On the Mode-Seeking Properties of Langevin Dynamics

Anonymous Author(s) Affiliation Address email

Abstract

The Langevin Dynamics framework, which aims to generate samples from the 1 score function of a probability distribution, is widely used for analyzing and 2 interpreting score-based generative modeling. While the convergence behavior of 3 Langevin Dynamics under unimodal distributions has been extensively studied in 4 the literature, in practice the data distribution could consist of multiple distinct 5 modes. In this work, we investigate Langevin Dynamics in producing samples 6 from multimodal distributions and theoretically study its mode-seeking properties. 7 8 We prove that under a variety of sub-Gaussian mixtures, Langevin Dynamics is unlikely to find all mixture components within a sub-exponential number of steps in 9 the data dimension. To reduce the mode-seeking tendencies of Langevin Dynamics, 10 we propose Chained Langevin Dynamics, which divides the data vector into patches 11 of constant size and generates every patch sequentially conditioned on the previous 12 patches. We perform a theoretical analysis of Chained Langevin Dynamics by 13 reducing it to sampling from a constant-dimensional distribution. We present 14 the results of several numerical experiments on synthetic and real image datasets, 15 supporting our theoretical results on the iteration complexities of sample generation 16 17 from mixture distributions using the chained and vanilla Langevin Dynamics.

18 1 Introduction

A central task in unsupervised learning involves learning the underlying probability distribution of 19 training data and efficiently generating new samples from the distribution. Score-based generative 20 modeling (SGM) (Song et al., 2020c) has achieved state-of-the-art performance in various learning 21 tasks including image generation (Song and Ermon, 2019, 2020; Ho et al., 2020; Song et al., 2020a; 22 Ramesh et al., 2022; Rombach et al., 2022), audio synthesis (Chen et al., 2020; Kong et al., 2020), 23 and video generation (Ho et al., 2022; Blattmann et al., 2023). In addition to the successful empirical 24 results, the convergence analysis of SGM has attracted significant attention in the recent literature 25 (Lee et al., 2022, 2023; Chen et al., 2023; Li et al., 2023, 2024). 26

Stochastic gradient Langevin dynamics (SGLD) (Welling and Teh, 2011), as a fundamental method-27 ology to implement and interpret SGM, can produce samples from the (Stein) score function of a 28 probability density, i.e., the gradient of the log probability density function with respect to data. It 29 has been widely recognized that a pitfall of SGLD is its slow mixing rate (Wooddard et al., 2009; 30 31 Raginsky et al., 2017; Lee et al., 2018). Specifically, Song and Ermon (2019) shows that under a multi-modal data distribution, the samples from Langevin dynamics may have an incorrect relative 32 density across the modes. Based on this finding, Song and Ermon (2019) proposes anneal Langevin 33 *dynamics*, which injects different levels of Gaussian noise into the data distribution and samples with 34 SGLD on the perturbed distribution. While outputting the correct relative density across modes can 35 36 be challenging for SGLD, a natural question is whether SGLD would be able to find all the modes of a multi-modal distribution. 37

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In this work, we study this question by analyzing the mode-seeking properties of SGLD. The notion 38 of mode-seekingness (Bishop, 2006; Ke et al., 2021; Li and Farnia, 2023) refers to the property that a 39 generative model captures only a subset of the modes of a multi-modal distribution. We note that 40 a similar problem, known as metastability, has been studied in the context of Langevin diffusion, 41 a continuous-time version of SGLD described by stochastic differential equation (SDE) (Bovier 42 et al., 2002, 2004; Gayrard et al., 2005). Specifically, Bovier et al. (2002) gave a sharp bound on 43 44 the mean hitting time of Langevin diffusion and proved that it may require exponential (in the space dimensionality d) time for transition between modes. Regarding discrete SGLD, Lee et al. (2018) 45 constructed a probability distribution whose density is close to a mixture of two well-separated 46 isotropic Gaussians, and proved that SGLD could not find one of the two modes within an exponential 47 number of steps. However, further exploration of mode-seeking tendencies of SGLD and its variants 48 such as annealed Langevin dynamics for general distributions is still lacking in the literature. 49

In this work, we theoretically formulate and demonstrate the potential mode-seeking tendency of 50 SGLD. We begin by analyzing the convergence under a variety of Gaussian mixture probability 51 distributions, under which SGLD could fail to visit all the mixture components within sub-exponential 52 steps (in the data dimension). Subsequently, we generalize this result to mixture distributions with 53 sub-Gaussian modes. This generalization extends our earlier result on Gaussian mixtures to a 54 significantly larger family of mixture models, as the sub-Gaussian family includes any distribution 55 over an ℓ_2 -norm-bounded support set. Furthermore, we extend our theoretical results to anneal 56 Langevin dynamics with bounded noise scales. 57

To reduce SGLD's large iteration complexity shown under a high-dimensional input vector, we 58 propose Chained Langevin Dynamics (Chained-LD). Since SGLD could suffer from the curse of 59 dimensionality, we decompose the sample $\mathbf{x} \in \mathbb{R}^d$ into d/Q patches $\mathbf{x}^{(1)}, \cdots, \mathbf{x}^{(d/Q)}$, each of 60 constant size Q, and sequentially generate every patch $\mathbf{x}^{(q)}$ for all $q \in [d/Q]$ statistically conditioned 61 on previous patches, i.e., $P(\mathbf{x}^{(q)} \mid \mathbf{x}^{(0)}, \cdots \mathbf{x}^{(q-1)})$. The combination of all patches generated from 62 63 the conditional distribution faithfully follows the probability density $P(\mathbf{x})$, while learning each patch 64 requires less cost due to the reduced dimension. We also provide a theoretical analysis of Chained-LD 65 by reducing the convergence of a *d*-dimensional sample to the convergence of each patch.

Finally, we present the results of several numerical experiments to validate our theoretical findings. 66 For synthetic experiments, we consider moderately high-dimensional Gaussian mixture models, 67 where the vanilla and annealed Langevin dynamics could not find all the components within a million 68 69 steps, while Chained-LD could capture all the components with correct frequencies in $\mathcal{O}(10^4)$ steps. 70 For experiments on real image datasets, we consider a mixture of two modes by using the original images from MNIST/Fashion-MNIST training dataset (black background and white digits/objects) 71 as the first mode and constructing the second mode by i.i.d. flipping the images (white background 72 and black digits/objects) with probability 0.5. Following from Song and Ermon (2019), we trained 73 a Noise Conditional Score Network (NCSN) to estimate the score function. Our numerical results 74 indicate that vanilla Langevin dynamics can fail to capture the two modes, as also observed by Song 75 and Ermon (2019). On the other hand, Chained-LD was capable of finding both modes regardless of 76 initialization. We summarize the contributions of this work as follows: 77

⁷⁸ • Theoretically studying the mode-seeking properties of vanilla and annealed Langevin dynamics,

Proposing Chained Langevin Dynamics (Chained-LD), which decomposes the sample into patches
 and sequentially generates each patch conditioned on previous patches,

• Providing a theoretical analysis of the convergence behavior of Chained-LD,

Numerically comparing the mode-seeking properties of vanilla, annealed, and chained Langevin dynamics.

Notations: We use [n] to denote the set $\{1, 2, \dots, n\}$. Also, in the paper, $\|\cdot\|$ refers to the ℓ_2 norm. We use $\mathbf{0}_n$ and $\mathbf{1}_n$ to denote a 0-vector and 1-vector of length n. We use I_n to denote the identity matrix of size $n \times n$. In the text, TV stands for the total variation distance.

87 2 Related Works

Langevin Dynamics: The convergence guarantees for Langevin diffusion, a continuous version of

⁸⁹ Langevin dynamics, are classical results extensively studied in the literature (Bhattacharya, 1978;

Roberts and Tweedie, 1996; Bakry and Émery, 1983; Bakry et al., 2008). Langevin dynamics, also 90 known as Langevin Monte Carlo, is a discretization of Langevin diffusion typically modeled as a 91 Markov Chain Monte Carlo (Welling and Teh, 2011). For unimodal distributions, e.g., the probability 92 density function that is log-concave or satisfies log-Sobolev inequality, the convergence of Langevin 93 dynamics is provably fast (Dalalyan, 2017; Durmus and Moulines, 2017; Vempala and Wibisono, 94 95 2019). However, for multimodal distributions, the non-asymptotic convergence analysis is much more challenging (Cheng et al., 2018). Raginsky et al. (2017) gave an upper bound on the convergence time 96 of Langevin dynamics for arbitrary non-log-concave distributions with certain regularity assumptions, 97 which, however, could be exponentially large without imposing more restrictive assumptions. Lee 98 et al. (2018) studied the special case of a mixture of Gaussians of equal variance and provided 99 heuristic analysis of sampling from general non-log-concave distributions. 100 Mode-Seekingness of Langevin Dynamics: The investigation of the mode-seekingness of gener-101 ative models starts with different generative adversarial network (GAN) (Goodfellow et al., 2014) 102 model formulations and divergence measures, from both the practical (Goodfellow, 2016; Poole 103 et al., 2016) and theoretical (Shannon et al., 2020; Li and Farnia, 2023) perspectives. In the context 104 of Langevin dynamics, mode-seekingness is closely related to a lower bound on the transition time

of Langevin dynamics, mode-seekingness is closely related to a lower bound on the transition time
 between two modes, e.g., two local maximums. Bovier et al. (2002, 2004); Gayrard et al. (2005)
 studied the mean hitting time of the continuous Langevin diffusion. Lee et al. (2018) proved the
 existence of a mixture of two Gaussian distributions whose covariance matrices differ by a constant
 factor, Langevin dynamics cannot find both modes in polynomial time.

Score-based Generative Modeling: Since Song et al. (2020b) proposed sliced score matching 110 which can train deep models to learn the score functions of implicit probability distributions on high-111 dimensional data, score-based generative modeling (SGM) has been going through a spurt of growth. 112 Annealed Langevin dynamics (Song and Ermon, 2019) estimates the noise score of the probability 113 density perturbed by Gaussian noise and utilizes stochastic gradient Langevin dynamics to generate 114 samples from a sequence of decreasing noise scales. Song and Ermon (2020) conducted a heuristic 115 analysis of the effect of noise levels on the performance of annealed Langevin dynamics. Denoising 116 diffusion probabilistic model (DDPM) (Ho et al., 2020) incorporates a step-by-step introduction of 117 random noise into data, followed by learning to reverse this diffusion process in order to generate 118 desired data samples from the noise. Song et al. (2020c) unified anneal Langevin dynamics and 119 DDPM via a stochastic differential equation. A recent line of work focuses on the non-asymptotic 120 convergence guarantees for SGM with an imperfect score estimation under various assumptions on 121 the data distribution (Block et al., 2020; De Bortoli et al., 2021; Lee et al., 2022; Chen et al., 2023; 122 Benton et al., 2023; Li et al., 2023, 2024). 123

