamental task across various fields. In 705 computer science, several studies have explored the convergence of the ball walk and hit-and-run 706 algorithms toward the uniform density on a convex body \mathcal{K} , or more generally, a log-concave den-707 sity (Dyer et al., 1991; Kannan et al., 1997; Kook et al., 2024; Lovász, 1999; Lovász & Simonovits, 708 1993; Lovász & Vempala, 2007; Smith, 1984). Unlike the first-order sampling methods examined 709 in this work, these algorithms require calls to zero-order oracles at each iteration. Additionally, the 710 performance of these geometric random walks is influenced by the skewed shape of \mathcal{K} , necessitating pre-processing steps to enhance efficiency. 711

712 Another category of samplers addresses the geometric structure of convex constraints, including 713 Riemannian Hamiltonian Monte Carlo (Brubaker et al., 2012; Gürbüzbalaban et al., 2022; Kook 714 et al., 2022), Gibbs sampling (Gelfand et al., 1992), and Mirror Langevin methods (Ahn & Chewi, 715 2021; Chewi et al., 2020; Hsieh et al., 2018; Zhang et al., 2020).

716 A particularly relevant approach to this work involves diffusion-type samplers (Brosse et al., 2017; 717 Bubeck et al., 2015; 2018; Lamperski, 2021; Lehec, 2021; 2023). These samplers, based on dis-718 cretizations of Itô diffusions, exhibit rapid convergence to the target density in continuous time and 719 have been extensively employed in unconstrained settings. While the underlying stochastic pro-720 cesses can be generalized to constrained settings, the discretization analysis heavily relies on the 721 smoothness of the target distribution, which is challenging to achieve in the constrained context considered here. 722

723 Among these works, Brosse et al. (2017) is the most closely aligned with the spirit of this work. 724 It considers the same setting as the current study, proposing the use of the Moreau envelope (Ex-725 ample 2.1) to derive a smooth surrogate density, followed by applying LMC to sample from this 726 density. Their approach supports our general framework, which integrates Euclidean projection with LMC, and provides convergence bounds in both the W_1 metric and total variation norm. Our 727 work expands on this approach by presenting a general scheme that accommodates a broader range 728 of projection methods and sampling techniques beyond LMC. This enhances the applicability of 729 algorithms such as KLMC, pRLMC, and pRKLMC in constrained settings. Furthermore, we estab-730 lish a lower bound for the distance between the surrogate distribution and the target distribution in 731 Proposition 2.2, demonstrating the tightness of the bound for the distance between the approxima-732 tion and the target density established in Proposition 2.1. This aspect has not been addressed in any 733 of the previously mentioned works on the diffusion-type samplers. 734

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PROOFS OF SECTION 2 В

In this part, we provide the proof of Lemma 2.1, Lemma 2.2, Proposition 2.1, and Proposition 2.2.

B.1 PROOFS OF LEMMA 2.1 AND LEMMA 2.2

Proof of Lemma 2.1. By Theorem 2.26 in Rockafellar & Wets (2009), the function $\ell_{\mathcal{K}}^{E,\lambda}(x)$ is con-742 vex and continuously differentiable, and $\ell_{\mathcal{K}}^{E,\lambda}(x)$ is $\frac{1}{\lambda^2}$ -smooth. The last claim follows from Lemma 743 744 1 in Brosse et al. (2017). 745

Proof of Lemma 2.2. By the definition of $\ell_{\kappa}^{G,\lambda}$, we have

$$\ell_{\mathcal{K}}^{G,\lambda} = \frac{1}{2\lambda^2} \left(1 - \frac{1}{g_{\mathcal{K}}(x)}\right)^2 \|x\|_2^2$$

752 By the definition of $U^{G,\lambda}$, the function $U^{G,\lambda}$ is *m*-strongly convex, continuously differentiable. 753 Moreover, it holds that 754

$$\|\nabla \ell_{\mathcal{K}}^{G,\lambda}(x) - \nabla \ell_{\mathcal{K}}^{G,\lambda}(y)\|_{2} \leqslant \frac{1}{\lambda^{2}} \|x - y\|_{2},$$

This implies $U^{G,\lambda}$ is $\frac{1}{\lambda^2} + M$ -smooth. Now we show that $\int_{\mathbb{R}^p} e^{-U^{G,\lambda}(s)} ds < \infty$. Let $R := \max_{x \in \mathcal{K}} \|x\|_2$, and there is a constant $l \in \mathbb{R}$ such that $f(x) \ge l, \forall x \in \mathbb{R}$. Define $R := \sup_{x \in \mathcal{K}} \|x\|_2$. Note that

$$\int_{\mathbb{R}^p} e^{-U^{G,\lambda}(x)} \mathrm{d}x = \int_{\mathcal{K}} e^{-f(x)} \mathrm{d}x + \int_{\mathcal{B}_2(0,R) \cap \mathcal{K}^c} e^{-U^{G,\lambda}(x)} \mathrm{d}x + \int_{\mathcal{B}_2(0,R)^c} e^{-U^{G,\lambda}(x)} \mathrm{d}x + \int$$

We now derive the bound for the second and the third terms on the right-hand side of the previous display. Since f is convex, there exists $a \in \mathbb{R}$ and $b \in \mathbb{R}^p$ such that

$$f(x) \ge a + \langle b, x \rangle$$

We then find

$$\int_{\mathcal{B}_2(0,R)\cap\mathcal{K}^c} e^{-U^{G,\lambda}(x)} \mathrm{d}x \leqslant e^{-a+bR} \int_{\mathcal{B}_2(0,R)} \mathrm{d}x \,.$$

