

---

# Separable PINN - Supplementary Materials

---

## 1 A Source Code

2 We provided a demo code to show the implementation details of our work. You can run the Jupyter  
3 notebook file to reproduce the (2+1)-d Klein-Gordon experiment.

## 4 B Pseudocode

5 Alg. 1 shows the pseudocode of forward pass and feature merging step of SPINN when solving  
6 a 3-dimensional PDE. The outer product and summation between feature vectors can be easily  
7 implemented and parallelized by the einsum function.

---

**Algorithm 1** JAX-style pseudocode for SPINN’s forward pass and feature merging.

---

```
1 import jax.numpy as jnp
2 from flax import linen as nn
3
4 class SPINN(nn.Module):
5     hidden_sizes: Sequence[int] # list of feature sizes for each layer
6
7     # forward pass and feature merging
8     @nn.compact
9     def forward(self, x, y, z): # 3-dimensional PDE
10         inputs = [x, y, z] # full batch of the sampled coordinates
11         feats = []
12         for X in inputs: # three individual MLPs for each input x, y, z
13             for hidden_size in self.hidden_sizes[:-1]:
14                 X = nn.Dense(hidden_size)(X)
15                 X = nn.activation.tanh(X)
16                 """
17                 the output feature size (self.hidden_sizes[-1]) corresponds to the
18                 rank of the predicted solution tensor (in case of 1-dimensional output)
19                 """
20                 X = nn.Dense(self.hidden_sizes[-1])(X)
21                 feats += [X] # final feature representations
22                 # feature merging (outer product between columns of feats)
23                 xy = jnp.einsum('fx, fy->fxy', feats[:, 0], feats[:, 1])
24                 return jnp.einsum('fxy, fz->xyz', xy, feats[:, 2])
```

---

## 8 C Proof of Theorem 1

9 Here we show the preliminary lemmas and proofs for Theorem 1. in the main paper. We start by  
10 defining a general tensor product between two Hilbert spaces.

11

12 **Definition 1.** Let  $\{v_\beta\}$  be an orthonormal basis for  $\mathcal{H}_2$ . **Tensor product** between Hilbert spaces  
13  $\mathcal{H}_1$  and  $\mathcal{H}_2$ , denoted by  $\mathcal{H}_1 \otimes \mathcal{H}_2$ , is a set of all antilinear mappings  $A : \mathcal{H}_2 \rightarrow \mathcal{H}_1$  such that  
14  $\sum_\beta \|Av_\beta\|^2 < \infty$  for every orthonormal basis for  $\mathcal{H}_2$ .

15 Then by Theorem 7.12 in Folland [3],  $\mathcal{H}_1 \otimes \mathcal{H}_2$  is also a Hilbert space with respect to norm  $\|\cdot\|$  and  
 16 associated inner product  $\langle \cdot, \cdot \rangle$ :

$$\|A\|^2 \equiv \sum_{\beta} \|Av_{\beta}\|^2, \quad (1)$$

$$\langle A, B \rangle \equiv \sum_{\beta} \langle Av_{\beta}, Bv_{\beta} \rangle, \quad (2)$$

17 where  $A, B \in \mathcal{H}_1 \otimes \mathcal{H}_2$ , and  $\{v_{\beta}\}$  is any orthonormal basis of  $\mathcal{H}_2$ .

18

19 **Lemma 1.** Let  $x \in \mathcal{H}_1$  and  $y, y' \in \mathcal{H}_2$ . Then,  $(x \otimes y)y' = \langle y, y' \rangle x$ .

20 *Proof.* Let  $A$  be a mapping  $A : y' \rightarrow \langle y, y' \rangle x$ , where  $\|y\| = 1$ . We expand  $y'$  with orthonormal basis  
 21  $\{y, z_{\beta}\}$ . i.e.,  $\{y, z_{\beta}\}$  is a basis for  $\mathcal{H}_2$ . Then,

$$\|Ay\|^2 + \sum_{\beta} \|Az_{\beta}\|^2 = \|Ay\|^2 + \sum_{\beta} \|\langle y, z_{\beta} \rangle x\|^2 \quad (3)$$

$$= \|Ay\|^2 < \infty \quad (4)$$

22 It is obvious that  $A$  is antilinear. Then by Definition 1,  $A \in \mathcal{H}_1 \otimes \mathcal{H}_2$  which is,  $Ay' = (x \otimes y)y'$ .  
 23 Therefore,  $(x \otimes y)y' = \langle y, y' \rangle x$ .

24

□

25 **Lemma 2.**  $\{u_{\alpha} \otimes v_{\beta}\}$  is an orthonormal basis for  $\mathcal{H}_1 \otimes \mathcal{H}_2$ .

26 *Proof.* Let  $A, B \in \mathcal{H}_1 \otimes \mathcal{H}_2$ , where  $B = u_{\alpha} \otimes v_{\beta}$ . Then by the definition of inner product in Eq. 2,

$$\langle A, B \rangle = \sum_i \langle Av_i, Bv_i \rangle \quad (5)$$

$$= \sum_i \langle Av_i, (u_{\alpha} \otimes v_{\beta})v_i \rangle \quad (6)$$

$$= \sum_i \langle Av_i, \langle v_{\beta}, v_i \rangle u_{\alpha} \rangle \quad (\because \text{Lemma 1}) \quad (7)$$

$$= \langle Av_1, \langle v_{\beta}, v_1 \rangle u_{\alpha} \rangle + \langle Av_2, \langle v_{\beta}, v_2 \rangle u_{\alpha} \rangle + \dots + \langle Av_{\beta}, \langle v_{\beta}, v_{\beta} \rangle u_{\alpha} \rangle + \dots \quad (8)$$

$$= \langle Av_{\beta}, u_{\alpha} \rangle \quad (9)$$

27 Now we check the Parseval identity:

$$\sum_{\alpha, \beta} |\langle A, u_{\alpha} \otimes v_{\beta} \rangle|^2 = \sum_{\alpha, \beta} |\langle Av_{\beta}, u_{\alpha} \rangle|^2 \quad (\because u_{\alpha} \otimes v_{\beta} = B) \quad (10)$$

$$= \sum_{\beta} \|Av_{\beta}\|^2 \quad (11)$$

$$= \|A\|^2. \quad (12)$$

28  $\therefore \{u_{\alpha} \otimes v_{\beta}\}$  is a basis.