124 **3** Preliminaries

125 3.1 Langevin Dynamics

Generative modeling aims to produce samples such that their distribution is close to the underlying true distribution P. For a continuously differentiable probability density $P(\mathbf{x})$ on \mathbb{R}^d , its score function is defined as the gradient of the log probability density function (PDF) $\nabla_{\mathbf{x}} \log P(\mathbf{x})$. Langevin diffusion is a stochastic process defined by the stochastic differential equation (SDE)

$$\mathrm{d}\mathbf{x}_t = -\nabla_{\mathbf{x}} \log P(\mathbf{x}_t) \,\mathrm{d}t + \sqrt{2} \,\mathrm{d}\mathbf{w}_t$$

where \mathbf{w}_t is the Wiener process on \mathbb{R}^d . To generate samples from Langevin diffusion, Welling and Teh (2011) proposed stochastic gradient Langevin dynamics (SGLD), a discretization of the SDE for *T* iterations. Each iteration of SGLD is defined as

$$\mathbf{x}_{t} = \mathbf{x}_{t-1} + \frac{\delta_{t}}{2} \nabla_{\mathbf{x}} \log P(\mathbf{x}_{t-1}) + \sqrt{\delta_{t}} \boldsymbol{\epsilon}_{t}, \tag{1}$$

where δ_t is the step size and $\epsilon_t \sim \mathcal{N}(\mathbf{0}_d, \mathbf{I}_d)$ is Gaussian noise. It has been widely recognized that Langevin diffusion could take exponential time to mix without additional assumptions on the probability density (Bovier et al., 2002, 2004; Gayrard et al., 2005; Raginsky et al., 2017; Lee et al., 2018). To combat the slow mixing, Song and Ermon (2019) proposed annealed Langevin dynamics by perturbing the probability density with Gaussian noise of variance σ^2 , i.e.,

$$P_{\sigma}(\mathbf{x}) := \int P(\mathbf{z}) \mathcal{N}(\mathbf{x} \mid \mathbf{z}, \sigma^2 \mathbf{I}_d) \, \mathrm{d}\mathbf{z},$$
(2)

and running SGLD on the perturbed data distribution $P_{\sigma_t}(\mathbf{x})$ with gradually decreasing noise levels $\{\sigma_t\}_{t\in[T]}$, i.e.,

$$\mathbf{x}_{t} = \mathbf{x}_{t-1} + \frac{\delta_{t}}{2} \nabla_{\mathbf{x}} \log P_{\sigma_{t}}(\mathbf{x}_{t-1}) + \sqrt{\delta_{t}} \boldsymbol{\epsilon}_{t},$$
(3)

where δ_t is the step size and $\epsilon_t \sim \mathcal{N}(\mathbf{0}_d, \mathbf{I}_d)$ is Gaussian noise. When the noise level σ is vanishingly small, the perturbed distribution is close to the true distribution, i.e., $P_{\sigma}(\mathbf{x}) \approx P(\mathbf{x})$. Since we do not have direct access to the (perturbed) score function, Song and Ermon (2019) proposed the Noise Conditional Score Network (NCSN) $\mathbf{s}_{\theta}(\mathbf{x}, \sigma)$ to jointly estimate the scores of all perturbed data distributions, i.e.,

$$\forall \sigma \in \{\sigma_t\}_{t \in [T]}, \ \mathbf{s}_{\theta}(\mathbf{x}, \sigma) \approx \nabla_{\mathbf{x}} \log P_{\sigma}(\mathbf{x}).$$

To train the NCSN, Song and Ermon (2019) adopted denoising score matching, which minimizes the
 following loss

$$\mathcal{L}\left(\boldsymbol{\theta}; \{\sigma_t\}_{t\in[T]}\right) := \frac{1}{2T} \sum_{t\in[T]} \sigma_t^2 \mathbb{E}_{\mathbf{x}\sim P} \mathbb{E}_{\tilde{\mathbf{x}}\sim\mathcal{N}(\mathbf{x},\sigma_t^2 \boldsymbol{I}_d)} \left[\left\| \mathbf{s}_{\boldsymbol{\theta}}(\tilde{\mathbf{x}},\sigma_t) - \frac{\tilde{\mathbf{x}} - \mathbf{x}}{\sigma_t^2} \right\|^2 \right].$$

Assuming the NCSN has enough capacity, $\mathbf{s}_{\theta^*}(\mathbf{x}, \sigma)$ minimizes the loss $\mathcal{L}\left(\boldsymbol{\theta}; \{\sigma_t\}_{t\in[T]}\right)$ if and only if $\mathbf{s}_{\theta^*}(\mathbf{x}, \sigma_t) = \nabla_{\mathbf{x}} \log P_{\sigma_t}(\mathbf{x})$ almost surely for all $t \in [T]$.

149 3.2 Multi-Modal Distributions

Our work focuses on multi-modal distributions. We use $P = \sum_{i \in [k]} w_i P^{(i)}$ to represent a mixture of k modes, where each mode $P^{(i)}$ is a probability density with frequency w_i such that $w_i > 0$ for all $i \in [k]$ and $\sum_{i \in [k]} w_i = 1$. In our theoretical analysis, we consider Gaussian mixtures and sub-Gaussian mixtures, i.e., every component $P^{(i)}$ is a Gaussian or sub-Gaussian distribution. A probability distribution $p(\mathbf{z})$ of dimension d is defined as a sub-Gaussian distribution with parameter ν^2 if, given the mean vector $\boldsymbol{\mu} := \mathbb{E}_{\mathbf{z} \sim p}[\mathbf{z}]$, the moment generating function (MGF) of p satisfies the following inequality for every vector $\boldsymbol{\alpha} \in \mathbb{R}^d$:

$$\mathbb{E}_{\mathbf{z}\sim p}\left[\exp\left(\boldsymbol{\alpha}^{T}(\mathbf{z}-\boldsymbol{\mu})\right] \leq \exp\left(\frac{\nu^{2} \|\boldsymbol{\alpha}\|_{2}^{2}}{2}\right).$$
(4)

We remark that sub-Gaussian distributions include a wide variety of distributions such as Gaussian distributions and any distribution within a bounded ℓ_2 -norm distance from the mean μ . From equation 2 we note that the perturbed distribution is the convolution of the original distribution and a Gaussian random variable, i.e., for random variables $\mathbf{z} \sim p$ and $\mathbf{t} \sim \mathcal{N}(\mathbf{0}_d, \mathbf{I}_d)$, their sum $\mathbf{z} + \mathbf{t} \sim p_{\sigma}$ follows the perturbed distribution with noise level σ . Therefore, a perturbed (sub)Gaussian distribution remains (sub)Gaussian. We formalize this property in Proposition 1 and defer the proof to Appendix A for completeness.

Proposition 1. Suppose the perturbed distribution of a d-dimensional probability distribution p with noise level σ is p_{σ} , then the mean of the perturbed distribution is the same as the original distribution, i.e., $\mathbb{E}_{\mathbf{z}\sim p_{\sigma}}[\mathbf{z}] = \mathbb{E}_{\mathbf{z}\sim p}[\mathbf{z}]$. If $p = \mathcal{N}(\boldsymbol{\mu}, \boldsymbol{\Sigma})$ is a Gaussian distribution, $p_{\sigma} = \mathcal{N}(\boldsymbol{\mu}, \boldsymbol{\Sigma} + \sigma^2 \mathbf{I}_d)$ is also a Gaussian distribution. If p is a sub-Gaussian distribution with parameter ν^2 , p_{σ} is a sub-Gaussian distribution with parameter $(\nu^2 + \sigma^2)$.

¹⁶⁹ 4 Theoretical Analysis of the Mode-Seeking Properties of Langevin Dynamics

In this section, we theoretically investigate the mode-seeking properties of vanilla and annealed
 Langevin dynamics. We begin with analyzing Langevin dynamics in Gaussian mixtures.

172 4.1 Langevin Dynamics in Gaussian Mixtures

Assumption 1. Consider a data distribution $P := \sum_{i=0}^{k} w_i P^{(i)}$ as a mixture of Gaussian distributions, where $1 \le k = o(d)$ and $w_i > 0$ is a positive constant such that $\sum_{i=0}^{k} w_i = 1$. Suppose that $P^{(i)} = \mathcal{N}(\boldsymbol{\mu}_i, \nu_i^2 \boldsymbol{I}_d)$ is a Gaussian distribution over \mathbb{R}^d for all $i \in \{0\} \cup [k]$ such that for all $i \in [k]$, $\nu_i < \nu_0$ and $\|\boldsymbol{\mu}_i - \boldsymbol{\mu}_0\|^2 \le \frac{\nu_0^2 - \nu_i^2}{2} \left(\log \left(\frac{\nu_i^2}{\nu_0^2} \right) - \frac{\nu_i^2}{2\nu_i^2} + \frac{\nu_0^2}{2\nu_i^2} \right) d$. Denote $\nu_{\max} := \max_{i \in [k]} \nu_i$.

Regarding the first requirement $\nu_i < \nu_0$, we first note that the probability density $p(\mathbf{z})$ of a Gaussian 177 distribution $\mathcal{N}(\boldsymbol{\mu}, \nu^2 \boldsymbol{I}_d)$ decays exponentially in terms of $\frac{\|\mathbf{z}-\boldsymbol{\mu}\|^2}{|\mathbf{z}|^2}$. When a state \mathbf{z} is sufficiently far 178 from all modes (i.e., $\|\mathbf{z}\| \gg \|\boldsymbol{\mu}_i\|$), the Gaussian distribution with the largest variance (i.e., $P^{(0)}$ in 179 Assumption 1) dominates all other modes because $\frac{\|\mathbf{z}-\boldsymbol{\mu}_0\|^2}{\nu_0^2} \approx \frac{\|\mathbf{z}\|^2}{\nu_0^2} \gg \frac{\|\mathbf{z}\|^2}{\nu_i^2} \approx \frac{\|\mathbf{z}-\boldsymbol{\mu}_i\|^2}{\nu_i^2}$. We call 180 such mode $P^{(0)}$ the *universal mode*. Therefore, if z is initialized far from all modes, it can only 181 converge to the universal mode because the gradient information of other modes is masked. Once 182 z enters the universal mode $P^{(0)}$, if the step size δ_t of Langevin dynamics is small (i.e., $\delta_t \leq \nu_0^2$), 183 it would take exponential steps to escape the local mode $P^{(0)}$; while if the step size is large (i.e., $\delta_t > \nu_0^2$), the state **z** would again be far from all modes and thus the universal mode $P^{(0)}$ dominates all other modes. Hence, **z** can only visit the universal mode unless the stochastic noise ϵ_t miraculously 184 185 186 leads it to the region of another mode. In addition, it can be verified that $\log\left(\frac{\nu_i^2}{\nu_0^2}\right) - \frac{\nu_i^2}{2\nu_0^2} + \frac{\nu_0^2}{2\nu_i^2}$ is a positive constant for $\nu_i < \nu_0$, thus the second requirement of Assumption 1 essentially represents 187 188 $\|\boldsymbol{\mu}_i - \boldsymbol{\mu}_0\|^2 \leq \mathcal{O}(d)$. We formalize the intuition in Theorem 1 and defer the proof to Appendix A.1. 189

Theorem 1. Consider a data distribution P satisfying Assumption 1. We follow Langevin dynamics for $T = \exp(\mathcal{O}(d))$ steps. Suppose the sample is initialized in $P^{(0)}$, then with probability at least $1 - T \cdot \exp(-\Omega(d))$, we have $\|\mathbf{x}_t - \boldsymbol{\mu}_i\|^2 > \frac{\nu_0^2 + \nu_{\max}^2}{2} d$ for all $t \in \{0\} \cup [T]$ and $i \in [k]$.

We note that $\|\mathbf{x}_t - \boldsymbol{\mu}_i\|^2 > \frac{\nu_0^2 + \nu_{\max}^2}{2} d$ is a strong notion of mode-seekingness, since the probability density of mode $P^{(i)} = \mathcal{N}(\boldsymbol{\mu}_i, \nu_i^2 \boldsymbol{I}_d)$ concentrates around the ℓ_2 -norm ball $\{\mathbf{z} : \|\mathbf{z} - \boldsymbol{\mu}_i\|^2 \le \nu_i^2 d\}$. This notion can also easily be translated into a lower bound in terms of other distance measures such as total variation distance and Wasserstein 2-distance. Moreover, in Theorem 2 we extend the result to annealed Langevin dynamics with bounded noise level, and the proof is deferred to Appendix A.2.

Theorem 2. Consider a data distribution P satisfying Assumption 1. We follow annealed Langevin dynamics for $T = \exp(\mathcal{O}(d))$ steps with noise levels $c_{\sigma} \ge \sigma_0 \ge \cdots \ge \sigma_T \ge 0$ for constant $c_{\sigma} > 0$. In addition, assume for all $i \in [k]$, $\|\boldsymbol{\mu}_i - \boldsymbol{\mu}_0\|^2 \le \frac{\nu_0^2 - \nu_i^2}{2} \left(\log \left(\frac{\nu_i^2 + c_{\sigma}^2}{\nu_0^2 + c_{\sigma}^2} \right) - \frac{\nu_i^2 + c_{\sigma}^2}{2\nu_i^2 + c_{\sigma}^2} + \frac{\nu_0^2 + c_{\sigma}^2}{2\nu_i^2 + c_{\sigma}^2} \right) d.$ Suppose that the sample is initialized in $P_{\sigma_0}^{(0)}$, then with probability at least $1 - T \cdot \exp(-\Omega(d))$, we have $\|\mathbf{x}_t - \boldsymbol{\mu}_i\|^2 > \frac{\nu_0^2 + \nu_{\max}^2 + 2\sigma_i^2}{2} d$ for all $t \in \{0\} \cup [T]$ and $i \in [k]$.