Thus,

$$\int_{\mathcal{B}_{2}(0,R)\cap\mathcal{K}^{c}} e^{-U^{G,\lambda}(x)} \mathrm{d}x \leqslant e^{-a+bR} \int_{\mathcal{B}_{2}(0,R)} \mathrm{d}x$$
$$\leqslant e^{-a+bR} \int_{\mathcal{B}_{2}(0,R)} \mathrm{d}x$$
$$= e^{-a+bR} \frac{\pi^{p/2}}{\Gamma(\frac{d}{2}+1)} R^{p}$$
$$<\infty.$$

Note that

$$\int_{\mathcal{B}_{2}(0,R)^{c}} e^{-U^{G,\lambda}(x)} \mathrm{d}x \leqslant e^{-l} \int_{\mathcal{B}_{2}(0,R)^{c}} e^{-\frac{1}{2\lambda^{2}} (\|x\|_{2} - R)^{2}} \mathrm{d}x$$

By Fubini's theorem and the fact that

$$e^{-\frac{1}{2\lambda^2}(\|x\|_2 - R)^2} = \int_0^\infty \frac{t}{\lambda^2} e^{-\frac{t^2}{2\lambda^2}} \mathbf{1}_{[\|x\|_2 - R, \infty)}(t) \mathrm{d}t,$$

we obtain

$$\begin{split} \int_{\mathcal{B}_2(0,R)^c} e^{-U^{G,\lambda}(x)} \mathrm{d}x \leqslant e^{-l} \int_0^\infty \frac{t}{\lambda^2} e^{-\frac{t}{2\lambda^2}} \int_{x: \|x\|_2 \ge R} \mathbf{1}_{[\|x\|_2 - R,\infty)}(t) \mathrm{d}x \mathrm{d}t \\ \leqslant e^{-l} \int_0^\infty \frac{t}{\lambda^2} e^{-\frac{t}{2\lambda^2}} \int_{\mathcal{B}_2(0,R+t)} \mathrm{d}x \mathrm{d}t \\ = e^{-l} \int_0^\infty \frac{t}{\lambda^2} e^{-\frac{t}{2\lambda^2}} \frac{\pi^{p/2}}{\Gamma(\frac{p}{2}+1)} (R+t)^d \mathrm{d}t \\ < \infty \,. \end{split}$$

798 Collecting pieces gives the desired result.

B.2 PROOFS OF PROPOSITION 2.1

To prove the Proposition 2.1, we need the following auxiliary lemma.

Lemma B.1. Assume f is convex and $\mathbf{0} = \arg \min_{\boldsymbol{\theta} \in \mathcal{K}} f(\boldsymbol{\theta})$. Under Assumption 2.1, it holds for any $\lambda \in (0, \frac{r}{2(p+k)})$ that

$$\int_{\mathcal{K}^c} \|x\|_2^k \nu^{\lambda}(x) \mathrm{d}x \leqslant \frac{3\lambda \left(p+k\right)}{r} \times \mu_k(\nu) \,,$$

where $\mu_k(\nu) = \int_{\mathbb{R}^p} \|\boldsymbol{\theta}\|_2^k \nu(\mathrm{d}\boldsymbol{\theta}).$

Proof of Lemma B.1. By Fubini's theorem and the fact that

$$\int_0^\infty \frac{t}{\lambda^2} e^{-\frac{t^2}{2\lambda^2}} \mathbf{1}_{[\|x - \mathsf{P}_{\mathcal{K}}(x)\|_2, +\infty)}(t) \mathrm{d}t = e^{-\frac{1}{2\lambda^2} \|x - \mathsf{P}_{\mathcal{K}}(x)\|_2^2}$$

it holds for every $k \ge 0$ that

$$\int_{\mathcal{K}^{c}} \|x\|_{2}^{k} e^{-U^{\lambda}(x)} \mathrm{d}x = \int_{0}^{\infty} \left(\underbrace{\int_{\mathcal{K}^{c}} \|x\|_{2}^{k} e^{-f(x)} \mathbf{1}(\|x - \mathsf{P}_{\mathcal{K}}(x)\|_{2} \leqslant t) \, \mathrm{d}x}_{\psi(t)} \right) \frac{t}{\lambda^{2}} e^{-\frac{t^{2}}{2\lambda^{2}}} \, \mathrm{d}t. \tag{B.1}$$

Define the set $\mathcal{K}_t := \{x \in \mathbb{R}^p : \|x - \mathsf{P}_{\mathcal{K}}(x)\|_2 \leq t\}$, we then have

$$\psi(t) = \int_{\mathcal{K}_t} \|x\|_2^k e^{-f(x)} \, \mathrm{d}x - \int_{\mathcal{K}} \|x\|_2^k e^{-f(x)} \, \mathrm{d}x.$$
(B.2)

On the other hand, it holds for every b > 0 that

$$\int_{\mathcal{K}_t} \|x\|_2^k e^{-f(x)} \, \mathrm{d}x = (1+b)^{p+k} \int_{\mathcal{K}_t/(1+b)} \|y\|_2^k e^{-f((1+b)y)} \, \mathrm{d}y$$
$$\leqslant (1+b)^{p+k} \int_{\mathcal{K}_t/(1+b)} \|y\|_2^k e^{-f(y)} \, \mathrm{d}y$$

where we have used the inequality

$$f((1+b)y) + bf(0) \ge (1+b)f\left(\frac{1}{1+b}(1+b)y + \frac{b}{b+1}0\right) = (1+b)f(y) \ge f(y) + bf(0),$$

implying that $f((1+b)y) \ge f(y)$ for every b > 0 and every $y \in \mathbb{R}^p$. In addition, it holds for any $x \in \mathcal{K}_t/(1+b)$ that

$$\|(1+b)x - \mathsf{P}_{\mathcal{K}}\big((1+b)x\big)\|_2 \leq t$$

which is equivalent to

$$\frac{1+b}{b} \|x - \frac{1}{1+b} \mathsf{P}_{\mathcal{K}} \left((1+b)x \right) \|_2 \leqslant \frac{t}{b} \,.$$