29

□

30 Now we begin the proof of Theorem 1. in the main paper.

31 **Theorem 1.** Let  $X, Y$  be compact subsets of  $\mathbb{R}^d$ . Choose  $u \in L^2(X \times Y)$ . Then, for arbitrary  $\varepsilon > 0$ ,  
 32 we can find a sufficiently large  $r > 0$  and neural networks  $f_j^{(\theta_1)}$  and  $f_j^{(\theta_2)}$  such that

$$\left\| u - \sum_{j=1}^r f_j^{(\theta_1)} f_j^{(\theta_2)} \right\|_{L^2(X \times Y)} < \varepsilon. \quad (13)$$

33

34 *Proof.* Let  $\{\phi_i\}$  and  $\{\psi_j\}$  be orthonormal basis for  $L^2(X)$  and  $L^2(Y)$  respectively. Then  $\{\phi_i\psi_j\}$   
 35 forms an orthonormal basis for  $L^2(X \times Y)$  ( $\because$  Lemma 2). Therefore, we can find a sufficiently large  
 36  $r$  such that

$$\left\| u - \sum_{i,j}^r a_{ij} \phi_i \psi_j \right\|_{L^2(X \times Y)} < \frac{\varepsilon}{2}, \quad (14)$$

where  $a_{ij}$  denotes

$$a_{ij} = \int_{X \times Y} u(x, y) \phi_i(x) \psi_j(y) dx dy.$$

37 On the other hand, by the universal approximation theorem [1], we can find neural networks  $f_j^{(\theta_1)}$   
 38 and  $f_j^{(\theta_2)}$  such that

$$\|\phi_i - f_j^{(\theta_1)}\|_{L^2(X)} \leq \frac{\varepsilon}{3^j \|u\|_{L^2(X \times Y)}} \quad \text{and} \quad \|\psi_j - f_j^{(\theta_2)}\|_{L^2(Y)} \leq \frac{\varepsilon}{3^j \|u\|_{L^2(X \times Y)}}. \quad (15)$$

39 We first consider the difference between  $u$  and  $\sum_{i,j}^r a_{ij} f_i^{(\theta_1)} f_j^{(\theta_2)}$ :

$$\left\| u - \sum_{i,j}^r a_{ij} f_i^{(\theta_1)} f_j^{(\theta_2)} \right\|_{L^2(X \times Y)} \quad (16)$$

$$\leq \left\| u - \sum_{i,j}^r a_{ij} \phi_i \psi_j \right\|_{L^2(X \times Y)} + \left\| \sum_{i,j}^r a_{ij} \phi_i \psi_j - \sum_{i,j}^r a_{ij} f_i^{(\theta_1)} f_j^{(\theta_2)} \right\|_{L^2(X \times Y)} \quad (17)$$

$$\equiv I + II \quad (18)$$

40 Since  $|I| < \varepsilon/2$  from (14), it is enough to estimate  $II$ . For this, we consider

$$\sum_{i,j}^r a_{ij} \phi_i \psi_j - \sum_{i,j}^r a_{ij} f_i^{(\theta_1)} f_j^{(\theta_2)} = \sum_{i,j}^r a_{ij} \phi_i (\psi_j - f_j^{(\theta_2)}) + \sum_{i,j}^r a_{ij} (\phi_i - f_i^{(\theta_1)}) f_j^{(\theta_2)} \quad (19)$$

$$\equiv II_1 + II_2. \quad (20)$$

41 We first compute  $II_1$ :

$$\|II_1\|_{L^2(X \times Y)}^2 = \int_{X \times Y} \left\{ \sum_{i,j}^r a_{ij} \phi_i (\psi_j - f_j^{(\theta_2)}) \right\}^2 dx dy \quad (21)$$

$$= \int_{X \times Y} \left\{ \sum_j^r \left( \sum_i^r a_{ij} \phi_i \right) (\psi_j - f_j^{(\theta_2)}) \right\}^2 dx dy \quad (22)$$

42 We set

$$A_j(x) = \sum_i^r a_{ij} \phi_i(x), \quad B_j(y) = \psi_j(y) - f_j^{(\theta_2)}(y) \quad (23)$$

43 to write  $II_1$  as

$$\|II_1\|_{L^2(X \times Y)}^2 = \int_{X \times Y} \left\{ \sum_j^r A_j(x) B_j(y) \right\}^2 dx dy. \quad (24)$$

44 We then apply Cauchy-Scharwz inequality to get

$$\|II_1\|_{L^2(X \times Y)}^2 \leq \int_{X \times Y} \left( \sum_j^r |A_j(x)|^2 \right) \left( \sum_j^r |B_j(y)|^2 \right) dx dy \quad (25)$$

$$= \left( \int_X \sum_j^r |A_j(x)|^2 dx \right) \left( \int_Y \sum_j^r |B_j(y)|^2 dy \right). \quad (26)$$

45 Now

$$\int_X \sum_j^r |A_j(x)|^2 dx = \int_X \sum_j^r \left( \sum_i^r a_{ij} \phi_i \right)^2 dx \quad (27)$$

$$= \sum_j^r \left\| \sum_i^r a_{ij} \phi_i \right\|_{L(X)}^2. \quad (28)$$

46 Since  $\{\phi_i\}$  is an orthonormal basis, we see that

$$\left\| \sum_i^r a_{ij} \phi_i \right\|_{L(X)} \leq \sum_i^r |a_{ij}| \|\phi_i\|_{L(X)} = \sum_i^r |a_{ij}|. \quad (29)$$

47 Therefore,

$$\int_X \sum_j^r |A_j(x)|^2 dx = \int_X \sum_j^r \left( \sum_i^r a_{ij} \phi_i \right)^2 dx \quad (30)$$

$$= \sum_j^r \int_X \left( \sum_i^r a_{ij} \phi_i \right)^2 dx \quad (31)$$

$$= \sum_j^r \left\| \sum_i^r a_{ij} \phi_i \right\|_{L(X)}^2 \quad (32)$$

$$= \sum_j^r \left( \sum_i^r |a_{ij}| \right)^2. \quad (33)$$

48 Finally, we recall

$$\left( \sum_i^r |a_{ij}| \right)^2 \leq 2 \sum_i^r |a_{ij}|^2 \quad (34)$$

49 to conclude

$$\int_X \sum_j^r |A_j(x)|^2 dx \leq 2 \sum_j^r \sum_i^r |a_{ij}|^2 < 2 \sum_{i,j}^{\infty} |a_{ij}|^2 = 2 \|u\|_{L^2(X \times Y)}^2. \quad (35)$$

50 On the other hand, we have from (15)

$$\int_Y \sum_j^r |B_j(y)|^2 dy = \sum_j^r \int_Y |\psi_j(y) - f_j^{(\theta_2)}(y)|^2 dy \quad (36)$$

$$= \sum_j^r \|\psi_j(y) - f_j^{(\theta_2)}(y)\|_{L^2(Y)}^2 \quad (37)$$

$$\leq \sum_j^r \frac{\varepsilon^2}{9^j \|u\|_{L^2(X \times Y)}^2} \quad (38)$$

$$< \frac{\varepsilon^2}{8 \|u\|_{L^2(X \times Y)}^2}. \quad (39)$$

51 Hence we have

$$\|II_1\|_{L^2(X \times Y)}^2 < \frac{\varepsilon^2}{8 \|u\|_{L^2(X \times Y)}^2} 2 \|u\|_{L^2(X \times Y)}^2 = \frac{\varepsilon^2}{4}. \quad (40)$$