203 4.2 Langevin Dynamics in Sub-Gaussian Mixtures

We further generalize our results to sub-Gaussian mixtures. We impose the following assumptions on the mixture. It is worth noting that these assumptions automatically hold for Gaussian mixtures.

Assumption 2. Consider a data distribution $P := \sum_{i=0}^{k} w_i P^{(i)}$ as a mixture of sub-Gaussian distributions, where $1 \le k = o(d)$ and $w_i > 0$ is a positive constant such that $\sum_{i=0}^{k} w_i = 1$. Suppose that $P^{(0)} = \mathcal{N}(\boldsymbol{\mu}_0, \nu_0^2 \boldsymbol{I}_d)$ is Gaussian and for all $i \in [k]$, $P^{(i)}$ satisfies

209 *i.* $P^{(i)}$ is a sub-Gaussian distribution of mean μ_i with parameter ν_i^2 ,

210 *ii.* $P^{(i)}$ is differentiable and $\nabla P^{(i)}(\boldsymbol{\mu}_i) = \mathbf{0}_d$,

iii. the score function of
$$P^{(i)}$$
 is L_i -Lipschitz such that $L_i \leq \frac{c_L}{\mu^2}$ for some constant $c_L > 0$,

212 *iv.*
$$\nu_0^2 > \max\left\{1, \frac{4(c_L^2 + c_\nu c_L)}{c_\nu (1 - c_\nu)}\right\} \frac{\nu_{\max}^2}{1 - c_\nu}$$
 for constant $c_\nu \in (0, 1)$, where $\nu_{\max} := \max_{i \in [k]} \nu_i$,

213 V.
$$\|\boldsymbol{\mu}_i - \boldsymbol{\mu}_0\|^2 \le \frac{(1-c_{\nu})\nu_0^2 - \nu_i^2}{2(1-c_{\nu})} \left(\log \frac{c_{\nu}\nu_i^2}{(c_L^2 + c_{\nu}c_L)\nu_0^2} - \frac{\nu_i^2}{2(1-c_{\nu})\nu_0^2} + \frac{(1-c_{\nu})\nu_0^2}{2\nu_i^2} \right) dt$$

We validate the feasibility of Assumption 2.v. in Lemma 9 in the Appendix. With Assumption 2, we
show the mode-seeking tendency of Langevin dynamics under sub-Gaussian distributions in Theorem
and defer the proof to Appendix A.3.

Algorithm 1 Chained Langevin Dynamics (Chained-LD)

Require: Patch size Q, dimension d, conditional score function estimator s_{θ} , number of iterations T, noise levels $\{\sigma_t\}_{t \in [TQ/d]}$, step size $\{\delta_t\}_{t \in [TQ/d]}$.

- 1: Initialize \mathbf{x}_0 , and divide \mathbf{x}_0 into d/Q patches $\mathbf{x}_0^{(1)}, \cdots \mathbf{x}_0^{(d/Q)}$ of equal size Q
- 2: for $q \leftarrow 1$ to d/Q do
- for $t \leftarrow 1$ to TQ/d do 3:
- $\mathbf{x}_{t}^{(q)} \leftarrow \mathbf{x}_{t-1}^{(q)} + \frac{\delta_{t}}{2} \mathbf{s}_{\boldsymbol{\theta}} \left(\mathbf{x}_{t}^{(q)} \mid \sigma_{t}, \mathbf{x}_{t}^{(1)}, \cdots, \mathbf{x}_{t}^{(q-1)} \right) + \sqrt{\delta_{t}} \boldsymbol{\epsilon}_{t}, \text{ where } \boldsymbol{\epsilon}_{t} \sim \mathcal{N}(\mathbf{0}_{Q}, \boldsymbol{I}_{Q})$ 4:
- 5:
- end for $\mathbf{x}_{0}^{(q)} \leftarrow \mathbf{x}_{TQ/d}^{(q)}$ 6:

7: end for

8: return $\mathbf{x}_{TQ/d}$

217 **Theorem 3.** Consider a data distribution P satisfying Assumption 2. We follow Langevin dynamics for $T = \exp(\mathcal{O}(d))$ steps. Suppose the sample is initialized in $P^{(0)}$, then with probability at least 218

 $1 - T \cdot \exp(-\mathcal{O}(d))$, we have $\|\mathbf{x}_t - \boldsymbol{\mu}_i\|^2 > \left(\frac{\nu_0^2}{2} + \frac{\nu_{\max}^2}{2(1 - c_{\nu})}\right) d$ for all $t \in \{0\} \cup [T]$ and $i \in [k]$. 219

Finally, we slightly modify Assumption 2 and extend our results to annealed Langevin dynamics 220 under sub-Gaussian mixtures in Theorem 4. The details of Assumption 3 and the proof of Theorem 4 221 are deferred to Appendix A.4. 222

Theorem 4. Consider a data distribution P satisfying Assumption 3. We follow annealed Langevin 223 dynamics for $T = \exp(\mathcal{O}(d))$ steps with noise levels $c_{\sigma} \geq \sigma_0 \geq \cdots \geq \sigma_T \geq 0$. Suppose 224 the sample is initialized in $P_{\sigma_0}^{(0)}$, then with probability at least $1 - T \cdot \exp(-\mathcal{O}(d))$, we have $\|\mathbf{x}_t - \boldsymbol{\mu}_i\|^2 > \left(\frac{\nu_0^2 + \sigma_t^2}{2} + \frac{\nu_{\max}^2 + \sigma_t^2}{2(1 - c_\nu)}\right) d$ for all $t \in \{0\} \cup [T]$ and $i \in [k]$. 225 226

Chained Langevin Dynamics 5 227

To reduce the mode-seeking tendencies of vanilla and annealed Langevin dynamics, we propose 228 Chained Langevin Dynamics (Chained-LD) in Algorithm 1. While vanilla and annealed Langevin 229 dynamics apply gradient updates to all coordinates of the sample in every step, we decompose the 230 sample into patches of constant size and generate each patch sequentially to alleviate the exponen-231 tial dependency on the dimensionality. More precisely, we divide a sample x into d/Q patches 232 $\mathbf{x}^{(1)}, \cdots \mathbf{x}^{(d/Q)}$ of some constant size Q, and apply annealed Langevin dynamics to sample each patch $\mathbf{x}^{(q)}$ (for $q \in [d/Q]$) from the conditional distribution $P(\mathbf{x}^{(q)} | \mathbf{x}^{(1)}, \cdots \mathbf{x}^{(q-1)})$. 233 234

An ideal conditional score function estimator s_{θ} could jointly estimate the scores of all perturbed 235 conditional patch distribution, i.e., $\forall \sigma \in \{\sigma_t\}_{t \in [TQ/d]}, q \in [d/Q],$ 236

$$\mathbf{s}_{\boldsymbol{\theta}}\left(\mathbf{x}^{(q)} \mid \sigma, \mathbf{x}^{(1)}, \cdots, \mathbf{x}^{(q-1)}\right) \approx \nabla_{\mathbf{x}^{(q)}} \log P_{\sigma}(\mathbf{x}^{(q)} \mid \mathbf{x}^{(1)}, \cdots, \mathbf{x}^{(q-1)}).$$

Following from Song and Ermon (2019), we use the denoising score matching to train the estimator. 237 For a given σ , the denoising score matching objective is 238

$$\ell(\boldsymbol{\theta};\sigma) := \frac{1}{2} \mathbb{E}_{\mathbf{x} \sim P} \mathbb{E}_{\tilde{\mathbf{x}} \sim \mathcal{N}(\mathbf{x},\sigma^{2}\boldsymbol{I}_{d})} \sum_{q \in [d/Q]} \left[\left\| \mathbf{s}_{\boldsymbol{\theta}} \left(\mathbf{x}^{(q)} \mid \sigma, \mathbf{x}^{(1)}, \cdots, \mathbf{x}^{(q-1)} \right) - \frac{\tilde{\mathbf{x}}^{(q)} - \mathbf{x}^{(q)}}{\sigma^{2}} \right\|^{2} \right].$$

Then, combining the objectives gives the following loss 239

$$\mathcal{L}\left(\boldsymbol{\theta}; \{\sigma_t\}_{t \in [TQ/d]}\right) := \frac{d}{TQ} \sum_{t \in [TQ/d]} \sigma_t^2 \ell(\boldsymbol{\theta}; \sigma_t).$$

As shown in Vincent (2011), an estimator s_{θ} with enough capacity minimizes the loss \mathcal{L} if and only if 240 s_{θ} outputs the scores of all perturbed conditional patch distribution almost surely. Ideally, if a sampler 241

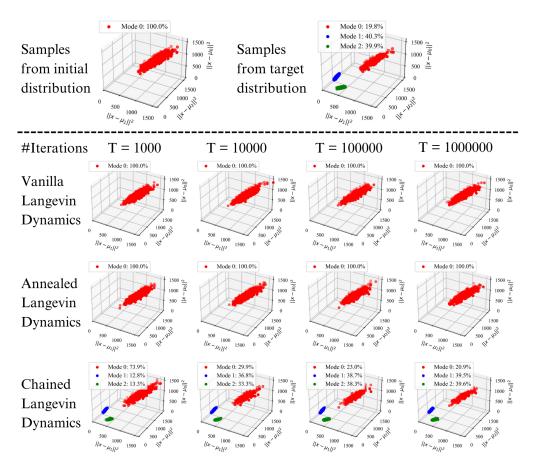


Figure 1: Samples from a mixture of three Gaussian modes generated by vanilla, annealed, and chained Langevin dynamics. Three axes are ℓ_2 distance from samples to the mean of the three modes. The samples are initialized in mode 0.

perfectly generates every patch, combining all patches gives a sample from the original distribution since $P(\mathbf{x}) = \prod_{q \in [d/Q]} P(\mathbf{x}^{(q)} | \mathbf{x}^{(1)}, \cdots \mathbf{x}^{(q-1)})$. In Theorem 5 we give a linear reduction from producing samples of dimension *d* using Chained-LD to learning the distribution of a *Q*-dimensional variable for constant *Q*. The proof of Theorem 5 is deferred to Appendix A.5.

Theorem 5. Consider a sampler algorithm taking the first q - 1 patches $\mathbf{x}^{(1)}, \dots, \mathbf{x}^{(q-1)}$ as input and outputing a sample of the next patch $\mathbf{x}^{(q)}$ with probability $\hat{P}(\mathbf{x}^{(q)} | \mathbf{x}^{(1)}, \dots, \mathbf{x}^{(q-1)})$ for all $q \in [d/Q]$. Suppose that for every $q \in [d/Q]$ and any given previous patches $\mathbf{x}^{(1)}, \dots, \mathbf{x}^{(q-1)}$, the sampler algorithm can achieve

$$TV\left(\hat{P}\left(\mathbf{x}^{(q)} \mid \mathbf{x}^{(1)}, \cdots, \mathbf{x}^{(q-1)}\right), P\left(\mathbf{x}^{(q)} \mid \mathbf{x}^{(1)}, \cdots, \mathbf{x}^{(q-1)}\right)\right) \leq \varepsilon \cdot \frac{Q}{d}$$

in $\tau(\varepsilon, d)$ iterations for some $\varepsilon > 0$. Then, equipped with the sampler algorithm, the Chained-LD algorithm in $\frac{d}{O} \cdot \tau(\varepsilon, d)$ iterations can achieve

$$TV\left(\hat{P}(\mathbf{x}), P(\mathbf{x})\right) \leq \varepsilon.$$

252 6 Numerical Results

In this section, we empirically evaluated the mode-seeking tendencies of vanilla, annealed, and chained Langevin dynamics. We performed numerical experiments on synthetic Gaussian mixture models and real image datasets including MNIST (LeCun, 1998) and Fashion-MNIST (Xiao et al., 2017). Details on the experiment setup are deferred to Appendix B.