844 Note that we can rewrite x as

$$x = \frac{1}{1+b} \mathsf{P}_{\mathcal{K}} \left((1+b)x \right) + \frac{b}{1+b} \frac{x - \frac{1}{1+b} \mathsf{P}_{\mathcal{K}} \left((1+b)x \right)}{\frac{b}{1+b}}$$

Notice that

$$\mathsf{P}_{\mathcal{K}}\big((1+b)x\big)\in\mathcal{K}$$

and when b = t/r we have

$$\frac{x - \frac{1}{1+b} \mathsf{P}_{\mathcal{K}} \big((1+b)x \big)}{\frac{b}{1+b}} \in \mathcal{K}$$

By the convexity of the set \mathcal{K} , this implies $x \in \mathcal{K}$. Therefore, we obtain

$$\int_{\mathcal{K}_t} \|x\|_2^k e^{-f(x)} \, \mathrm{d}x \leqslant \left(1 + \frac{t}{r}\right)^{p+k} \int_{\mathcal{K}} \|y\|_2^k e^{-f(y)} \, \mathrm{d}y \leqslant e^{(p+k)t/r} \int_{\mathcal{K}} \|y\|_2^k e^{-f(y)} \, \mathrm{d}y \,.$$

861 Combining this inequality with display B.2, we get

$$\psi(t) \leqslant \left(e^{(p+k)t/r} - 1\right) \int_{\mathcal{K}} \|y\|_2^k e^{-f(y)} \,\mathrm{d}y$$

864 Substituting this upper bound in display B.1 leads to

$$\begin{split} \int_{\mathcal{K}^c} \|x\|_2^k e^{-U^{\lambda}(x)} \mathrm{d}x &\leqslant \int_0^\infty \frac{t}{\lambda^2} e^{-t^2/2\lambda^2} \left(e^{(p+k)t/r} - 1 \right) \mathrm{d}t \times \int_{\mathcal{K}} \|y\|_2^k e^{-f(y)} \, \mathrm{d}y \\ &= \int_0^\infty \left(e^{(p+k)t/r} - 1 \right) \mathrm{d} \left(-e^{-t^2/2\lambda^2} \right) \times \int_{\mathcal{K}} \|y\|_2^k e^{-f(y)} \, \mathrm{d}y \\ &= \frac{(p+k)}{r} \int_0^\infty \exp\left\{ -\frac{t^2}{2\lambda^2} + \frac{(p+k)t}{r} \right\} \mathrm{d}t \times \int_{\mathcal{K}} \|y\|_2^k e^{-f(y)} \, \mathrm{d}y \\ &= \frac{(p+k)}{r} \exp\left\{ \frac{\lambda^2(p+k)^2}{2r^2} \right\} \int_0^\infty \exp\left\{ -\frac{(t-(\lambda^2/r)(p+k))^2}{2\lambda^2} \right\} \mathrm{d}t \times \int_{\mathcal{K}} \|y\|_2^k e^{-f(y)} \, \mathrm{d}y \\ &\leqslant \frac{\sqrt{2\pi\lambda}\left(p+k\right)}{r} \exp\left\{ \frac{\lambda^2(p+k)^2}{2r^2} \right\} \times \int_{\mathcal{K}} \|y\|_2^k e^{-f(y)} \, \mathrm{d}y \, . \end{split}$$

When $\lambda < \frac{r}{2(p+k)}$, it holds that

$$\int_{\mathcal{K}^c} \|x\|_2^k e^{-U^\lambda(x)} \mathrm{d}x \leqslant \frac{3\lambda \left(p+k\right)}{r} \times \int_{\mathcal{K}} \|y\|_2^k e^{-f(y)} \,\mathrm{d}y. \tag{B.3}$$

Finally, note that the normalizing constant of ν^{λ} can be lower bounded as follows

$$\int_{\mathbb{R}^p} e^{-U^{\lambda}(x)} \mathrm{d}x \ge \int_{\mathcal{K}} e^{-U^{\lambda}(x)} \mathrm{d}x = \int_{\mathcal{K}} e^{-f(x)} \mathrm{d}x \,.$$

The desired result follows readily by dividing the two sides of display B.3 by $\int_{\mathbb{R}^p} e^{-U^{\lambda}(x)} dx$ and employing this inequality.