52 Likewise,

$$\|II_2\|_{L^2(X \times Y)}^2 < \frac{\varepsilon^2}{4}. \quad (41)$$

53 Therefore,

$$\|II\|_{L^2(X \times Y)}^2 < \frac{\varepsilon^2}{2}. \quad (42)$$

54 We go back to (16) with the estimates (14) and (42) to derive

$$\left\| u - \sum_{i,j}^r a_{ij} f_i^{(\theta_1)} f_j^{(\theta_2)} \right\|_{L^2(X \times Y)} < \frac{\varepsilon}{2} + \frac{\varepsilon}{2} = \varepsilon. \quad (43)$$

55 Finally, we reorder the index to rewrite

$$\sum_{i,j}^r a_{ij} f_i^{(\theta_1)} f_j^{(\theta_2)} = \sum_i^{\tilde{r}} b_i \tilde{f}_i^{(\theta_1)} \tilde{f}_i^{(\theta_2)} \quad (44)$$

$$= \sum_i^{\tilde{r}} \left\{ b_i \tilde{f}_i^{(\theta_1)} \right\} \tilde{f}_i^{(\theta_2)} \quad (45)$$

$$= \sum_i^{\tilde{r}} g_i^{(\theta_1)} \tilde{f}_i^{(\theta_2)} \quad (46)$$

56 Without loss of generality, we rewrite  $\tilde{r}, g_i^{(\theta_1)}, \tilde{f}_i^{(\theta_2)}$  into  $r, f_i^{(\theta_1)}, f_i^{(\theta_2)}$  respectively, to complete the  
57 proof.  $\square$

## 58 D Training with Physics-Informed Loss

59 After SPINN predicts an output function with the methods described above, the rest of the training  
60 procedure follows the same process used in conventional PINN training [11], except we use forward-  
61 mode AD to compute PDE residuals (standard back-propagation, a.k.a. reverse-mode AD for  
62 parameter updates). With the slight abuse of notation, our predicted solution function is denoted as  
63  $\hat{u}^{(\theta)}(x, t)$  from onwards, explicitly expressing time coordinates. Given an underlying PDE (or ODE),  
64 the initial, and the boundary conditions, SPINN is trained with a ‘physics-informed’ loss function:

$$\min_{\theta} \mathcal{L}(\hat{u}^{(\theta)}(x, t)) = \min_{\theta} \lambda_{\text{pde}} \mathcal{L}_{\text{pde}} + \lambda_{\text{ic}} \mathcal{L}_{\text{ic}} + \lambda_{\text{bc}} \mathcal{L}_{\text{bc}}, \quad (47)$$

$$\mathcal{L}_{\text{pde}} = \int_{\Gamma} \int_{\Omega} \|\mathcal{N}[\hat{u}^{(\theta)}](x, t)\|^2 dx dt, \quad (48)$$

$$\mathcal{L}_{\text{ic}} = \int_{\Omega} \|\hat{u}^{(\theta)}(x, 0) - u_{\text{ic}}(x)\|^2 dx, \quad (49)$$

$$\mathcal{L}_{\text{bc}} = \int_{\Gamma} \int_{\partial\Omega} \|\mathcal{B}[\hat{u}^{(\theta)}](x, t) - u_{\text{bc}}(x, t)\|^2 dx dt, \quad (50)$$

65 where  $\Omega$  is an input domain,  $\mathcal{N}, \mathcal{B}$  are generic differential operators and  $u_{\text{ic}}, u_{\text{bc}}$  are initial, boundary  
66 conditions, respectively.  $\lambda$  are weighting factors for each loss term. When calculating the PDE  
67 loss ( $\mathcal{L}_{\text{pde}}$ ) with Monte-Carlo integral approximation, we sampled collocation points from factorized  
68 coordinates and used forward-mode AD. The remaining  $\mathcal{L}_{\text{ic}}$  and  $\mathcal{L}_{\text{bc}}$  are then computed with initial  
69 and boundary coordinates to regress the given conditions. By minimizing the objective loss in Eq. 47,  
70 the model output is enforced to satisfy the given equation, the initial, and the boundary conditions.

## 71 E FLOPs Estimation

72 The FLOPs for evaluating the derivatives can be systematically calculated by disassembling the  
73 computational graph into elementary operations such as additions and multiplications. Given a  
74 computational graph of forward pass for computing the primals, AD augments each elementary  
75 operation into other elementary operations. The FLOPs in the forward pass can be precisely calculated  
76 since it consists of a series of matrix multiplications and additions. We used the method described  
77 in [4] to estimate FLOPs for evaluating the derivatives. Table 1 shows the number of additions

78 (ADDS) and multiplications (MULTS) in each evaluation process. Note that FLOPs is a summation  
 79 of ADDS and MULTS by definition.

80 One thing to note here is that this is a theoretical estimation. Theoretically, the number of JVP  
 81 evaluations for computing the gradient with respect to the input coordinates is  $Nd$ , when  $N$  is  
 82 the number of coordinates for each axis and  $d$  is the input dimension (see section 4.3 in the main  
 83 paper). However, our actual implementation of gradient calculation involves re-computing the  
 84 feature representations  $f^{(\theta_i)}$ , which makes the complexity of network propagations from  $\mathcal{O}(Nd)$  to  
 85  $\mathcal{O}(Nd^2)$ . Ideally, these feature vectors can be computed only once and stored to be used later for  
 86 gradient computations. Although it is still significantly more efficient than the conventional PINN’s  
 87 complexity ( $Nd^2 \ll N^d$ ), there is a room to bridge the gap between theoretical FLOPs and actual  
 88 training runtime by further software optimization.

Table 1: The number of elementary operations for evaluating forward pass, first and second-order derivatives. The calculation is based on  $64^3$  collocation points in a 3-d system and the vanilla MLP settings used for diffusion, Helmholtz, and Klein-Gordon equations. We assumed that each derivative is evaluated on every coordinate axis.

	SPINN (ours)		PINN (baseline)	
	ADDS ( $\times 10^6$ )	MULTS ( $\times 10^6$ )	ADDS ( $\times 10^6$ )	MULTS ( $\times 10^6$ )
forward pass	20	20	21,609	21,609
1st-order derivative	40	40	86,638	43,419
2nd-order derivative	80	80	130,057	87,040
MFLOPs (total)	280		390,370	

## 89 F Experimental Details and Results

90 In this section, we provide experimental details, numerical results, and visualizations for each  
 91 experiment in the main paper.