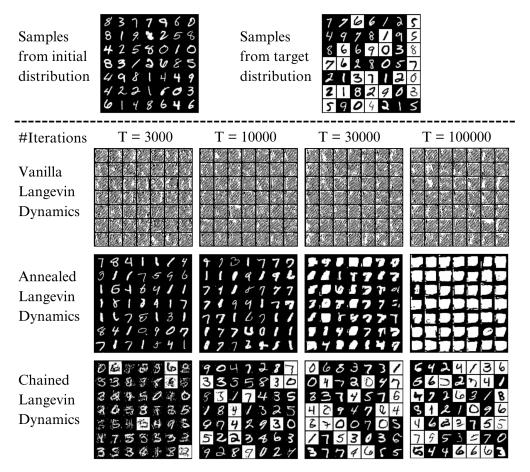


Figure 2: Samples from a mixture distribution of the original and flipped images from the MNIST dataset generated by vanilla, annealed, and chained Langevin dynamics. The samples are initialized as original images from MNIST.

Synthetic Gaussian mixture model: We define the data distribution P as a mixture of three Gaussian components in dimension d = 100, where mode 0 defined as $P^{(0)} = \mathcal{N}(\mathbf{0}_d, 3\mathbf{I}_d)$ is the universal mode with the largest variance, and mode 1 and mode 2 are respectively defined as $P^{(1)} = \mathcal{N}(\mathbf{1}_d, \mathbf{I}_d)$ and $P^{(2)} = \mathcal{N}(-\mathbf{1}_d, \mathbf{I}_d)$. The frequencies of the three modes are 0.2, 0.4 and 0.4, i.e.,

$$P = 0.2P^{(0)} + 0.4P^{(1)} + 0.4P^{(2)} = 0.2\mathcal{N}(\mathbf{0}_d, 3\mathbf{I}_d) + 0.4\mathcal{N}(\mathbf{1}_d, \mathbf{I}_d) + 0.4\mathcal{N}(-\mathbf{1}_d, \mathbf{I}_d).$$

As shown in Figure 1, vanilla and annealed Langevin dynamics cannot find mode 1 or 2 within 10⁶ iterations if the sample is initialized in mode 0, while chained Langevin dynamics can find the other two modes in 1000 steps and correctly recover their frequencies as gradually increasing the number of iterations. In Appendix B.1 we present additional experiments on samples initialized in mode 1 or 2, which also verify the mode-seeking tendencies of vanilla and annealed Langevin dynamics.

Image datasets: We construct the distribution as a mixture of two modes by using the original images 266 from MNIST/Fashion-MNIST training dataset (black background and white digits/objects) as the 267 first mode and constructing the second mode by i.i.d. randomly flipping an image (white background 268 and black digits/objects) with probability 0.5. Regarding the neural network architecture of the score 269 function estimator, for vanilla and annealed Langevin dynamics we use U-Net (Ronneberger et al., 270 2015) following from Song and Ermon (2019). For chained Langevin dynamics, we proposed to use 271 Recurrent Neural Network (RNN) architectures. We note that for a sequence of inputs, the output of 272 RNN from the previous step is fed as input to the current step. Therefore, in the scenario of chained 273 Langevin dynamics, the hidden state of RNN contains information about the previous patches and 274 allows the network to estimate the conditional score function $\nabla_{\mathbf{x}^{(q)}} \log P(\mathbf{x}^{(q)} \mid \mathbf{x}^{(1)}, \cdots, \mathbf{x}^{(q-1)})$. 275 More implementation details are deferred to Appendix B.2. 276

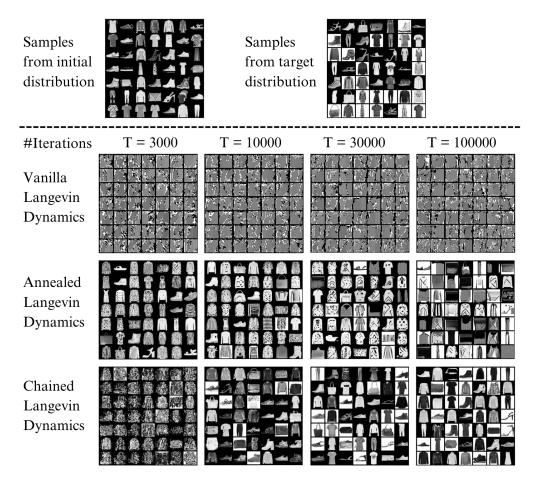


Figure 3: Samples from a mixture distribution of the original and flipped images from the Fashion-MNIST dataset generated by vanilla, annealed, and chained Langevin dynamics. The samples are initialized as original images from Fashion-MNIST.

The numerical results on image datasets are shown in Figures 2 and 3. Vanilla Langevin dynamics fails to generate reasonable samples, as also observed in Song and Ermon (2019). When the sample is initialized as original images from the datasets, annealed Langevin dynamics tends to generate samples from the same mode, while chained Langevin dynamics can generate samples from both modes. Additional experiments are deferred to Appendix B.2.

282 7 Conclusion

In this work, we theoretically and numerically studied the mode-seeking properties of vanilla and annealed Langevin dynamics sampling methods under a multi-modal distribution. We characterized Gaussian and sub-Gaussian mixture models under which Langevin dynamics are unlikely to find all the components within a sub-exponential number of iterations. To reduce the mode-seeking tendency of vanilla Langevin dynamics, we proposed Chained Langevin Dynamics (Chained-LD) and analyzed its convergence behavior. Studying the connections between Chained-LD and denoising diffusion models will be an interesting topic for future exploration.

290 Limitations

Our RNN-based implementation of Chained-LD is currently limited to image data generation tasks. An interesting future direction is to extend the application of Chained-LD to other domains such as audio and text data. Another future direction could be to study the convergence of Chained-LD under an imperfect score estimation which we did not address in our analysis.

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413 A Theoretical Analysis on the Mode-Seeking Tendency of Langevin Dynamics

We begin by introducing some well-established lemmas used in our proof. We first provide the proof of Proposition 1 for completeness:

⁴¹⁶ *Proof of Proposition 1.* By the definition in equation 2, we have

$$p_{\sigma}(\mathbf{z}) = \int p(\mathbf{t}) \mathcal{N}(\mathbf{z} \mid \mathbf{t}, \sigma^2 \mathbf{I}_d) \, \mathrm{d}\mathbf{t} = \int p(\mathbf{t}) \mathcal{N}(\mathbf{z} - \mathbf{t} \mid \mathbf{0}_d, \sigma^2 \mathbf{I}_d) \, \mathrm{d}\mathbf{t}.$$

For random variables $\mathbf{t} \sim p$ and $\mathbf{y} \sim \mathcal{N}(\mathbf{0}_d, \mathbf{I}_d)$, their sum $\mathbf{z} = \mathbf{t} + \mathbf{y} \sim p_\sigma$ follows the perturbed distribution with noise level σ . Therefore,

$$\mathbb{E}_{\mathbf{z} \sim p_{\sigma}}[\mathbf{z}] = \mathbb{E}_{(\mathbf{t}+\mathbf{y}) \sim p_{\sigma}}[\mathbf{t}+\mathbf{y}] = \mathbb{E}_{\mathbf{t} \sim p}[\mathbf{t}] + \mathbb{E}_{\mathbf{y} \sim \mathcal{N}(\mathbf{0}_{d}, \mathbf{I}_{d})}[\mathbf{y}] = \mathbb{E}_{\mathbf{t} \sim p}[\mathbf{t}].$$

If $\mathbf{t} \sim p = \mathcal{N}(\boldsymbol{\mu}, \boldsymbol{\Sigma})$ follows a Gaussian distribution, we have $\mathbf{z} = \mathbf{t} + \mathbf{y} \sim p_{\sigma} = \mathcal{N}(\boldsymbol{\mu}, \boldsymbol{\Sigma} + \sigma^2 \boldsymbol{I}_d)$. If p is a sub-Gaussian distribution with parameter ν^2 , we have $\mathbf{z} = \mathbf{t} + \mathbf{y} \sim p_{\sigma}$ is a sub-Gaussian

420 If p is a sub-statistical of the parameter ν , we have $2 = 0 + g = p_{\sigma}$ is a sub-421 distribution with parameter ($\nu^2 + \sigma^2$). Hence we obtain Proposition 1.

⁴²² We use the following lemma on the tail bound for multivariate Gaussian random variables.

- Lemma 1 (Lemma 1, Laurent and Massart (2000)). Suppose that a random variable $\mathbf{z} \sim \mathcal{N}(\mathbf{0}_d, \mathbf{I}_d)$.
- 424 Then for any $\lambda > 0$,

$$\mathbb{P}\left(\left\|\mathbf{z}\right\|^{2} \ge d + 2\sqrt{d\lambda} + 2\lambda\right) \le \exp(-\lambda),$$
$$\mathbb{P}\left(\left\|\mathbf{z}\right\|^{2} \le d - 2\sqrt{d\lambda}\right) \le \exp(-\lambda).$$

- We also use a tail bound for one-dimensional Gaussian random variables and provide the proof here for completeness.
- **Lemma 2.** Suppose a random variable $Z \sim \mathcal{N}(0, 1)$. Then for any t > 0,

$$\mathbb{P}(Z \ge t) = \mathbb{P}(Z \le -t) \le \frac{\exp(-t^2/2)}{\sqrt{2\pi}t}.$$

428 Proof of Lemma 2. Since $\frac{z}{t} \ge 1$ for all $z \in [t, \infty)$, we have

$$\mathbb{P}(Z \ge t) = \frac{1}{\sqrt{2\pi}} \int_t^\infty \exp\left(-\frac{z^2}{2}\right) \, \mathrm{d}z \le \frac{1}{\sqrt{2\pi}} \int_t^\infty \frac{z}{t} \exp\left(-\frac{z^2}{2}\right) \, \mathrm{d}z = \frac{\exp(-t^2/2)}{\sqrt{2\pi}t}.$$

Since the Gaussian distribution is symmetric, we have $\mathbb{P}(Z \ge t) = \mathbb{P}(Z \le -t)$. Hence we obtain the desired bound.

431 A.1 Proof of Theorem 1: Langevin Dynamics under Gaussian Mixtures

Without loss of generality, we assume that $\mu_0 = \mathbf{0}_d$ for simplicity. Let r and n respectively denote the rank and nullity of the vector space $\{\boldsymbol{\mu}_i\}_{i\in[k]}$, then we have r + n = d and $0 \le r \le k = o(d)$. Denote $\mathbf{R} \in \mathbb{R}^{d \times r}$ an orthonormal basis of the vector space $\{\boldsymbol{\mu}_i\}_{i\in[k]}$, and denote $\mathbf{N} \in \mathbb{R}^{d \times n}$ an orthonormal basis of the null space of $\{\boldsymbol{\mu}_i\}_{i\in[k]}$. Now consider decomposing the sample \mathbf{x}_t by

$$\mathbf{r}_t := \mathbf{R}^T \mathbf{x}_t$$
, and $\mathbf{n}_t := \mathbf{N}^T \mathbf{x}_t$

436 where $\mathbf{r}_t \in \mathbb{R}^r$, $\mathbf{n}_t \in \mathbb{R}^n$. Then we have

$$\mathbf{x}_t = \mathbf{R}\mathbf{r}_t + \mathbf{N}\mathbf{n}_t.$$

437 Similarly, we decompose the noise ϵ_t into

$$\boldsymbol{\epsilon}_t^{(\mathbf{r})} := \mathbf{R}^T \boldsymbol{\epsilon}_t, \text{ and } \boldsymbol{\epsilon}_t^{(\mathbf{n})} := \mathbf{N}^T \boldsymbol{\epsilon}_t,$$

438 where $\boldsymbol{\epsilon}_t^{(\mathbf{r})} \in \mathbb{R}^r, \, \boldsymbol{\epsilon}_t^{(\mathbf{n})} \in \mathbb{R}^n.$ Then we have

$$\boldsymbol{\epsilon}_t = \mathbf{R}\boldsymbol{\epsilon}_t^{(\mathbf{r})} + \mathbf{N}\boldsymbol{\epsilon}_t^{(\mathbf{n})}.$$