We are now ready to prove Proposition 2.1

Proof of Proposition 2.1. By Theorem 6.15 in Villani (2009), it holds that

$$\mathsf{W}_q^q(\nu,\nu^{\lambda}) \leqslant \int_{\mathbb{R}^p} \|x\|_2^q |\nu(x) - \nu^{\lambda}(x)| \mathrm{d}x$$

We note that

$$\begin{split} \int_{\mathbb{R}^p} \|x\|_2^q |\nu(x) - \nu^{\lambda}(x)| \, \mathrm{d}x &= \int_{\mathcal{K}} \|x\|_2^q |\nu(x) - \nu^{\lambda}(x)| \mathrm{d}x + \int_{\mathcal{K}^c} \|x\|_2^q \nu^{\lambda}(x) \mathrm{d}x \\ &= \left(1 - \frac{\int_{\mathcal{K}} e^{-f(y)} \, \mathrm{d}y}{\int_{\mathbb{R}^p} e^{-U^{\lambda}(y)} \, \mathrm{d}y}\right) \int_{\mathcal{K}} \|x\|_2^q \nu(x) \mathrm{d}x + \int_{\mathcal{K}^c} \|x\|_2^q \nu^{\lambda}(x) \mathrm{d}x \\ &= \frac{\int_{\mathcal{K}^c} e^{-U^{\lambda}(y)} \, \mathrm{d}y}{\int_{\mathbb{R}^p} e^{-U^{\lambda}(y)} \, \mathrm{d}y} \times \int_{\mathcal{K}} \|x\|_2^q \nu(x) \mathrm{d}x + \int_{\mathcal{K}^c} \|x\|_2^q \nu^{\lambda}(x) \mathrm{d}x \\ &= \int_{\mathcal{K}^c} \nu^{\lambda}(y) \, \mathrm{d}y \times \mu_q(\nu) + \int_{\mathcal{K}^c} \|x\|_2^q \nu^{\lambda}(x) \mathrm{d}x \,. \end{split}$$

When $\lambda \leq \frac{r}{2(p+q)}$, we apply Lemma B.1 with k = 0 and k = q, which leads us to the inequality

$$\int_{\mathbb{R}^p} \|x\|_2^q |\nu(x) - \nu^{\lambda}(x)| \, \mathrm{d}x \leqslant \frac{3\lambda \left(p\mu_0(\nu) + (p+q)\right)\mu_q(\nu)}{r} = \frac{3\lambda(2p+q)}{r} \times \mu_q(\nu).$$

This completes the proof.

B.3 PROOF OF PROPOSITION 2.2

To prove Proposition 2.2, we need the following auxiliary lemma.

Lemma B.2. Under the assumptions stated in Proposition 2.2, for any $k \ge 1$ and any $\lambda \in$ $\left(0, \frac{r}{2(p+k)} \wedge \frac{1}{\sqrt{M}}\right)$, it holds that

$$\int_{\mathcal{K}^c} \|x\|_2^k \nu^{\lambda}(\mathrm{d}x) / \mu^k(\nu) \ge \frac{1}{2} \sqrt{\frac{\pi}{2}} e^{-\frac{3}{4}Mr^2} \frac{p+k}{r+3\lambda p} \lambda$$

Proof of Lemma B.2. By the definition of ν^{λ} and $\mu^{k}(\nu)$, we have

$$\int_{\mathcal{K}^c} \|x\|_2^k \nu^{\lambda}(\mathrm{d}x) / \mu^k(\nu) = \frac{\int_{\mathcal{K}^c} \|x\|_2^k e^{-f(x) - \frac{1}{2\lambda^2}(\|x\|_2 - r)^2} \mathrm{d}x}{\int_{\mathbb{R}^p} e^{-U^{\lambda}(x)} \mathrm{d}x} \cdot \frac{\int_{\mathcal{K}} e^{-f(x)} \mathrm{d}x}{\int_{\mathcal{K}} \|x\|_2^k e^{-f(x)} \mathrm{d}x} \,. \tag{B.4}$$

Note that f is M-smooth and $f(\mathbf{0}) = 0$, it then holds that

$$f(x) \leqslant \frac{M}{2} \|x\|_2^2$$

This implies

$$\int_{\mathcal{K}^c} \|x\|_2^k e^{-f(x) - \frac{1}{2\lambda^2} (\|x\|_2 - r)^2} \mathrm{d}x \ge \int_{\mathcal{K}^c} \|x\|_2^k e^{-\frac{M}{2} \|x\|_2^2 - \frac{1}{2\lambda^2} (\|x\|_2 - r)^2} \mathrm{d}x$$

Employing the integration in polar coordinates gives

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$$\int_{\mathcal{K}^c} \|x\|_2^k e^{-\frac{M}{2}\|x\|_2^2 - \frac{1}{2\lambda^2}(\|x\|_2 - r)^2} \mathrm{d}x = \omega_{p-1} \int_r^{+\infty} t^{k+p-1} e^{-\frac{M}{2}t^2 - \frac{1}{2\lambda^2}(t-r)^2} \mathrm{d}t,$$

where
$$\omega_{p-1} = \int_{-\pi}^{\pi} \int_{0}^{\pi} \cdots \int_{0}^{\pi} \prod_{j=1}^{p-2} (\sin \theta_{j})^{p-j-1} d\theta_{1} \cdots d\theta_{p-1}$$
. Let $t = r + \lambda z$, it then follows that

$$\int_{\mathcal{K}^{c}} \|x\|_{2}^{k} e^{-f(x) - \frac{1}{2\lambda^{2}} (\|x\|_{2} - r)^{2}} dx$$

$$\geqslant \lambda \omega_{p-1} \int_{0}^{+\infty} (r + \lambda z)^{k+p-1} e^{-\frac{M}{2} (r + \lambda z)^{2} - \frac{1}{2} z^{2}} dz$$

$$\geqslant \lambda \omega_{p-1} r^{k+p-1} \int_{0}^{+\infty} e^{-\frac{M}{2} (r + \lambda z)^{2} - \frac{1}{2} z^{2}} dz$$

$$= \lambda \omega_{p-1} r^{k+p-1} \int_{0}^{+\infty} e^{-\left(\sqrt{\frac{1+M\lambda^{2}}{2}}x + \frac{M\lambda r}{\sqrt{2(1+M\lambda^{2})}}\right)^{2}} dx e^{-\frac{Mr^{2}}{2} \left(1 + \frac{M\lambda^{2}}{1+M\lambda^{2}}\right)}.$$