### 92 F.1 Diffusion Equation

93 The diffusion equation is one of the most representative parabolic PDEs, often used for modeling the  
 94 heat diffusion process. We especially choose a nonlinear diffusion equation where it can be written  
 95 as:

$$\partial_t u = \alpha (\|\nabla u\|^2 + u\Delta u), \quad x \in \Omega, t \in \Gamma, \quad (51)$$

$$u(x, 0) = u_{ic}(x), \quad x \in \Omega, \quad (52)$$

$$u(x, t) = 0, \quad x \in \partial\Omega, t \in \Gamma. \quad (53)$$

96 We used diffusivity  $\alpha = 0.05$ , spatial domain  $\Omega = [-1, 1]^2$ , temporal domain  $\Gamma = [0, 1]$  and used  
 97 superposition of three Gaussian functions for the initial condition  $u_{ic}$ . We obtained the reference  
 98 solution ( $101 \times 101 \times 101$  resolution) through a widely-used PDE solver platform FEniCS [9]. Note  
 99 that FEniCS is a FEM-based solver. We particularly set the initial condition to be a superposition of  
 100 three gaussian functions:

$$u_{ic}(x, y) = 0.25 \exp[-10\{(x - 0.2)^2 + (y - 0.3)^2\}] + 0.4 \exp[-15\{(x + 0.1)^2 + (y + 0.5)^2\}] \\ + 0.3 \exp[-20\{(x + 0.5)^2 + y^2\}]. \quad (54)$$

101 For our model, we used three body networks of 4 hidden layers with 64/32 hidden feature/output  
 102 size each. For the baseline model, we used a single MLP of 5 hidden layers with 128 hidden feature  
 103 sizes. We used Adam [6] optimizer with a learning rate of 0.001 and trained for 50,000 iterations for  
 104 every experiment. All weight factors  $\lambda$  in the loss function in Eq. 47 are set to 1. The final reported  
 105 errors are extracted where the total loss was minimum across the entire training iteration. We also  
 106 resampled the input points every 100 epochs. Tab. 3 shows the numerical results, and the visualized  
 107 solutions are provided in Fig. 1.

108 **F.2 Helmholtz Equation**

109 The Helmholtz equation is a time-independent wave equation that takes the form:

$$\Delta u + k^2 u = q, \quad x \in \Omega, \quad (55)$$

$$u(x) = 0, \quad x \in \partial\Omega, \quad (56)$$

110 where the spatial domain is  $\Omega = [-1, 1]^3$ . For a given source term  $q = -(a_1\pi)^2 u - (a_2\pi)^2 u -$   
 111  $(a_3\pi)^2 u + k^2 u$ , we devised a manufactured solution  $u = \sin(a_1\pi x_1) \sin(a_2\pi x_2) \sin(a_3\pi x_3)$ , where  
 112 we take  $k = 1, a_1 = 4, a_2 = 4, a_3 = 3$ .

113 For our model, we used three body networks of 4 hidden layers with 64/32 hidden feature/output  
 114 size each. For the baseline model, we used a single MLP of 5 hidden layers with 128 hidden feature  
 115 sizes. We used Adam [6] optimizer with a learning rate of 0.001 and trained for 50,000 iterations for  
 116 every experiment. All weight factors  $\lambda$  in the loss function in Eq. 47 are set to 1. The final reported  
 117 errors are extracted where the total loss was minimum across the entire training iteration. We also  
 118 resampled the input points every 100 epochs. Tab. 4 shows the numerical results and the visualized  
 119 solutions are provided in Fig. 2.

120 **F.3 Klein-Gordon Equation**

121 The Klein-Gordon equation is a nonlinear hyperbolic PDE, which arises in diverse applied physics  
 122 for modeling relativistic wave propagation. The inhomogeneous Klein-Gordon equation is given by

$$\partial_{tt}u - \Delta u + u^2 = f, \quad x \in \Omega, t \in \Gamma, \quad (57)$$

$$u(x, 0) = x_1 + x_2, \quad x \in \Omega, \quad (58)$$

$$u(x, t) = u_{bc}(x), \quad x \in \partial\Omega, t \in \Gamma, \quad (59)$$

123 where we chose the spatial/temporal domain to be  $\Omega = [-1, 1]^2$  and  $\Gamma = [0, 10]$ , respectively. For  
 124 error measurement, we used a manufactured solution  $u = (x_1 + x_2) \cos(2t) + x_1 x_2 \sin(2t)$  and  $f,$   
 125  $u_{bc}$  are extracted from this exact solution.

126 For our model, we used three body networks of 4 hidden layers with 64/32 hidden feature/output  
 127 size each. For the baseline model, we used a single MLP of 5 hidden layers with 128 hidden feature  
 128 sizes. We used Adam [6] optimizer with a learning rate of 0.001 and trained for 50,000 iterations for  
 129 every experiment. All weight factors  $\lambda$  in the loss function in Eq. 47 are set to 1. The final reported  
 130 errors are extracted where the total loss was minimum across the entire training iteration. We also  
 131 resampled the input points every 100 epochs. Tab. 5 shows the numerical results.

132 We used the same settings used in (2+1)-d Klein-Gordon experiment for the (3+1)-d experiment  
 133 except the manufactured solution was chosen as:

$$u = (x_1 + x_2 + x_3) \cos(2t) + x_1 x_2 x_3 \sin(2t), \quad (60)$$

134 where  $f, u_{bc}$  are extracted from this exact solution. Tab. 6 shows the numerical results. The  
 135 collocation sizes of  $2^3, 18^4$  were the maximum value for PINN and PINN with modified MLP,  
 136 respectively.

137 **F.4 (2+1)-d Navier-Stokes Equation**

138 Navier-Stokes equation is a nonlinear time-dependent PDE that describes the motion of a viscous  
 139 fluid. Various engineering fields rely on this equation, such as modeling the weather, airflow, or ocean  
 140 currents. The vorticity form for incompressible fluid can be written as below:

$$\partial_t \omega + u \cdot \nabla \omega = \nu \Delta \omega, \quad x \in \Omega, t \in \Gamma, \quad (61)$$

$$\nabla \cdot u = 0, \quad x \in \Omega, t \in \Gamma, \quad (62)$$

$$\omega(x, 0) = \omega_0(x), \quad x \in \Omega, \quad (63)$$

141 where  $u \in \mathbb{R}^2$  is the velocity field,  $\omega = \nabla \times u$  is the vorticity,  $\omega_0$  is the initial vorticity, and  $\nu$  is  
 142 the viscosity. We used the viscosity 0.01 and made the spatial/temporal domain  $\Omega = [0, 2\pi]^2$  and  
 143  $\Gamma = [0, 1]$ , respectively. Note that Eq. 61 models decaying turbulence since there is no forcing term  
 144 and Eq. 62 is the incompressible fluid condition. The reference solution is generated by JAX-CFD

145 solver [7] which specifically used the pseudo-spectral method. The initial condition was generated  
 146 using the gaussian random field with a maximum velocity of 5. The resolution of the obtained  
 147 solution is  $100 \times 128 \times 128$  ( $N_t \times N_x \times N_y$ ), and we tested our model on this data.