Since a linear combination of a Gaussian random variable still follows Gaussian distribution, by $\epsilon_t \sim \mathcal{N}(\mathbf{0}_d, \mathbf{I}_d), \mathbf{R}^T \mathbf{R} = \mathbf{I}_r$, and $\mathbf{N}^T \mathbf{N} = \mathbf{I}_n$ we obtain

$$\boldsymbol{\epsilon}_t^{(\mathbf{r})} \sim \mathcal{N}(\mathbf{0}_r, \boldsymbol{I}_r), ext{ and } \boldsymbol{\epsilon}_t^{(\mathbf{n})} \sim \mathcal{N}(\mathbf{0}_n, \boldsymbol{I}_n).$$

By the definition of Langevin dynamics in equation 1, the two components of \mathbf{x}_t follow from the update rule:

$$\mathbf{n}_{t} = \mathbf{n}_{t-1} + \frac{\delta_{t}}{2} \mathbf{N}^{T} \nabla_{\mathbf{x}} \log P(\mathbf{x}_{t-1}) + \sqrt{\delta_{t}} \boldsymbol{\epsilon}_{t}^{(\mathbf{n})},$$
(5)
$$\mathbf{r}_{t} = \mathbf{r}_{t-1} + \frac{\delta_{t}}{2} \mathbf{R}^{T} \nabla_{\mathbf{x}} \log P(\mathbf{x}_{t-1}) + \sqrt{\delta_{t}} \boldsymbol{\epsilon}_{t}^{(\mathbf{r})}.$$

443 It is worth noting that since $\mathbf{N}^T \boldsymbol{\mu}_i = \mathbf{0}_n$. To show $\|\mathbf{x}_t - \boldsymbol{\mu}_i\|^2 > \frac{\nu_0^2 + \nu_{\max}^2}{2}d$, it suffices to prove

$$\|\mathbf{n}_t\|^2 > \frac{\nu_0^2 + \nu_{\max}^2}{2} d.$$

We start by proving that the initialization of the state x_0 has a large norm on the null space with high probability in the following proposition.

Proposition 2. Suppose that a sample \mathbf{x}_0 is initialized in the distribution $P^{(0)}$, i.e., $\mathbf{x}_0 \sim P^{(0)}$, then for any constant $\nu_{\max} < \nu_0$, with probability at least $1 - \exp(-\Omega(d))$, we have $\|\mathbf{n}_0\|^2 \ge \frac{3\nu_0^2 + \nu_{\max}^2}{4}d$.

448 Proof of Proposition 2. Since $\mathbf{x}_0 \sim P^{(0)} = \mathcal{N}(\mathbf{0}_d, \nu_0^2 \mathbf{I}_d)$ and $\mathbf{N}^T \mathbf{N} = \mathbf{I}_n$, we know $\mathbf{n}_0 = \mathbf{N}^T \mathbf{x}_0 \sim$ 449 $\mathcal{N}(\mathbf{0}_n, \nu_0^2 \mathbf{I}_n)$. Therefore, by Lemma 1 we can bound

$$\begin{split} \mathbb{P}\left(\left\|\mathbf{n}_{0}\right\|^{2} \leq \frac{3\nu_{0}^{2} + \nu_{\max}^{2}}{4}d\right) &= \mathbb{P}\left(\frac{\left\|\mathbf{n}_{0}\right\|^{2}}{\nu_{0}^{2}} \leq d - 2\sqrt{d \cdot \left(\frac{\nu_{0}^{2} - \nu_{\max}^{2}}{8\nu_{0}^{2}}\right)^{2}d}\right) \\ &\leq \mathbb{P}\left(\frac{\left\|\mathbf{n}_{0}\right\|^{2}}{\nu_{0}^{2}} \leq n - 2\sqrt{n\left(\frac{\nu_{0}^{2} - \nu_{\max}^{2}}{8\nu_{0}^{2}}\right)^{2}\frac{d}{2}}\right) \\ &\leq \exp\left(-\left(\frac{\nu_{0}^{2} - \nu_{\max}^{2}}{8\nu_{0}^{2}}\right)^{2}\frac{d}{2}\right), \end{split}$$

where the second last step follows from the assumption d - n = r = o(d). Hence we complete the proof of Proposition 2.

Then, with the assumption that the initialization satisfies $\|\mathbf{n}_0\|^2 \ge \frac{3\nu_0^2 + \nu_{\max}^2}{4}d$, the following proposition shows that $\|\mathbf{n}_t\|$ remains large with high probability.

Proposition 3. Consider a data distribution P satisfies the constraints specified in Theorem 1. We follow the Langevin dynamics for $T = \exp(\mathcal{O}(d))$ steps. Suppose that the initial sample satisfies $\|\mathbf{n}_0\|^2 \ge \frac{3\nu_0^2 + \nu_{\max}^2}{4}d$, then with probability at least $1 - T \cdot \exp(-\Omega(d))$, we have that $\|\mathbf{n}_t\|^2 > \frac{\nu_0^2 + \nu_{\max}^2}{2}d$ for all $t \in \{0\} \cup [T]$.

Proof of Proposition 3. To establish a lower bound on $||\mathbf{n}_t||$, we consider different cases of the step size δ_t . Intuitively, when δ_t is large enough, \mathbf{n}_t will be too noisy due to the introduction of random noise $\sqrt{\delta_t} \boldsymbol{\epsilon}_t^{(\mathbf{n})}$ in equation 5. While for small δ_t , the update of \mathbf{n}_t is bounded and thus we can iteratively analyze \mathbf{n}_t . We first handle the case of large δ_t in the following lemma.

462 **Lemma 3.** If $\delta_t > \nu_0^2$, with probability at least $1 - \exp(-\Omega(d))$, for \mathbf{n}_t satisfying equation 5, we 463 have $\|\mathbf{n}_t\|^2 \ge \frac{3\nu_0^2 + \nu_{\max}^2}{4} d$ regardless of the previous state \mathbf{x}_{t-1} .

464 Proof of Lemma 3. Denote $\mathbf{v} := \mathbf{n}_{t-1} + \frac{\delta_t}{2} \mathbf{N}^T \nabla_{\mathbf{x}} \log P(\mathbf{x}_{t-1})$ for simplicity. Note that \mathbf{v} is fixed 465 for any given \mathbf{x}_{t-1} . We decompose $\boldsymbol{\epsilon}_t^{(\mathbf{n})}$ into a vector aligning with \mathbf{v} and another vector orthogonal to **v**. Consider an orthonormal matrix $\mathbf{M} \in \mathbb{R}^{n \times (n-1)}$ such that $\mathbf{M}^T \mathbf{v} = \mathbf{0}_{n-1}$ and $\mathbf{M}^T \mathbf{M} = \mathbf{I}_{n-1}$. By denoting $\mathbf{u} := \boldsymbol{\epsilon}_t^{(\mathbf{n})} - \mathbf{M}\mathbf{M}^T \boldsymbol{\epsilon}_t^{(\mathbf{n})}$ we have $\mathbf{M}^T \mathbf{u} = \mathbf{0}_{n-1}$, thus we obtain

$$\begin{split} \|\mathbf{n}_{t}\|^{2} &= \left\|\mathbf{v} + \sqrt{\delta_{t}} \boldsymbol{\epsilon}_{t}^{(\mathbf{n})}\right\|^{2} \\ &= \left\|\mathbf{v} + \sqrt{\delta_{t}} \mathbf{u} + \sqrt{\delta_{t}} \mathbf{M} \mathbf{M}^{T} \boldsymbol{\epsilon}_{t}^{(\mathbf{n})}\right\|^{2} \\ &= \left\|\mathbf{v} + \sqrt{\delta_{t}} \mathbf{u}\right\|^{2} + \left\|\sqrt{\delta_{t}} \mathbf{M} \mathbf{M}^{T} \boldsymbol{\epsilon}_{t}^{(\mathbf{n})}\right\|^{2} \\ &\geq \left\|\sqrt{\delta_{t}} \mathbf{M} \mathbf{M}^{T} \boldsymbol{\epsilon}_{t}^{(\mathbf{n})}\right\|^{2} \\ &\geq \nu_{0}^{2} \left\|\mathbf{M}^{T} \boldsymbol{\epsilon}_{t}^{(\mathbf{n})}\right\|^{2}. \end{split}$$

Since $\boldsymbol{\epsilon}_t^{(\mathbf{n})} \sim \mathcal{N}(\mathbf{0}_n, \boldsymbol{I}_n)$ and $\mathbf{M}^T \mathbf{M} = \boldsymbol{I}_{n-1}$, we obtain $\mathbf{M}^T \boldsymbol{\epsilon}_t^{(\mathbf{n})} \sim \mathcal{N}(\mathbf{0}_{n-1}, \boldsymbol{I}_{n-1})$. Therefore, by Lemma 1 we can bound

$$\begin{split} \mathbb{P}\left(\left\|\mathbf{n}_{t}\right\|^{2} \leq \frac{3\nu_{0}^{2} + \nu_{\max}^{2}}{4}d\right) \leq \mathbb{P}\left(\left\|\mathbf{M}^{T}\boldsymbol{\epsilon}_{t}^{(\mathbf{n})}\right\|^{2} \leq \frac{3\nu_{0}^{2} + \nu_{\max}^{2}}{4\nu_{0}^{2}}d\right) \\ &= \mathbb{P}\left(\left\|\mathbf{M}^{T}\boldsymbol{\epsilon}_{t}^{(\mathbf{n})}\right\|^{2} \leq d - 2\sqrt{d \cdot \left(\frac{\nu_{0}^{2} - \nu_{\max}^{2}}{8\nu_{0}^{2}}\right)^{2}d}\right) \\ &\leq \mathbb{P}\left(\left\|\mathbf{M}^{T}\boldsymbol{\epsilon}_{t}^{(\mathbf{n})}\right\|^{2} \leq (n - 1) - 2\sqrt{(n - 1)\left(\frac{\nu_{0}^{2} - \nu_{\max}^{2}}{8\nu_{0}^{2}}\right)^{2}\frac{d}{2}}\right) \\ &\leq \exp\left(-\left(\frac{\nu_{0}^{2} - \nu_{\max}^{2}}{8\nu_{0}^{2}}\right)^{2}\frac{d}{2}\right), \end{split}$$

where the second last step follows from the assumption d - n = r = o(d). Hence we complete the proof of Lemma 3.

We then consider the case when $\delta_t \leq \nu_0^2$. Let $\mathbf{r} := \mathbf{R}^T \mathbf{x}$ and $\mathbf{n} := \mathbf{N}^T \mathbf{x}$, then $\mathbf{x} = \mathbf{R}\mathbf{r} + \mathbf{N}\mathbf{n}$. We first show that when $\|\mathbf{n}\|^2 \geq \frac{\nu_0^2 + \nu_{\max}^2}{2} d$, $P^{(i)}(\mathbf{x})$ is exponentially smaller than $P^{(0)}(\mathbf{x})$ for all $i \in [k]$ in the following lemma.

475 **Lemma 4.** Given that $\|\mathbf{n}\|^2 \ge \frac{\nu_0^2 + \nu_{\max}^2}{2} d$ and $\|\boldsymbol{\mu}_i\|^2 \le \frac{\nu_0^2 - \nu_i^2}{2} \left(\log\left(\frac{\nu_i^2}{\nu_0^2}\right) - \frac{\nu_i^2}{2\nu_0^2} + \frac{\nu_0^2}{2\nu_i^2} \right) d$ for all 476 $i \in [k]$, we have $\frac{P^{(i)}(\mathbf{x})}{P^{(0)}(\mathbf{x})} \le \exp(-\Omega(d))$ for all $i \in [k]$.