The last step follows from a change of variables. Using the fact that $\lambda < 1/\sqrt{M}$ and $(x+y)^2 \leq$ $2x^2 + 2y^2, \forall x, y \ge 0$, we then obtain

$$\begin{split} \int_{\mathcal{K}^c} \|x\|_2^k e^{-f(x) - \frac{1}{2\lambda^2} (\|x\|_2^2 - r)^2} \mathrm{d}x &\geqslant \lambda \omega_{p-1} r^{k+p-1} \int_0^{+\infty} e^{-(1+M\lambda^2)x^2} \mathrm{d}x e^{-\frac{Mr^2}{2} (1 + \frac{M\lambda^2}{1+M\lambda^2})} \\ &= \lambda \omega_{p-1} r^{k+p-1} \frac{1}{2} \sqrt{\frac{\pi}{1+M\lambda^2}} e^{-\frac{Mr^2}{2} (1 + \frac{M\lambda^2}{1+M\lambda^2})} \\ &\geqslant \lambda \omega_{p-1} r^{k+p-1} \frac{1}{2} \sqrt{\frac{\pi}{2}} e^{-\frac{3}{4}Mr^2} \,. \end{split}$$

Moreover, by display B.3, we find

$$\int_{\mathbb{R}^p} e^{-U^{\lambda}(x)} \mathrm{d}x = \int_{\mathcal{K}} e^{-U^{\lambda}(x)} \mathrm{d}x + \int_{\mathcal{K}^c} e^{-U^{\lambda}(x)} \mathrm{d}x$$

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$$\int_{\mathcal{K}} g^{p} = \int_{\mathcal{K}} \int_{\mathcal{K}} g^{p} = \int_{\mathcal{K}} g^{p} = \int_{\mathcal{K}} g^{p} dx.$$

972 Furthermore, using the integration in polar coordinates again gives

$$\int_{\mathcal{K}} \|x\|_2^k e^{-f(x)} \mathrm{d}x \leqslant \omega_{p-1} \int_0^r t^{k+p-1} \mathrm{d}t$$
$$= \omega_{p-1} \frac{r^{k+p}}{k+p} \,.$$

Collecting pieces and plugging them into display B.4 gives

$$\int_{\mathcal{K}^c} \|x\|_2^k \nu^{\lambda}(\mathrm{d}x) / \mu^k(\nu) \ge \frac{1}{2} \sqrt{\frac{\pi}{2}} e^{-\frac{3}{4}Mr^2} \frac{p+k}{r+3\lambda p} \lambda$$

as desired.

We are now ready to prove Proposition 2.2.

Proof of Proposition 2.2. By Proposition 7.29 in Villani (2009), it holds for any $q \ge 1$ that

$$\mathsf{W}_q^q(\nu^{\lambda},\nu) \geqslant \left| \left(\int_{\mathbb{R}^p} \|x\|_2^q \nu^{\lambda}(\mathrm{d}x) \right)^{1/q} - \left(\int_{\mathbb{R}^p} \|x\|_2^q \nu(\mathrm{d}x) \right)^{1/q} \right|^q.$$

992 Dividing both sides by $\mu_q(\nu)$ gives

$$\begin{split} \mathsf{W}_{q}^{q}(\nu^{\lambda},\nu)/\mu_{q}(\nu) &\geqslant \left| \left(\int_{\mathbb{R}^{p}} \|x\|_{2}^{q} \nu^{\lambda}(\mathrm{d}x)/\mu_{q}(\nu) \right)^{1/q} - 1 \right|^{q} \\ &= \left| \left(\frac{\int_{\mathcal{K}} \|x\|_{2}^{q} \nu^{\lambda}(\mathrm{d}x) + \int_{\mathcal{K}^{c}} \|x\|_{2}^{q} \nu^{\lambda}(\mathrm{d}x)}{\int_{\mathbb{R}^{p}} \|x\|_{2}^{q} \nu(\mathrm{d}x)} \right)^{1/q} - 1 \right|. \end{split}$$

By the definitions of ν and ν^{λ} , we have

$$\frac{\int_{\mathcal{K}} \|x\|_2^q \nu^{\lambda}(\mathrm{d}x)}{\int_{\mathbb{R}^p} \|x\|_2^q \nu(\mathrm{d}x)} = \frac{\int_{\mathcal{K}} \|x\|_2^q e^{-f(x)} \mathrm{d}x}{\int_{\mathbb{R}^p} e^{-U^{\lambda}} \mathrm{d}x} \cdot \frac{\int_{\mathcal{K}} e^{-f(x)} \mathrm{d}x}{\int_{\mathcal{K}} \|x\|_2^q e^{-f(x)} \mathrm{d}x}$$

$$\int_{\mathcal{K}} e^{-f(x)} \mathrm{d}x$$

$$= \frac{1}{\int_{\mathcal{K}} e^{-f(x)} \mathrm{d}x + \int_{\mathcal{K}^c} e^{-f(x)} \mathrm{d}x}$$

1006 Combining this with display **B**.3 then yields

$$\frac{\int_{\mathcal{K}} \|x\|_2^q \nu^{\lambda}(\mathrm{d}x)}{\int_{\mathbb{R}^p} \|x\|_2^q \nu(\mathrm{d}x)} \geqslant \frac{r}{r+3\lambda p} \,.$$