148 For our model, we used three body networks (modified MLP) of 3 hidden layers with 128/256 hidden  
 149 feature/output sizes each. We divided the temporal domain into ten time windows to adopt the time  
 150 marching method [8, 13]. We used Adam [6] optimizer with a learning rate of 0.002 and each time  
 151 window is trained for 100,000 iterations. Followed by causal PINN, the PDE (residual) loss and  
 152 initial condition loss function are written as follows.

$$\mathcal{L}_{\text{pde}} = \frac{\lambda_w}{N_c} \sum^{N_c} |\partial_t w + u_x \partial_x w + u_y \partial_y w - \nu(\partial_{xx} w + \partial_{yy} w)|^2 + \frac{\lambda_c}{N_c} \sum^{N_c} |\partial_x u_x + \partial_y u_y|^2, \quad (64)$$

$$\mathcal{L}_{\text{ic}} = \frac{\lambda_{\text{ic}}}{N_{\text{ic}}} \sum^{N_{\text{ic}}} (|u_x - u_{x0}|^2 + |u_y - u_{y0}|^2 + |w - w_0|^2), \quad (65)$$

153 where  $u_x, u_y$  are  $x, y$  components of predicted velocity,  $w = \partial_x u_y - \partial_y u_x$ ,  $N_c$  is the collocation size,  
 154  $N_{\text{ic}}$  is the number of coordinates for initial condition and  $u_{x0}, u_{y0}, w_0$  are the initial conditions. We  
 155 chose the weighting factors  $\lambda_w = 1$ ,  $\lambda_c = 5,000$ , and  $\lambda_{\text{ic}} = 10,000$ . We also resampled the input  
 156 points every 100 epochs. The periodic boundary condition can be explicitly enforced by positional  
 157 encoding [2], and we specifically used the following encoding function only for the spatial input  
 158 coordinates.

$$\gamma(x) = [1, \sin(x), \sin(2x), \sin(3x), \sin(4x), \sin(5x), \cos(x), \cos(2x), \cos(3x), \cos(4x), \cos(5x)]^\top. \quad (66)$$

159 Unlike other experiments, the solution of the Navier-Stokes equation is a 2-dimensional vector-valued  
 160 function  $u : \mathbb{R}^3 \rightarrow \mathbb{R}^2$ . We can rewrite the feature merging equation Eq. 5 in the main paper to  
 161 construct SPINN into a 2-dimensional vector function:

$$u_1 = \sum_{j=1}^r \prod_{i=1}^d f_j^{(\theta_i)}(x_i), \quad (67)$$

$$u_2 = \sum_{j=r+1}^{2r} \prod_{i=1}^d f_j^{(\theta_i)}(x_i). \quad (68)$$

162 For example in our (2+1)-d Navier-Stokes equation setting,  $r = 128$  since the network output feature  
 163 size is 256. This can be applied to any  $m$ -dimensional vector function if we use a larger output feature  
 164 size. More formally, if we want to construct an  $m$ -dimensional vector function with SPINN of rank  $r$ ,  
 165 the  $k$ -th element of the function output can be written as

$$u_k = \sum_{j=(k-1)r+1}^{kr} \prod_{i=1}^d f_j^{(\theta_i)}(x_i), \quad (69)$$

166 where the output feature size of each body network is  $mr$ .

## 167 F.5 (3+1)-d Navier-Stokes Equation

168 The vorticity form of (3+1)-d Navier-Stokes equation is given as:

$$\partial_t \omega + (u \cdot \nabla) \omega = (\omega \cdot \nabla) u + \nu \Delta \omega + F, \quad x \in \Omega, t \in \Gamma, \quad (70)$$

$$\nabla \cdot u = 0, \quad x \in \Omega, t \in \Gamma, \quad (71)$$

$$\omega(x, 0) = \omega_0(x), \quad x \in \Omega. \quad (72)$$

169 We constructed the spatial/temporal domain to be  $\Omega = [0, 2\pi]^3$  and  $\Gamma = [0, 5]$ , respectively. We  
 170 constructed the analytic solution for (3+1)-d Navier-Stokes equation introduced in Taylor et al. [12].

171 The manufactured velocity and vorticity are

$$u_x = 2e^{-9\nu t} \cos(2x) \sin(2y) \sin(z), \quad (73)$$

$$u_y = -e^{-9\nu t} \sin(2x) \cos(2y) \sin(z), \quad (74)$$

$$u_z = -2e^{-9\nu t} \sin(2x) \sin(2y) \cos(z), \quad (75)$$

$$\omega_x = -3e^{-9\nu t} \sin(2x) \cos(2y) \cos(z), \quad (76)$$

$$\omega_y = 6e^{-9\nu t} \cos(2x) \sin(2y) \cos(z), \quad (77)$$

$$\omega_z = -6e^{-9\nu t} \cos(2x) \cos(2y) \sin(z), \quad (78)$$

172 where we chose the viscosity to be  $\nu = 0.05$ . Each forcing term  $F$  in the Eq. 70 is then given as

$$F_x = -6e^{-18\nu t} \sin(4y) \sin(2z), \quad (79)$$

$$F_y = -6e^{-18\nu t} \sin(4x) \sin(2z), \quad (80)$$

$$F_z = 6e^{-18\nu t} \sin(4x) \sin(4y). \quad (81)$$

173 We constructed SPINN to be three body networks (modified MLP) of 5 hidden layers with 64/384  
 174 hidden feature/output sizes each. We used Adam [6] optimizer with a learning rate of 0.001 and  
 175 trained for 50,000 iterations. The weight factors in the loss function in Eq. 48 are chosen as  $\lambda_{\text{pde}} = 1$ ,  
 176  $\lambda_{\text{ic}} = 10$ , and  $\lambda_{\text{bc}} = 1$ . We also weighted the incompressibility loss (Eq. 71) with 100. The visualized  
 177 solution vector field is shown in Fig. 4.

## 178 G Additional Experiments

### 179 G.1 (5+1)-d Heat Equation

180 We tested our model on (5+1)-d heat equation to verify the effectiveness of our model on higher  
 181 dimensional PDE:

$$\frac{\partial u(t, x)}{\partial t} = \Delta u(t, x), \quad x \in [-1, 1]^5, t \in [0, 1], \quad (82)$$

182 where the manufactured solution is chosen to be  $\|x\|^2 + 10t$ . When trained with  $8^6$  collocation points,  
 183 SPINN achieved a relative  $L_2$  error of 0.0074 within 2 minutes.

### 184 G.2 Fine Tuning with L-BFGS

185 We also conducted some experiments to explore the use of L-BFGS when training SPINN. We  
 186 found that training with Adam first and then fine-tuning with L-BFGS showed a slight increase in  
 187 accuracy. Note that this training strategy is used by other works [5, 10] and is known to be effective  
 188 in some cases. Tab. 2 shows the numerical results on three 3-d PDEs. Understanding the effect of the  
 189 optimization algorithm is still an open question in PINNs, we believe that investigating this issue in  
 190 the context of SPINN would be a valuable direction for future study.