477 Proof of Lemma 4. For all $i \in [k]$, define $\rho_i(\mathbf{x}) := \frac{P^{(i)}(\mathbf{x})}{P^{(0)}(\mathbf{x})}$, then

$$\begin{split} \rho_i(\mathbf{x}) &= \frac{P^{(i)}(\mathbf{x})}{P^{(0)}(\mathbf{x})} = \frac{(2\pi\nu_i^2)^{-d/2}\exp\left(-\frac{1}{2\nu_i^2}\|\mathbf{x}-\boldsymbol{\mu}_i\|^2\right)}{(2\pi\nu_0^2)^{-d/2}\exp\left(-\frac{1}{2\nu_0^2}\|\mathbf{x}\|^2\right)} \\ &= \left(\frac{\nu_0^2}{\nu_i^2}\right)^{d/2}\exp\left(\frac{1}{2\nu_0^2}\|\mathbf{x}\|^2 - \frac{1}{2\nu_i^2}\|\mathbf{x}-\boldsymbol{\mu}_i\|^2\right) \\ &= \left(\frac{\nu_0^2}{\nu_i^2}\right)^{d/2}\exp\left(\left(\frac{1}{2\nu_0^2} - \frac{1}{2\nu_i^2}\right)\|\mathbf{Nn}\|^2 + \left(\frac{\|\mathbf{Rr}\|^2}{2\nu_0^2} - \frac{\|\mathbf{Rr}-\boldsymbol{\mu}_i\|^2}{2\nu_i^2}\right)\right) \\ &= \left(\frac{\nu_0^2}{\nu_i^2}\right)^{d/2}\exp\left(\left(\frac{1}{2\nu_0^2} - \frac{1}{2\nu_i^2}\right)\|\mathbf{n}\|^2 + \left(\frac{\|\mathbf{r}\|^2}{2\nu_0^2} - \frac{\|\mathbf{r}-\mathbf{R}^T\boldsymbol{\mu}_i\|^2}{2\nu_i^2}\right)\right), \end{split}$$

where the last step follows from the definition that $\mathbf{R} \in \mathbb{R}^{d \times r}$ an orthonormal basis of the vector space $\{\boldsymbol{\mu}_i\}_{i \in [k]} \text{ and } \mathbf{N}^T \mathbf{N} = \boldsymbol{I}_n$. Since $\nu_0^2 > \nu_i^2$, the quadratic term $\frac{\|\mathbf{r}\|^2}{2\nu_0^2} - \frac{\|\mathbf{r} - \mathbf{R}^T \boldsymbol{\mu}_i\|^2}{2\nu_i^2}$ is maximized at 480 $\mathbf{r} = \frac{\nu_0^2 \mathbf{R}^T \boldsymbol{\mu}_i}{\nu_0^2 - \nu_i^2}$. Therefore,

$$\frac{\|\mathbf{r}\|^2}{2\nu_0^2} - \frac{\|\mathbf{r} - \mathbf{R}^T \boldsymbol{\mu}_i\|^2}{2\nu_i^2} \le \frac{\nu_0^4 \|\mathbf{R}^T \boldsymbol{\mu}_i\|^2}{2\nu_0^2(\nu_0^2 - \nu_i^2)^2} - \frac{1}{2\nu_i^2} \left(\frac{\nu_0^2}{\nu_0^2 - \nu_i^2} - 1\right)^2 \|\mathbf{R}^T \boldsymbol{\mu}_i\|^2 = \frac{\|\boldsymbol{\mu}_i\|^2}{2(\nu_0^2 - \nu_i^2)}.$$

481 Hence, for $\|\mathbf{n}\|^2 \ge \frac{\nu_0^2 + \nu_{\max}^2}{2} d$ and $\|\boldsymbol{\mu}_i\|^2 \le \frac{\nu_0^2 - \nu_i^2}{2} \left(\log \left(\frac{\nu_i^2}{\nu_0^2} \right) - \frac{\nu_i^2}{2\nu_0^2} + \frac{\nu_0^2}{2\nu_i^2} \right) d$, we have

$$\begin{split} \rho_{i}(\mathbf{x}) &= \left(\frac{\nu_{0}^{2}}{\nu_{i}^{2}}\right)^{d/2} \exp\left(\left(\frac{1}{2\nu_{0}^{2}} - \frac{1}{2\nu_{i}^{2}}\right) \|\mathbf{n}\|^{2} + \left(\frac{\|\mathbf{r}\|^{2}}{2\nu_{0}^{2}} - \frac{\|\mathbf{r} - \mathbf{R}^{T}\boldsymbol{\mu}_{i}\|^{2}}{2\nu_{i}^{2}}\right)\right) \\ &\leq \left(\frac{\nu_{0}^{2}}{\nu_{i}^{2}}\right)^{d/2} \exp\left(\left(\frac{1}{2\nu_{0}^{2}} - \frac{1}{2\nu_{i}^{2}}\right) \frac{\nu_{0}^{2} + \nu_{i}^{2}}{2} d + \frac{\|\boldsymbol{\mu}_{i}\|^{2}}{2(\nu_{0}^{2} - \nu_{i}^{2})}\right) \\ &= \exp\left(-\left(\log\left(\frac{\nu_{i}^{2}}{\nu_{0}^{2}}\right) - \frac{\nu_{i}^{2}}{2\nu_{0}^{2}} + \frac{\nu_{0}^{2}}{2\nu_{i}^{2}}\right) \frac{d}{2} + \frac{\|\boldsymbol{\mu}_{i}\|^{2}}{2(\nu_{0}^{2} - \nu_{i}^{2})}\right) \\ &\leq \exp\left(-\left(\log\left(\frac{\nu_{i}^{2}}{\nu_{0}^{2}}\right) - \frac{\nu_{i}^{2}}{2\nu_{0}^{2}} + \frac{\nu_{0}^{2}}{2\nu_{i}^{2}}\right) \frac{d}{4}\right). \end{split}$$

Notice that for function $f(z) = \log z - \frac{z}{2} + \frac{1}{2z}$, we have f(1) = 0 and $\frac{d}{dz}f(z) = \frac{1}{z} - \frac{1}{2} - \frac{1}{2z^2} = -\frac{1}{2}\left(\frac{1}{z}-1\right)^2 < 0$ when $z \in (0,1)$. Thus, $\log\left(\frac{\nu_i^2}{\nu_0^2}\right) - \frac{\nu_i^2}{2\nu_0^2} + \frac{\nu_0^2}{2\nu_i^2}$ is a positive constant for $\nu_i < \nu_0$, i.e., $\rho_i(\mathbf{x}) = \exp(-\Omega(d))$. Therefore we finish the proof of Lemma 4.

Lemma 4 implies that when $||\mathbf{n}||$ is large, the Gaussian mode $P^{(0)}$ dominates other modes $P^{(i)}$. To bound $||\mathbf{n}_t||$, we first consider a simpler case that $||\mathbf{n}_{t-1}||$ is large. Intuitively, the following lemma proves that when the previous state \mathbf{n}_{t-1} is far from a mode, a single step of Langevin dynamics with bounded step size is not enough to find the mode.

Lemma 5. Suppose $\delta_t \leq \nu_0^2$ and $\|\mathbf{n}_{t-1}\|^2 > 36\nu_0^2 d$, then for \mathbf{n}_t following from equation 5, we have 490 $\|\mathbf{n}_t\|^2 \geq \nu_0^2 d$ with probability at least $1 - \exp(-\Omega(d))$.

⁴⁹¹ *Proof of Lemma 5.* From the recursion of n_t in equation 5 we have

$$\mathbf{n}_{t} = \mathbf{n}_{t-1} + \frac{\delta_{t}}{2} \mathbf{N}^{T} \nabla_{\mathbf{x}} \log P(\mathbf{x}_{t-1}) + \sqrt{\delta_{t}} \boldsymbol{\epsilon}_{t}^{(\mathbf{n})}$$

$$= \mathbf{n}_{t-1} - \frac{\delta_{t}}{2} \sum_{i=0}^{k} \frac{P^{(i)}(\mathbf{x}_{t-1})}{P(\mathbf{x}_{t-1})} \cdot \frac{\mathbf{N}^{T}(\mathbf{x}_{t-1} - \boldsymbol{\mu}_{i})}{\nu_{i}^{2}} + \sqrt{\delta_{t}} \boldsymbol{\epsilon}_{t}^{(\mathbf{n})}$$

$$= \left(1 - \frac{\delta_{t}}{2} \sum_{i=0}^{k} \frac{P^{(i)}(\mathbf{x}_{t-1})}{P(\mathbf{x}_{t-1})} \cdot \frac{1}{\nu_{i}^{2}}\right) \mathbf{n}_{t-1} + \sqrt{\delta_{t}} \boldsymbol{\epsilon}_{t}^{(\mathbf{n})}.$$
(6)

 $\text{By Lemma 4, we have } \frac{P^{(i)}(\mathbf{x}_{j-1})}{P^{(0)}(\mathbf{x}_{j-1})} \leq \exp(-\Omega(d)) \text{ for all } i \in [k], \text{ therefore } i \in [k]$

$$1 - \frac{\delta_t}{2} \sum_{i=0}^k \frac{P^{(i)}(\mathbf{x}_{t-1})}{P(\mathbf{x}_{t-1})} \cdot \frac{1}{\nu_i^2} \ge 1 - \frac{\delta_t}{2} \cdot \frac{1}{\nu_0^2} - \frac{\delta_t}{2} \sum_{i \in [k]} \frac{w_i P^{(i)}(\mathbf{x}_{t-1})}{w_0 P^{(0)}(\mathbf{x}_{t-1})} \cdot \frac{1}{\nu_i^2} \ge 1 - \frac{1}{2} - \exp(-\Omega(d)) > \frac{1}{3}.$$
(7)

493 On the other hand, from $\boldsymbol{\epsilon}_t^{(\mathbf{n})} \sim \mathcal{N}(\mathbf{0}_n, \boldsymbol{I}_n)$ we know $\frac{\langle \mathbf{n}_{t-1}, \boldsymbol{\epsilon}_t^{(\mathbf{n})} \rangle}{\|\mathbf{n}_{t-1}\|} \sim \mathcal{N}(0, 1)$ for any fixed $\mathbf{n}_{t-1} \neq \mathbf{0}_n$, 494 hence by Lemma 2 we have

$$\mathbb{P}\left(\frac{\langle \mathbf{n}_{t-1}, \boldsymbol{\epsilon}_t^{(\mathbf{n})} \rangle}{\|\mathbf{n}_{t-1}\|} \ge \frac{\sqrt{d}}{4}\right) = \mathbb{P}\left(\frac{\langle \mathbf{n}_{t-1}, \boldsymbol{\epsilon}_t^{(\mathbf{n})} \rangle}{\|\mathbf{n}_{t-1}\|} \le -\frac{\sqrt{d}}{4}\right) \le \frac{4}{\sqrt{2\pi d}} \exp\left(-\frac{d}{32}\right)$$
(8)

495 Combining equation 6, equation 7 and equation 8 gives that

$$\begin{split} \|\mathbf{n}_{t}\|^{2} &\geq \left(\frac{1}{3}\right)^{2} \|\mathbf{n}_{t-1}\|^{2} - 2\nu_{0} |\langle \mathbf{n}_{t-1}, \boldsymbol{\epsilon}_{t}^{(\mathbf{n})} \rangle| \\ &\geq \frac{1}{9} \|\mathbf{n}_{t-1}\|^{2} - \frac{\nu_{0}\sqrt{d}}{2} \|\mathbf{n}_{t-1}\| \\ &\geq \frac{1}{9} \cdot 36\nu_{0}^{2}d - \frac{\nu_{0}\sqrt{d}}{2} \cdot 6\nu_{0}\sqrt{d} \\ &= \nu_{0}^{2}d \end{split}$$

with probability at least $1 - \frac{8}{\sqrt{2\pi d}} \exp\left(-\frac{d}{32}\right) = 1 - \exp(-\Omega(d))$. This proves Lemma 5.