1010 By Lemma B.2 and Lemma B.1, we have

$$\frac{1}{2}\sqrt{\frac{\pi}{2}}e^{-\frac{3}{4}Mr^2}\frac{p+q}{r+3\lambda p}\lambda \leqslant \frac{\int_{\mathcal{K}^c} \|x\|_2^q\nu^\lambda(\mathrm{d} x)}{\int_{\mathbb{R}^p} \|x\|_2^q\nu(\mathrm{d} x)} \leqslant \frac{3\lambda(p+q)}{r}\,.$$

1014 Collecting pieces gives the desired result.

1018 C PROOFS OF SECTION 3

Proof of Theorem 3.1. The triangle inequality for the Wasserstein distance provides us with

$$W_2(\nu_n^{\mathsf{LMC}},\nu) \leqslant W_2(\nu_n^{\mathsf{LMC}},\nu^{\lambda}) + W_2(\nu^{\lambda},\nu) \,. \tag{C.1}$$

1023 By Theorem 9 in Durmus et al. (2019), it holds that

$$\mathsf{W}_2(\nu_n^{\mathsf{LMC}},\nu^{\lambda}) \leqslant e^{-\frac{mnh}{2}} \mathsf{W}_2(\nu_0^{\mathsf{LMC}},\nu^{\lambda}) + \sqrt{\frac{2M^{\lambda}ph}{m}}$$
(C.2)

provided that $h \leq 1/M^{\lambda}$. Using the triangle inequality for the Wasserstein distance again gives

$$\mathsf{W}_2(\nu_n^{\mathsf{LMC}},\nu^{\lambda}) \leqslant e^{-\frac{mn\hbar}{2}} \mathsf{W}_2(\nu_0^{\mathsf{LMC}},\nu) + \sqrt{\frac{2M^{\lambda}ph}{m}} + \mathsf{W}_2(\nu,\nu^{\lambda}) \,.$$

Plugging this back to display C.1 then provides us with

$$\mathsf{W}_2(\nu_n^{\mathsf{LMC}},\nu) \leqslant e^{-\frac{mnh}{2}} \mathsf{W}_2(\nu_0^{\mathsf{LMC}},\nu) + \sqrt{\frac{2M^\lambda ph}{m}} + 2\mathsf{W}_2(\nu,\nu^\lambda) \,.$$

1035 Combining this with Proposition 2.1, for any $\lambda \in (0, \frac{r}{2(p+2)})$, we obtain

$$\mathsf{W}_2(\nu_n^{\mathsf{LMC}},\nu) \leqslant e^{-\frac{mnh}{2}} \mathsf{W}_2(\nu_0^{\mathsf{LMC}},\nu) + \sqrt{\frac{2M^\lambda ph}{m}} + 2\Big(\frac{3\mu_2(\nu)\lambda(2p+2)}{r}\Big)^{\frac{1}{2}}.$$

1039 This completes the proof of the first claim.

We then proceed to prove the second claim. By the triangle inequality and the monotonicity of Wasserstein distance, we have

$$\begin{split} \mathsf{W}_1(\nu_n^{\mathsf{LMC}},\nu) &\leqslant \mathsf{W}_1(\nu_n^{\mathsf{LMC}},\nu^\lambda) + \mathsf{W}_1(\nu^\lambda,\nu) \\ &\leqslant \mathsf{W}_2(\nu_n^{\mathsf{LMC}},\nu^\lambda) + \mathsf{W}_1(\nu^\lambda,\nu) \,. \end{split}$$

1046 When ν_0^{LMC} is the Dirac mass at the minimizer of the function f, by Proposition 1 in Durmus & 1047 Moulines (2019), it holds that

$$\mathsf{W}_2(\nu_0^{\mathsf{LMC}},\nu^\lambda)\leqslant \sqrt{\frac{p}{m}}$$

Combining this with previous display, display C.2, and Proposition 2.1 then gives

$$\begin{split} \mathsf{W}_1(\nu_n^{\mathsf{LMC}},\nu) \leqslant e^{-\frac{mnh}{2}} \mathsf{W}_2(\nu_0^{\mathsf{LMC}},\nu^{\lambda}) + \sqrt{\frac{2M^{\lambda}ph}{m}} + \mathsf{W}_1(\nu^{\lambda},\nu) \\ \leqslant e^{-\frac{mnh}{2}} \sqrt{\frac{p}{m}} + \sqrt{\frac{2M^{\lambda}ph}{m}} + \frac{3\mu_1(\nu)\lambda(2p+1)}{r} \end{split}$$

1058 as desired.

Proof of Theorem 3.2. The proof employs a similar technique to that used in the proof of Theorem 3.1. Note that $M^{\lambda}h \leq 1/10$, by the triangle inequality and Theorem 1 from Yu & Dalalyan (2024), we have