Table 2: Numerical result of 3-d PDEs with L-BFGS fine tuning. The number of training collocation points is  $64^3$ .

	Diffusion	Helmholtz	Klein-Gordon
Adam	0.0041	0.0360	0.0013
Adam + L-BFGS	0.0041	0.0308	0.0010

191

Table 3: Full results of **diffusion equation**.  $N_c$  is the number of collocation points.

model	$N_c$	relative $L_2$ error	runtime (ms/iter.)	memory (MB)
PINN	$16^3$	0.0095	3.98	1,022
	$32^3$	0.0082	12.82	2,942
	$64^3$	0.0081	95.22	18,122
PINN + modified MLP	$16^3$	0.0048	14.94	1,918
	$32^3$	0.0043	29.91	4,990
	$54^3$	0.0041	134.64	22,248
SPINN	$16^3$	0.0447	1.45	766
	$32^3$	0.0115	1.76	766
	$64^3$	0.0075	1.90	766
	$128^3$	0.0061	2.09	894
	$256^3$	0.0061	10.54	2,174
SPINN + modified MLP	$16^3$	0.0390	2.17	766
	$32^3$	0.0067	2.44	768
	$64^3$	0.0041	2.59	768
	$128^3$	<b>0.0036</b>	3.06	896
	$256^3$	<b>0.0036</b>	12.13	2,176

Table 4: Full results of the **Helmholtz equation**.  $N_c$  is the number of collocation points.

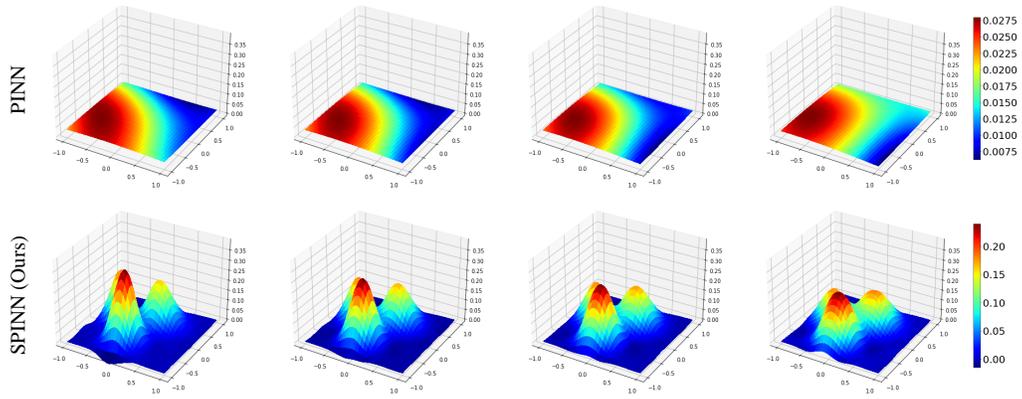
model	$N_c$	relative $L_2$ error	runtime (ms/iter.)	memory (MB)
PINN	$16^3$	0.9819	4.84	2,810
	$32^3$	0.9757	14.84	2,938
	$64^3$	0.9723	110.23	18,118
PINN + modified MLP	$16^3$	0.4770	18.32	7,034
	$32^3$	0.5176	35.02	9,082
	$54^3$	0.4770	159.90	22,244
SPINN	$16^3$	0.1177	1.54	762
	$32^3$	0.0809	1.71	762
	$64^3$	0.0592	1.85	762
	$128^3$	0.0449	1.89	762
	$256^3$	0.0435	3.84	1,146
SPINN + modified MLP	$16^3$	0.1161	2.24	764
	$32^3$	0.0595	2.50	764
	$64^3$	0.0360	2.57	764
	$128^3$	<b>0.0300</b>	2.76	764
	$256^3$	0.0311	5.50	1,148

Table 5: Full results of the **(2+1)-d Klein-Gordon equation**.  $N_c$  is the number of collocation points.

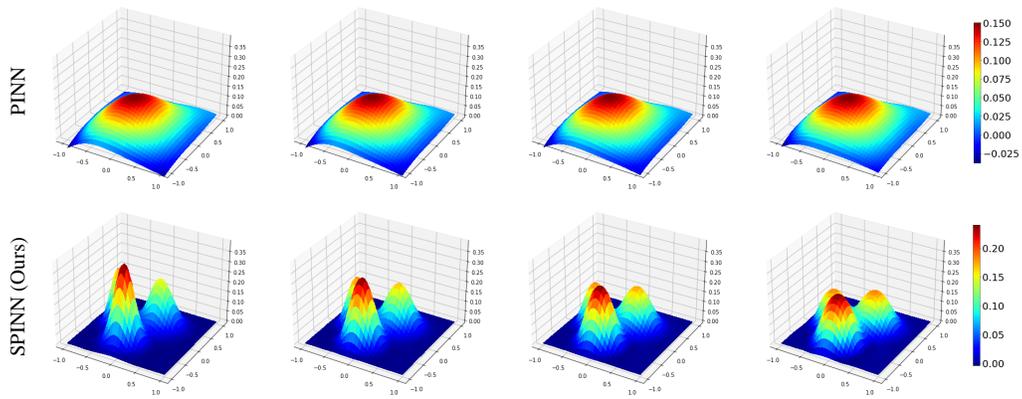
model	$N_c$	relative $L_2$ error	runtime (ms/iter.)	memory (MB)
PINN	$16^3$	0.0343	4.70	2,810
	$32^3$	0.0281	14.95	2,938
	$64^3$	0.0299	112.00	18,118
PINN + modified MLP	$16^3$	0.0158	17.87	7,036
	$32^3$	0.0185	34.61	9,082
	$54^3$	0.0163	159.20	22,246
SPINN	$16^3$	0.0193	1.55	762
	$32^3$	0.0060	1.71	762
	$64^3$	0.0045	1.82	762
	$128^3$	0.0040	1.85	890
	$256^3$	0.0039	3.98	1,658
SPINN + modified MLP	$16^3$	0.0062	2.20	764
	$32^3$	0.0020	2.41	764
	$64^3$	0.0013	2.57	764
	$128^3$	<b>0.0008</b>	2.79	892
	$256^3$	0.0009	5.61	1,660

Table 6: Full results of the **(3+1)-d Klein-Gordon equation**.  $N_c$  is the number of collocation points.