We then proceed to bound $\|\mathbf{n}_t\|$ iteratively for $\|\mathbf{n}_{t-1}\|^2 \leq 36\nu_0^2 d$. Recall that equation 5 gives

$$\mathbf{n}_t = \mathbf{n}_{t-1} + \frac{\delta_t}{2} \mathbf{N}^T \nabla_{\mathbf{x}} \log P(\mathbf{x}_{t-1}) + \sqrt{\delta_t} \boldsymbol{\epsilon}_t^{(\mathbf{n})}.$$

We notice that the difficulty of solving \mathbf{n}_t exhibits in the dependence of $\log P(\mathbf{x}_{t-1})$ on \mathbf{r}_{t-1} . Since $P = \sum_{i=0}^k w_i P^{(i)} = \sum_{i=0}^k w_i \mathcal{N}(\boldsymbol{\mu}_i, \nu_i^2 \boldsymbol{I}_d)$, we can rewrite the score function as

$$\nabla_{\mathbf{x}} \log P(\mathbf{x}) = \frac{\nabla_{\mathbf{x}} P(\mathbf{x})}{P(\mathbf{x})} = -\sum_{i=0}^{k} \frac{P^{(i)}(\mathbf{x})}{P(\mathbf{x})} \cdot \frac{\mathbf{x} - \boldsymbol{\mu}_{i}}{\nu_{i}^{2}} = -\frac{\mathbf{x}}{\nu_{0}^{2}} + \sum_{i \in [k]} \frac{P^{(i)}(\mathbf{x})}{P(\mathbf{x})} \left(\frac{\mathbf{x}}{\nu_{0}^{2}} - \frac{\mathbf{x} - \boldsymbol{\mu}_{i}}{\nu_{i}^{2}}\right).$$
(9)

Now, instead of directly working with \mathbf{n}_t , we consider a surrogate recursion $\hat{\mathbf{n}}_t$ such that $\hat{\mathbf{n}}_0 = \mathbf{n}_0$ and for all $t \ge 1$,

$$\hat{\mathbf{n}}_t = \hat{\mathbf{n}}_{t-1} - \frac{\delta_t}{2\nu_0^2} \hat{\mathbf{n}}_{t-1} + \sqrt{\delta_t} \boldsymbol{\epsilon}_t^{(\mathbf{n})}.$$
(10)

The advantage of the surrogate recursion is that $\hat{\mathbf{n}}_t$ is independent of \mathbf{r} , thus we can obtain the closed-form solution to $\hat{\mathbf{n}}_t$. Before we proceed to bound $\hat{\mathbf{n}}_t$, we first show that $\hat{\mathbf{n}}_t$ is sufficiently close to the original recursion \mathbf{n}_t in the following lemma.

Lemma 6. For any
$$t \ge 1$$
, given that $\delta_j \le \nu_0^2$ and $\frac{\nu_0^2 + \nu_{\max}^2}{2} d \le \|\mathbf{n}_{j-1}\|^2 \le 36\nu_0^2 d$ for all $j \in [t]$ and
 $\|\boldsymbol{\mu}_i\|^2 \le \frac{\nu_0^2 - \nu_i^2}{2} \left(\log \left(\frac{\nu_i^2}{\nu_0^2} \right) - \frac{\nu_i^2}{2\nu_0^2} + \frac{\nu_0^2}{2\nu_i^2} \right) d$ for all $i \in [k]$, we have $\|\hat{\mathbf{n}}_t - \mathbf{n}_t\| \le \frac{t}{\exp(\Omega(d))} \sqrt{d}$.

⁵⁰⁷ *Proof of Lemma 6.* Upon comparing equation 5 and equation 10, by equation 9 we have that for all $j \in [t]$,

$$\begin{aligned} \|\hat{\mathbf{n}}_{j} - \mathbf{n}_{j}\| &= \left\| \hat{\mathbf{n}}_{j-1} - \frac{\delta_{j}}{2\nu_{0}^{2}} \hat{\mathbf{n}}_{j-1} - \mathbf{n}_{j-1} - \frac{\delta_{j}}{2} \mathbf{N}^{T} \nabla_{\mathbf{x}} \log P(\mathbf{x}_{j-1}) \right\| \\ &= \left\| \left(1 - \frac{\delta_{j}}{2\nu_{0}^{2}} \right) (\hat{\mathbf{n}}_{j-1} - \mathbf{n}_{j-1}) + \frac{\delta_{j}}{2} \sum_{i \in [k]} \frac{P^{(i)}(\mathbf{x}_{j-1})}{P(\mathbf{x}_{j-1})} \left(\frac{1}{\nu_{i}^{2}} - \frac{1}{\nu_{0}^{2}} \right) \mathbf{n}_{j-1} \right\| \\ &\leq \left(1 - \frac{\delta_{j}}{2\nu_{0}^{2}} \right) \|\hat{\mathbf{n}}_{j-1} - \mathbf{n}_{j-1}\| + \sum_{i \in [k]} \frac{\delta_{j}}{2} \frac{P^{(i)}(\mathbf{x}_{j-1})}{P(\mathbf{x}_{j-1})} \left(\frac{1}{\nu_{i}^{2}} - \frac{1}{\nu_{0}^{2}} \right) \|\mathbf{n}_{j-1} \\ &\leq \|\hat{\mathbf{n}}_{j-1} - \mathbf{n}_{j-1}\| + \sum_{i \in [k]} \frac{\delta_{j}}{2} \frac{P^{(i)}(\mathbf{x}_{j-1})}{P^{(0)}(\mathbf{x}_{j-1})} \left(\frac{1}{\nu_{i}^{2}} - \frac{1}{\nu_{0}^{2}} \right) 6\nu_{0}\sqrt{d}. \end{aligned}$$

By Lemma 4, we have $\frac{P^{(i)}(\mathbf{x}_{j-1})}{P^{(0)}(\mathbf{x}_{j-1})} \leq \exp(-\Omega(d))$ for all $i \in [k]$, hence we obtain a recursive bound

$$\|\hat{\mathbf{n}}_{j} - \mathbf{n}_{j}\| \le \|\hat{\mathbf{n}}_{j-1} - \mathbf{n}_{j-1}\| + \frac{1}{\exp(\Omega(d))}\sqrt{d}.$$

510 Finally, by $\hat{\mathbf{n}}_0 = \mathbf{n}_0$, we have

$$\|\hat{\mathbf{n}}_{t} - \mathbf{n}_{t}\| = \sum_{j \in [t]} (\|\hat{\mathbf{n}}_{j} - \mathbf{n}_{j}\| - \|\hat{\mathbf{n}}_{j-1} - \mathbf{n}_{j-1}\|) \le \frac{t}{\exp(\Omega(d))} \sqrt{d}.$$

511 Hence we obtain Lemma 6.

We then proceed to analyze $\hat{\mathbf{n}}_t$, The following lemma gives us the closed-form solution of $\hat{\mathbf{n}}_t$. We slightly abuse the notations here, e.g., $\prod_{i=c_1}^{c_2} \left(1 - \frac{\delta_i}{2\nu_0^2}\right) = 1$ and $\sum_{j=c_1}^{c_2} \delta_j = 0$ for $c_1 > c_2$.

Lemma 7. For all
$$t \ge 0$$
, $\hat{\mathbf{n}}_t \sim \mathcal{N}\left(\prod_{i=1}^t \left(1 - \frac{\delta_i}{2\nu_0^2}\right)\mathbf{n}_0, \sum_{j=1}^t \prod_{i=j+1}^t \left(1 - \frac{\delta_i}{2\nu_0^2}\right)^2 \delta_j \mathbf{I}_n\right)$, where
the mean and covariance satisfy $\prod_{i=1}^t \left(1 - \frac{\delta_i}{2\nu_0^2}\right)^2 + \frac{1}{\nu_0^2} \sum_{j=1}^t \prod_{i=j+1}^t \left(1 - \frac{\delta_i}{2\nu_0^2}\right)^2 \delta_j \ge 1$.

Proof of Lemma 7. We prove the two properties by induction. When t = 0, they are trivial. Suppose they hold for t - 1, then for the distribution of $\hat{\mathbf{n}}_t$, we have

$$\begin{split} \hat{\mathbf{n}}_{t} &= \hat{\mathbf{n}}_{t-1} - \frac{\delta_{t}}{2\nu_{0}^{2}} \hat{\mathbf{n}}_{t-1} + \sqrt{\delta_{t}} \boldsymbol{\epsilon}_{t}^{(\mathbf{n})} \\ &\sim \mathcal{N}\left(\left(1 - \frac{\delta_{t}}{2\nu_{0}^{2}} \right) \prod_{i=1}^{t-1} \left(1 - \frac{\delta_{i}}{2\nu_{0}^{2}} \right) \mathbf{n}_{0}, \ \left(1 - \frac{\delta_{t}}{2\nu_{0}^{2}} \right)^{2} \sum_{j=1}^{t-1} \prod_{i=j+1}^{t-1} \left(1 - \frac{\delta_{i}}{2\nu_{0}^{2}} \right)^{2} \delta_{j} \boldsymbol{I}_{n} + \delta_{t} \boldsymbol{I}_{n} \right) \\ &= \mathcal{N}\left(\prod_{i=1}^{t} \left(1 - \frac{\delta_{i}}{2\nu_{0}^{2}} \right) \mathbf{n}_{0}, \ \sum_{j=1}^{t} \prod_{i=j+1}^{t} \left(1 - \frac{\delta_{i}}{2\nu_{0}^{2}} \right)^{2} \delta_{j} \boldsymbol{I}_{n} \right). \end{split}$$

518 For the second property,

$$\begin{split} \prod_{i=1}^{t} \left(1 - \frac{\delta_i}{2\nu_0^2}\right)^2 + \frac{1}{\nu_0^2} \sum_{j=1}^{t} \prod_{i=j+1}^{t} \left(1 - \frac{\delta_i}{2\nu_0^2}\right)^2 \delta_j \\ &= \left(1 - \frac{\delta_t}{2\nu_0^2}\right)^2 \left(\prod_{i=1}^{t-1} \left(1 - \frac{\delta_i}{2\nu_0^2}\right)^2 + \frac{1}{\nu_0^2} \sum_{j=1}^{t-1} \prod_{i=j+1}^{t-1} \left(1 - \frac{\delta_i}{2\nu_0^2}\right)^2 \delta_j\right) + \frac{1}{\nu_0^2} \delta_t \\ &\geq \left(1 - \frac{\delta_t}{2\nu_0^2}\right)^2 + \frac{1}{\nu_0^2} \delta_t = 1 + \frac{\delta_t^2}{4\nu_0^4} \ge 1. \end{split}$$

519 Hence we finish the proof of Lemma 7.

Armed with Lemma 7, we are now ready to establish the lower bound on $\|\hat{\mathbf{n}}_t\|$. For simplicity, denote $\alpha := \prod_{i=1}^t \left(1 - \frac{\delta_i}{2\nu_0^2}\right)^2$ and $\beta := \frac{1}{\nu_0^2} \sum_{j=1}^t \prod_{i=j+1}^t \left(1 - \frac{\delta_i}{2\nu_0^2}\right)^2 \delta_j$. By Lemma 7 we know $\hat{\mathbf{n}}_t \sim \mathcal{N}(\alpha \mathbf{n}_0, \beta \nu_0^2 \mathbf{I}_n)$, so we can write $\hat{\mathbf{n}}_t = \alpha \mathbf{n}_0 + \sqrt{\beta}\nu_0 \epsilon$, where $\epsilon \sim \mathcal{N}(\mathbf{0}_n, \mathbf{I}_n)$.

523 **Lemma 8.** Given that $\|\hat{\mathbf{n}}_0\|^2 \ge \frac{3\nu_0^2 + \nu_{\max}^2}{4}d$, we have $\|\hat{\mathbf{n}}_t\|^2 \ge \frac{5\nu_0^2 + 3\nu_{\max}^2}{8}d$ with probability at least 524 $1 - \exp(-\Omega(d))$.