$$\begin{array}{ll} & \mathsf{W}_{2}(\nu_{n}^{\mathsf{pRLMC}},\nu) \leqslant \mathsf{W}_{2}(\nu^{\lambda},\nu) + \mathsf{W}_{2}(\nu_{n}^{\mathsf{pRLMC}},\nu^{\lambda}) \\ & \mathsf{W}_{2}(\nu_{n}^{\mathsf{pRLMC}},\nu) \leqslant \mathsf{W}_{2}(\nu^{\lambda},\nu) + \mathsf{W}_{2}(\nu_{n}^{\mathsf{pRLMC}},\nu^{\lambda}) \\ & \leq \left(1 + \sqrt{\kappa M^{\lambda}h} \big(0.82(M^{\lambda}h)^{Q-1} + 0.94M^{\lambda}h/R\big)\big) \big) e^{-mnh/2} \mathsf{W}_{2}(\nu_{0}^{\mathsf{pRLMC}},\nu^{\lambda}) \\ & + \sqrt{\kappa M^{\lambda}h} \Big(3.98(M^{\lambda}h)^{Q-1} + 6.91M^{\lambda}h/\sqrt{R}\Big) \sqrt{p/m} + \mathsf{W}_{2}(\nu^{\lambda},\nu) \\ & \leq \left(1 + \sqrt{\kappa M^{\lambda}h} \big(0.82(M^{\lambda}h)^{Q-1} + 0.94M^{\lambda}h/R)\big) \right) e^{-mnh/2} \big(\mathsf{W}_{2}(\nu_{0}^{\mathsf{pRLMC}},\nu) + \mathsf{W}_{2}(\nu,\nu^{\lambda})) \\ & + \sqrt{\kappa M^{\lambda}h} \Big(3.98(M^{\lambda}h)^{Q-1} + 6.91M^{\lambda}h/\sqrt{R}\Big) \sqrt{p/m} + \mathsf{W}_{2}(\nu^{\lambda},\nu) \,. \end{array}$$

1074 Combining this with Proposition 2.1, for any $\lambda \in (0, \frac{r}{2(p+2)})$, we obtain

$$\begin{split} \mathsf{W}_2(\nu_n^{\mathsf{pRLMC}},\nu) &\leqslant \left(1 + \sqrt{\kappa M^\lambda h} \big(0.82(M^\lambda h)^{Q-1} + 0.94M^\lambda h/R\big) \big) e^{-mnh/2} \mathsf{W}_2(\nu_0^{\mathsf{pRLMC}},\nu) \\ &+ \sqrt{\kappa M^\lambda h} \big(3.98(M^\lambda h)^{Q-1} + 6.91M^\lambda h/\sqrt{R}\big) \sqrt{p/m} \end{split}$$

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$$\left(2 + \sqrt{\kappa M^{\lambda} h} \left(0.82(M^{\lambda} h)^{Q-1} + 0.94M^{\lambda} h/R\right)\right) \sqrt{3\mu_2(\nu)\lambda(2p+2)/r} \right)$$

as desired. Invoking Proposition 2.1 and the monotonicity of the Wasserstein distance, since ν_0^{pRLMC} is the Dirac mass at the minimizer of the function f, it holds that

$$\begin{array}{ll} \mathsf{W}_{1}(\nu_{n}^{\mathsf{pRLMC}},\nu) \leqslant \mathsf{W}_{2}(\nu_{n}^{\mathsf{pRLMC}},\nu^{\lambda}) + \mathsf{W}_{1}(\nu^{\lambda},\nu) \\ & \mathsf{W}_{1}(\nu_{n}^{\mathsf{pRLMC}},\nu) \leqslant \mathsf{W}_{2}(\nu_{n}^{\mathsf{pRLMC}},\nu^{\lambda}) + \mathsf{W}_{1}(\nu^{\lambda},\nu) \\ & \leqslant \left(1 + \sqrt{\kappa M^{\lambda}h} (0.82(M^{\lambda}h)^{Q-1} + 0.94M^{\lambda}h/R)\right) e^{-mnh/2} \sqrt{p/m} \\ & + \sqrt{\kappa M^{\lambda}h} (3.98(M^{\lambda}h)^{Q-1} + 6.91M^{\lambda}h/\sqrt{R}) \sqrt{p/m} \\ & \mathsf{1087} \\ & \mathsf{1088} \\ & \mathsf{1089} \\ \end{array}$$
 as desired
$$\begin{array}{l} \mathsf{I} \mathsf{V}_{1}(\nu_{n}^{\mathsf{pRLMC}},\nu) \leqslant \mathsf{W}_{2}(\nu_{n}^{\mathsf{pRLMC}},\nu^{\lambda}) + \mathsf{W}_{1}(\nu^{\lambda},\nu) \\ & \leqslant \left(1 + \sqrt{\kappa M^{\lambda}h} (0.82(M^{\lambda}h)^{Q-1} + 0.94M^{\lambda}h/R)\right) \sqrt{3\mu_{2}(\nu)\lambda(2p+2)/r} \\ & \mathsf{I} \mathsf{V}_{2}(\nu) \mathsf{V}_{2}(\nu) \mathsf{V}_{2}(\nu) \mathsf{V}_{2}(\nu) \mathsf{V}_{2}(\nu) \mathsf{V}_{2}(\nu) \\ & \mathsf{V}_{2}(\nu) \mathsf{V}_{2}(\nu) \mathsf{V}_{2}(\nu) \mathsf{V}_{2}(\nu) \mathsf{V}_{2}(\nu) \mathsf{V}_{2}(\nu) \\ & \mathsf{V}_{2}(\nu) \\ & \mathsf{V}_{2}(\nu) \mathsf{V$$

as desired.