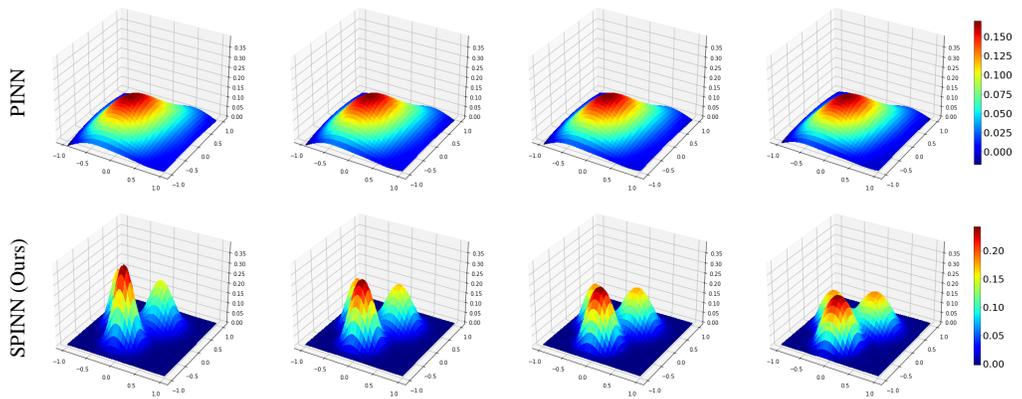
model	$N_c$	relative $L_2$ error	runtime (ms/iter.)	memory (MB)
PINN	$16^4$	0.0129	43.51	5,246
	$23^4$	0.0121	154.24	22,244
PINN + modified MLP	$16^4$	0.0061	100.06	17,534
	$18^4$	0.0059	174.00	22,246
SPINN	$16^4$	0.0122	2.45	890
	$32^4$	0.0095	2.98	892
	$64^4$	0.0093	9.22	2,172
SPINN + modified MLP	$16^4$	0.0064	3.48	892
	$32^4$	0.0022	3.66	892
	$64^4$	<b>0.0012</b>	10.96	2,172



(a) Training time: 5 sec



(b) Training time: 30 sec



(c) Training time: 1 min

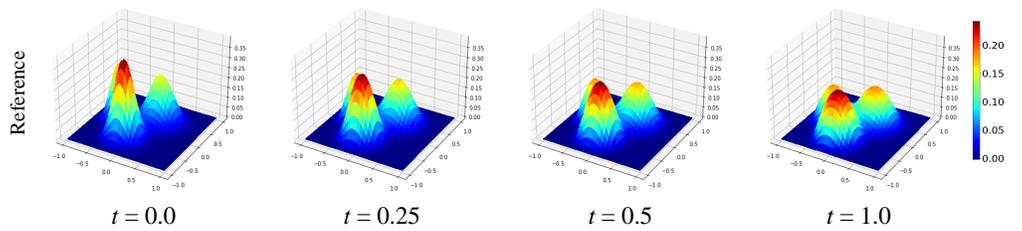


Figure 1: Visualized solution of **nonlinear diffusion equation** obtained by the baseline PINN and SPINN, both trained on  $64^3$  collocation points.

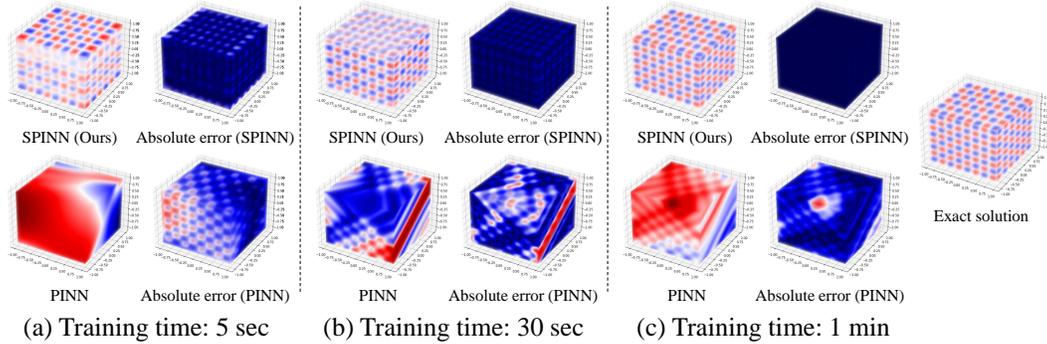


Figure 2: Visualized solution of **Helmholtz equation** obtained by the baseline PINN and SPINN, both trained on  $64^3$  collocation points.

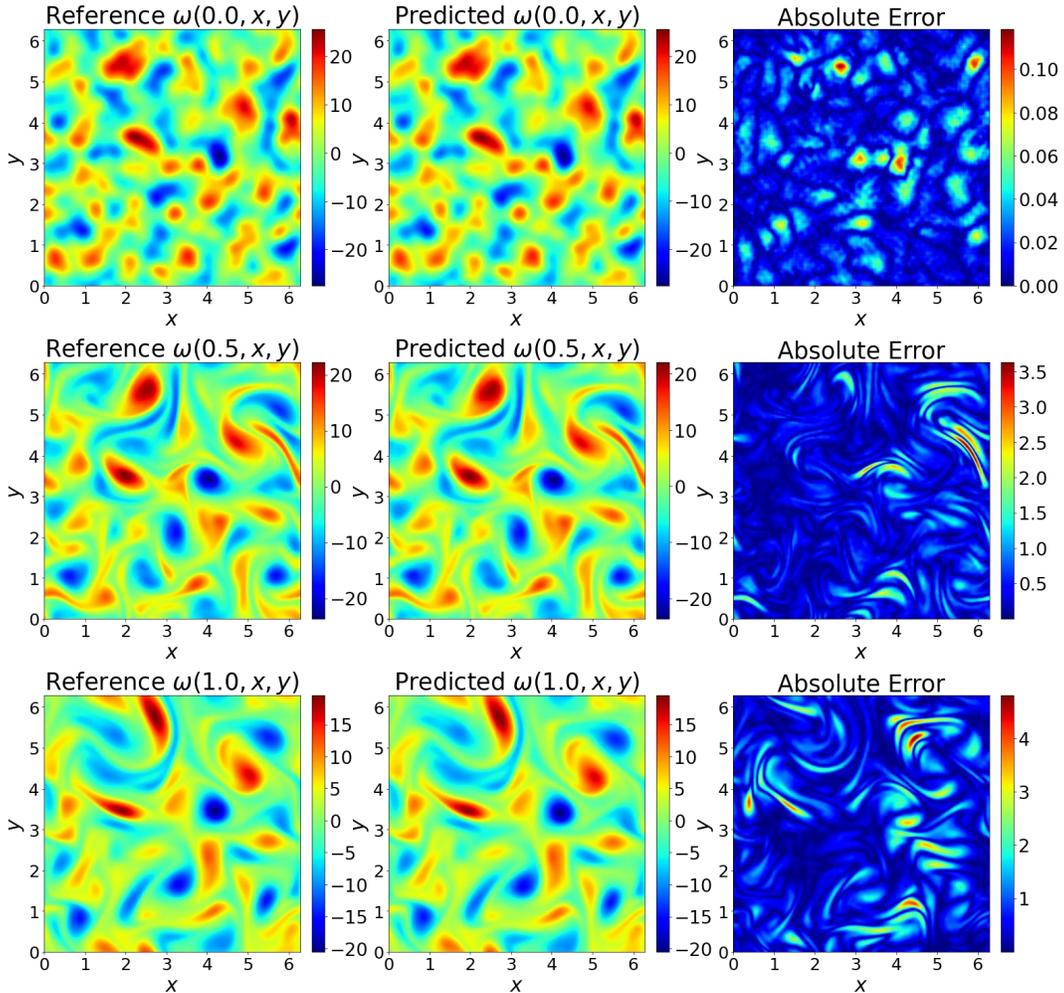


Figure 3: Visualized vorticity maps of **(2+1)-d Navier-Stokes equation** experiment predicted by SPINN. Three snapshots at timestamps  $t = 0.0, 0.5, 1.0$  are presented.