Proof of Theorem 3.3. The proof relies a similar technique to that used in the proof of Theorem 3.1. The key distinction lies in the utilization of Theorem 3 from Yu et al. (2023) to bound the W_2 distance between ν_n^{KLMC} and ν^{λ} . By the triangle inequality and Theorem 3 from Yu et al. (2023), we have

$$\begin{aligned} & \mathsf{W}_{2}(\nu_{n}^{\mathsf{KLMC}},\nu) \leqslant \mathsf{W}_{2}(\nu^{\lambda},\nu) + \mathsf{W}_{2}(\nu_{n}^{\mathsf{KLMC}},\nu^{\lambda}) \\ & \mathsf{W}_{2}(\nu_{n}^{\mathsf{KLMC}},\nu) \leqslant 2\varrho^{n} \mathsf{W}_{2}(\nu_{0}^{\mathsf{KLMC}},\nu^{\lambda}) + 0.05\sqrt{\frac{\varrho^{n} \mathbb{E}[U^{\lambda}(\boldsymbol{\vartheta}_{0}^{\mathsf{KLMC}}) - f(\mathbf{0})]}{m}} + 0.9\gamma h\sqrt{\frac{\kappa p}{m}} + \mathsf{W}_{2}(\nu^{\lambda},\nu) \\ & \mathsf{W}_{2}(\nu_{0}^{\mathsf{KLMC}},\nu) + \mathsf{W}_{2}(\nu,\nu^{\lambda})) + 0.05\sqrt{\frac{\varrho^{n} \mathbb{E}[U^{\lambda}(\boldsymbol{\vartheta}_{0}^{\mathsf{KLMC}}) - f(\mathbf{0})]}{m}} + 0.9\gamma h\sqrt{\frac{\kappa p}{m}} \\ & + \mathsf{W}_{2}(\nu^{\lambda},\nu) \,. \end{aligned}$$

Combining this with Proposition 2.1, for any $\lambda \in (0, \frac{r}{2(p+2)})$, we obtain

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$$W_{2}(\nu_{n}^{\mathsf{KLMC}},\nu) \leq 2\varrho^{n}W_{2}(\nu_{0}^{\mathsf{KLMC}},\nu) + 3\sqrt{\frac{3\mu_{2}(\nu)\lambda(2p+2)}{r}} + 0.05\sqrt{\frac{\varrho^{n}\mathbb{E}[U^{\lambda}(\vartheta_{0}^{\mathsf{KLMC}}) - f(\mathbf{0})]}{m}} + 0.9\gamma h\sqrt{\frac{\kappa p}{m}}$$

as desired.

Invoking Proposition 2.1 and the monotonicity of the Wasserstein distance, since $\vartheta_0^{\text{KLMC}} = 0$, it holds that

$$\begin{split} \mathsf{W}_1(\nu_n^{\mathsf{KLMC}},\nu) &\leqslant \mathsf{W}_2(\nu_n^{\mathsf{KLMC}},\nu^\lambda) + \mathsf{W}_1(\nu^\lambda,\nu) \\ &\leqslant 2\varrho^n \sqrt{\frac{p}{m}} + \frac{3\mu_1(\nu)\lambda(2p+1)}{r} + 0.9\gamma h \sqrt{\frac{\kappa p}{m}} \end{split}$$

as desired.

Proof of Theorem 3.4. By the triangle inequality and Theorem 2 from Yu & Dalalyan (2024), we have MARCENERGIAN () A MARCENERGIAN

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$$W_{2}(\nu_{n}^{\text{pRKLMC}},\nu) \leqslant W_{2}(\nu^{\lambda},\nu) + W_{2}(\nu_{n}^{\text{pRKLMC}},\nu^{\lambda})$$
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$$\leqslant 1.8\varrho^{n}W_{2}(\nu_{0}^{\text{pRKLMC}},\nu^{\lambda}) + 0.28\sqrt{\frac{\varrho^{n}\mathbb{E}[U^{\lambda}(\vartheta_{0}^{\text{pRKLMC}}) - f(\mathbf{0})]}{m}}$$
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$$+ 44.78 \sqrt{\frac{(\gamma M^{\lambda} h^2)^3}{R^2}} + (\gamma M^{\lambda} h^2)^{2Q-1} \sqrt{\frac{\kappa p}{m}} + \mathsf{W}_2(\nu^{\lambda}, \nu)$$

$$44.78\sqrt{(\gamma M^{\lambda}h^2)^3}$$

Combining this with Proposition 2.1, for any $\lambda \in (0, \frac{r}{2(p+2)})$, we obtain

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W_2(\nu_n^{\mathsf{pRKLMC}}, \nu) \leq 1.8 \varrho^n W_2(\nu_0^{\mathsf{pRKLMC}}, \nu) + 0.28 \sqrt{\frac{\varrho^n \mathbb{E}[U^{\lambda}(\boldsymbol{\vartheta}_0^{\mathsf{pRKLMC}}) - f(\mathbf{0})]}{m}}
$$\frac{1}{m}$$
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$$+ 44.78 \sqrt{\frac{(\gamma M^{\lambda} h^2)^3}{R^2}} + (\gamma M^{\lambda} h^2)^{2Q-1} \sqrt{\frac{\kappa p}{m}} + 2.8 \sqrt{\frac{3\mu_2(\nu)\lambda(2p-1)}{r}} + \frac{3\mu_2(\nu)\lambda(2p-1)}{r}$$

as desired.

1143 Invoking Proposition 2.1 and the monotonicity of the Wasserstein distance, since $\vartheta_0^{\text{pRKLMC}} = 0$, it 1144 holds that 1145 holds that

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W₁(
$$\nu_n^{\text{pRKLMC}}, \nu$$
) \leq W₂($\nu_n^{\text{pRKLMC}}, \nu^{\lambda}$) + W₁(ν^{λ}, ν)
 \leq 1.8 $\varrho^n \sqrt{\frac{p}{m}} + \frac{3\mu_1(\nu)\lambda(2p+1)}{r} + 44.78\sqrt{\frac{(\gamma M^{\lambda}h^2)^3}{R^2} + (\gamma M^{\lambda}h^2)^{2Q-1}}\sqrt{\frac{\kappa p}{m}}$

as desired.