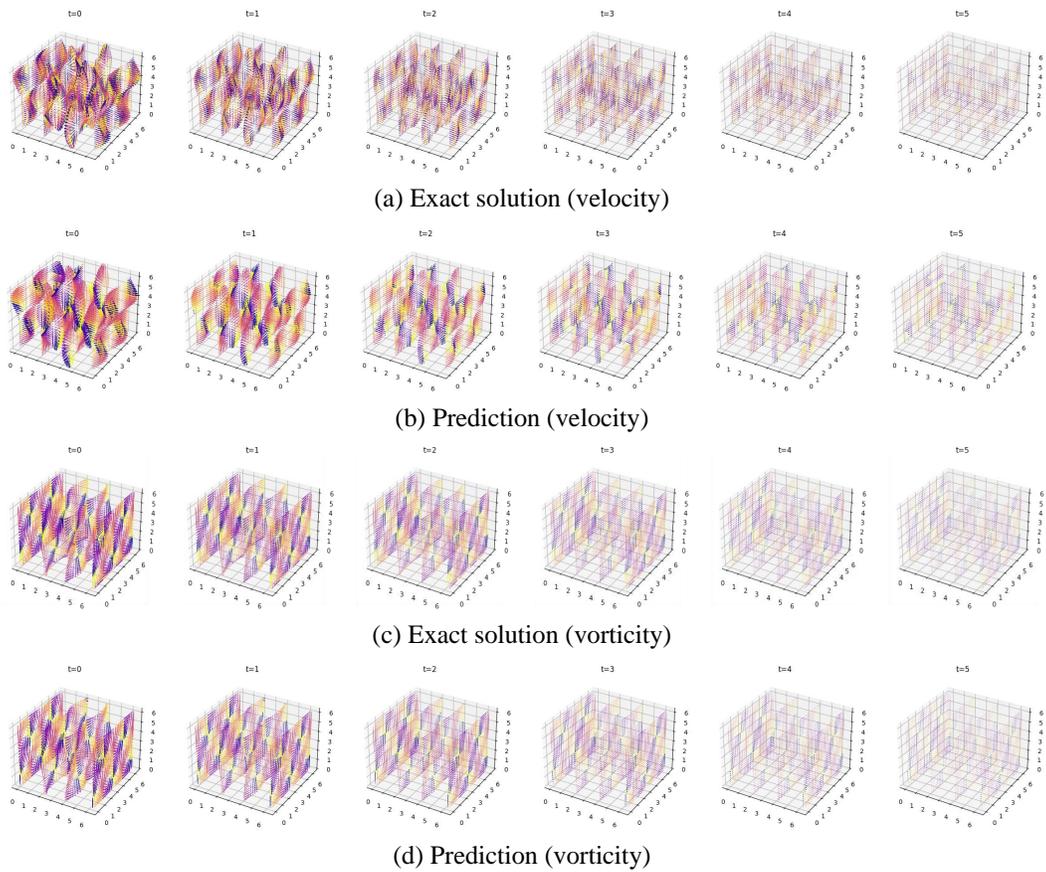


Figure 4: Visualized solution of **(3+1)-d Navier-Stokes equation** obtained by SPINN, trained on  $32^4$  collocation points. Each arrow is colored by its zenith angle.

## References

- 192
- 193 [1] George Cybenko. Approximation by superpositions of a sigmoidal function. *Mathematics of*  
194 *control, signals and systems*, 2(4):303–314, 1989.
- 195 [2] Suchuan Dong and Naxian Ni. A method for representing periodic functions and enforcing  
196 exactly periodic boundary conditions with deep neural networks. *Journal of Computational*  
197 *Physics*, 435:110242, 2021.
- 198 [3] Gerald B Folland. *A course in abstract harmonic analysis*, volume 29. CRC press, 2016.
- 199 [4] Andreas Griewank and Andrea Walther. *Evaluating derivatives: principles and techniques of*  
200 *algorithmic differentiation*. SIAM, 2008.
- 201 [5] Xiaowei Jin, Shengze Cai, Hui Li, and George Em Karniadakis. Nsfnets (navier-stokes flow  
202 nets): Physics-informed neural networks for the incompressible navier-stokes equations. *Journal*  
203 *of Computational Physics*, 426:109951, 2021.
- 204 [6] Diederik P. Kingma and Jimmy Ba. Adam: A method for stochastic optimization. In *Internat-*  
205 *ional Conference on Learning Representations*, 2015.
- 206 [7] Dmitrii Kochkov, Jamie A Smith, Ayya Alieva, Qing Wang, Michael P Brenner, and Stephan  
207 Hoyer. Machine learning–accelerated computational fluid dynamics. *Proceedings of the*  
208 *National Academy of Sciences*, 118(21):e2101784118, 2021.
- 209 [8] Aditi Krishnapriyan, Amir Gholami, Shandian Zhe, Robert Kirby, and Michael W Mahoney.  
210 Characterizing possible failure modes in physics-informed neural networks. *Advances in Neural*  
211 *Information Processing Systems*, 34:26548–26560, 2021.
- 212 [9] Anders Logg, Kent-Andre Mardal, and Garth Wells. *Automated solution of differential equations*  
213 *by the finite element method: The FEniCS book*, volume 84. Springer Science & Business  
214 Media, 2012.
- 215 [10] Xuhui Meng, Zhen Li, Dongkun Zhang, and George Em Karniadakis. Ppinn: Parareal physics-  
216 informed neural network for time-dependent pdes. *Computer Methods in Applied Mechanics*  
217 *and Engineering*, 370:113250, 2020.
- 218 [11] Maziar Raissi, Paris Perdikaris, and George E Karniadakis. Physics-informed neural networks:  
219 A deep learning framework for solving forward and inverse problems involving nonlinear partial  
220 differential equations. *Journal of Computational physics*, 378:686–707, 2019.
- 221 [12] Geoffrey Ingram Taylor and Albert Edward Green. Mechanism of the production of small  
222 eddies from large ones. *Proceedings of the Royal Society of London. Series A-Mathematical*  
223 *and Physical Sciences*, 158(895):499–521, 1937.
- 224 [13] Colby L Wight and Jia Zhao. Solving allen-cahn and cahn-hilliard equations using the adaptive  
225 physics informed neural networks. *arXiv preprint arXiv:2007.04542*, 2020.