525 *Proof of Lemma 8.* By $\hat{\mathbf{n}}_t = \alpha \mathbf{n}_0 + \sqrt{\beta} \nu_0 \boldsymbol{\epsilon}$ we have

$$\|\hat{\mathbf{n}}_t\|^2 = \alpha^2 \|\mathbf{n}_0\|^2 + \beta \nu_0^2 \|\boldsymbol{\epsilon}\|^2 + 2\alpha \sqrt{\beta} \nu_0 \langle \mathbf{n}_0, \boldsymbol{\epsilon} \rangle$$

526 By Lemma 1 we can bound

$$\begin{split} \mathbb{P}\left(\left\|\boldsymbol{\epsilon}\right\|^{2} \leq \frac{3\nu_{0}^{2} + \nu_{\max}^{2}}{4\nu_{0}^{2}}d\right) &= \mathbb{P}\left(\left\|\boldsymbol{\epsilon}\right\|^{2} \leq d - 2\sqrt{d \cdot \left(\frac{\nu_{0}^{2} - \nu_{\max}^{2}}{8\nu_{0}^{2}}\right)^{2}d}\right) \\ &\leq \mathbb{P}\left(\left\|\boldsymbol{\epsilon}\right\|^{2} \leq (n - 1) - 2\sqrt{(n - 1)\left(\frac{\nu_{0}^{2} - \nu_{\max}^{2}}{8\nu_{0}^{2}}\right)^{2}\frac{d}{2}}\right) \\ &\leq \exp\left(-\left(\frac{\nu_{0}^{2} - \nu_{\max}^{2}}{8\nu_{0}^{2}}\right)^{2}\frac{d}{2}\right), \end{split}$$

where the second last step follows from the assumption d - n = r = o(d). Since $\epsilon \sim \mathcal{N}(\mathbf{0}_n, \mathbf{I}_n)$, we know $\frac{\langle \mathbf{n}_0, \epsilon \rangle}{\| \mathbf{n}_0 \|} \sim \mathcal{N}(0, 1)$. Therefore by Lemma 2,

$$\mathbb{P}\left(\frac{\langle \mathbf{n}_{0}, \boldsymbol{\epsilon} \rangle}{\|\mathbf{n}_{0}\|} \leq -\frac{\nu_{0}^{2} - \nu_{\max}^{2}}{4\nu_{0}\sqrt{3\nu_{0}^{2} + \nu_{\max}^{2}}}\sqrt{d}\right) \leq \frac{4\nu_{0}\sqrt{3\nu_{0}^{2} + \nu_{\max}^{2}}}{\sqrt{2\pi}(\nu_{0}^{2} - \nu_{\max}^{2})\sqrt{d}}\exp\left(-\frac{(\nu_{0}^{2} - \nu_{\max}^{2})^{2}d}{32\nu_{0}^{2}(3\nu_{0}^{2} + \nu_{\max}^{2})}\right)$$

529 Conditioned on $\|\hat{\mathbf{n}}_{0}\|^{2} \geq \frac{3\nu_{0}^{2} + \nu_{\max}^{2}}{4}d$, $\|\boldsymbol{\epsilon}\|^{2} > \frac{3\nu_{0}^{2} + \nu_{\max}^{2}}{4\nu_{0}^{2}}d$ and $\frac{1}{\|\mathbf{n}_{0}\|}\langle\mathbf{n}_{0},\boldsymbol{\epsilon}\rangle > -\frac{\nu_{0}^{2} - \nu_{\max}^{2}}{4\nu_{0}\sqrt{3\nu_{0}^{2} + \nu_{\max}^{2}}}\sqrt{d}$, 530 since Lemma 7 gives $\alpha^{2} + \beta \geq 1$ we have

$$\begin{split} \|\hat{\mathbf{n}}_{t}\|^{2} &= \alpha^{2} \|\mathbf{n}_{0}\|^{2} + \beta\nu_{0}^{2} \|\boldsymbol{\epsilon}\|^{2} + 2\alpha\sqrt{\beta}\nu_{0}\langle\mathbf{n}_{0},\boldsymbol{\epsilon}\rangle \\ &\geq \alpha^{2} \|\mathbf{n}_{0}\|^{2} + \beta\nu_{0}^{2} \|\boldsymbol{\epsilon}\|^{2} - 2\alpha\sqrt{\beta}\nu_{0} \|\mathbf{n}_{0}\| \frac{\nu_{0}^{2} - \nu_{\max}^{2}}{4\nu_{0}\sqrt{3\nu_{0}^{2} + \nu_{\max}^{2}}}\sqrt{d} \\ &\geq \alpha^{2} \|\mathbf{n}_{0}\|^{2} + \beta\nu_{0}^{2} \|\boldsymbol{\epsilon}\|^{2} - 2\alpha\sqrt{\beta}\nu_{0} \|\mathbf{n}_{0}\| \|\boldsymbol{\epsilon}\| \cdot \frac{\nu_{0}^{2} - \nu_{\max}^{2}}{6\nu_{0}^{2} + 2\nu_{\max}^{2}} \\ &\geq \left(1 - \frac{\nu_{0}^{2} - \nu_{\max}^{2}}{6\nu_{0}^{2} + 2\nu_{\max}^{2}}\right) \left(\alpha^{2} \|\mathbf{n}_{0}\|^{2} + \beta\nu_{0}^{2} \|\boldsymbol{\epsilon}\|^{2}\right) \\ &\geq \frac{5\nu_{0}^{2} + 3\nu_{\max}^{2}}{6\nu_{0}^{2} + 2\nu_{\max}^{2}} \left(\alpha^{2} + \beta\right) \cdot \frac{3\nu_{0}^{2} + \nu_{\max}^{2}}{4} d \\ &\geq \frac{5\nu_{0}^{2} + 3\nu_{\max}^{2}}{8} d. \end{split}$$

Hence by union bound, we complete the proof of Lemma 8.

⁵³² Upon having all the above lemmas, we are now ready to establish Proposition 3 by induction. Suppose ⁵³³ the theorem holds for all T values of $1, \dots, T-1$. We consider the following 3 cases:

• If there exists some $t \in [T]$ such that $\delta_t > \nu_0^2$, by Lemma 3 we know that with probability at least $1 - \exp(-\Omega(d))$, we have $\|\mathbf{n}_t\|^2 \ge \frac{3\nu_0^2 + \nu_{\max}^2}{4}d$, thus the problem reduces to the two sub-arrays $\mathbf{n}_0, \cdots, \mathbf{n}_{t-1}$ and $\mathbf{n}_t, \cdots, \mathbf{n}_T$, which can be solved by induction.

• Suppose $\delta_t \leq \nu_0^2$ for all $t \in [T]$. If there exists some $t \in [T]$ such that $\|\mathbf{n}_{t-1}\|^2 > 36\nu_0^2 d$, by Lemma 5 we know that with probability at least $1 - \exp(-\Omega(d))$, we have $\|\mathbf{n}_t\|^2 \geq \nu_0^2 d > \frac{3\nu_0^2 + \nu_{\max}^2}{4} d$, thus the problem similarly reduces to the two sub-arrays $\mathbf{n}_0, \cdots, \mathbf{n}_{t-1}$ and $\mathbf{n}_t, \cdots, \mathbf{n}_T$, which can be solved by induction.

• Suppose
$$\delta_t \leq \nu_0^2$$
 and $\|\mathbf{n}_{t-1}\|^2 \leq 36\nu_0^2 d$ for all $t \in [T]$. Conditioned on $\|\mathbf{n}_{t-1}\|^2 > \frac{\nu_0^2 + \nu_{\max}^2}{2} d$ for all $t \in [T]$, by Lemma 6 we have that for $T = \exp(\mathcal{O}(d))$,

$$\|\hat{\mathbf{n}}_T - \mathbf{n}_T\| < \left(\sqrt{\frac{5\nu_0^2 + 3\nu_{\max}^2}{8}} - \sqrt{\frac{\nu_0^2 + \nu_{\max}^2}{2}}\right)\sqrt{d}.$$

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By Lemma 8 we have that with probability at least $1 - \exp(-\Omega(d))$,

$$\|\hat{\mathbf{n}}_T\|^2 \ge \frac{5\nu_0^2 + 3\nu_{\max}^2}{8}d$$

Combining the two inequalities implies the desired bound

$$\|\mathbf{n}_T\| \ge \|\hat{\mathbf{n}}_T\| - \|\hat{\mathbf{n}}_T - \mathbf{n}_T\| > \sqrt{\frac{\nu_0^2 + \nu_{\max}^2}{2}d}.$$

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Hence by induction we obtain
$$\|\mathbf{n}_t\|^2 > \frac{\nu_0^2 + \nu_{\max}^2}{2} d$$
 for all $t \in [T]$ with probability at least

$$(1 - (T - 1)\exp(-\Omega(d))) \cdot (1 - \exp(-\Omega(d))) \ge 1 - T\exp(-\Omega(d)).$$

⁵⁴⁶ Therefore we complete the proof of Proposition 3.

547 Finally, combining Propositions 2 and 3 finishes the proof of Theorem 1.

548 A.2 Proof of Theorem 2: Annealed Langevin Dynamics under Gaussian Mixtures

To establish Theorem 2, we first note from Proposition 1 that perturbing a Gaussian distribution $\mathcal{N}(\boldsymbol{\mu}, \nu^2 \boldsymbol{I}_d)$ with noise level σ results in a Gaussian distribution $\mathcal{N}(\boldsymbol{\mu}, (\nu^2 + \sigma^2)\boldsymbol{I}_d)$. Therefore, for a Gaussian mixture $P = \sum_{i=0}^k w_i P^{(i)} = \sum_{i=0}^k w_i \mathcal{N}(\boldsymbol{\mu}_i, \nu_i^2 \boldsymbol{I}_d)$, the perturbed distribution of noise level σ is

$$P_{\sigma} = \sum_{i=0}^{\kappa} w_i \mathcal{N}(\boldsymbol{\mu}_i, (\nu_i^2 + \sigma^2) \boldsymbol{I}_d).$$

553 Similar to the proof of Theorem 1, we decompose

$$\mathbf{x}_t = \mathbf{R}\mathbf{r}_t + \mathbf{N}\mathbf{n}_t$$
, and $\boldsymbol{\epsilon}_t = \mathbf{R}\boldsymbol{\epsilon}_t^{(\mathbf{r})} + \mathbf{N}\boldsymbol{\epsilon}_t^{(\mathbf{n})}$,

where $\mathbf{R} \in \mathbb{R}^{d \times r}$ an orthonormal basis of the vector space $\{\boldsymbol{\mu}_i\}_{i \in [k]}$ and $\mathbf{N} \in \mathbb{R}^{d \times n}$ an orthonormal basis of the null space of $\{\boldsymbol{\mu}_i\}_{i \in [k]}$. Now, we prove Theorem 2 by applying the techniques developed in Appendix A.1 via substituting ν^2 with $\nu^2 + \sigma_t^2$ at time step t.

First, by Proposition 2, suppose that the sample is initialized in the distribution $P_{\sigma_0}^{(0)}$, then with probability at least $1 - \exp(-\Omega(d))$, we have

$$\|\mathbf{n}_0\|^2 \ge \frac{3(\nu_0^2 + \sigma_0^2) + (\nu_{\max}^2 + \sigma_0^2)}{4}d = \frac{3\nu_0^2 + \nu_{\max}^2 + 4\sigma_0^2}{4}d.$$
 (11)

Then, with the assumption that the initialization satisfies $\|\mathbf{n}_0\|^2 \ge \frac{3\nu_0^2 + \nu_{\max}^2 + 4\sigma_0^2}{4}d$, the following proposition similar to Proposition 3 shows that $\|\mathbf{n}_t\|$ remains large with high probability.

Proposition 4. Consider a data distribution P satisfies the constraints specified in Theorem 2. We follow annealed Langevin dynamics for $T = \exp(\mathcal{O}(d))$ steps with noise level $c_{\sigma} \ge \sigma_0 \ge$ $\sigma_1 \ge \sigma_2 \ge \cdots \ge \sigma_T \ge 0$ for some constant $c_{\sigma} > 0$. Suppose that the initial sample satisfies $\|\mathbf{n}_0\|^2 \ge \frac{3\nu_0^2 + \nu_{\max}^2 + 4\sigma_0^2}{4}d$, then with probability at least $1 - T \cdot \exp(-\Omega(d))$, we have that $\|\mathbf{n}_t\|^2 >$ $\frac{\nu_0^2 + \nu_{\max}^2 + 2\sigma_t^2}{2}d$ for all $t \in \{0\} \cup [T]$.

Proof of Proposition 4. We prove Proposition 4 by induction. Suppose the theorem holds for all T values of $1, \dots, T-1$. We consider the following 3 cases:

- If there exists some $t \in [T]$ such that $\delta_t > \nu_0^2 + \sigma_t^2$, by Lemma 3 we know that with probability at least $1 - \exp(-\Omega(d))$, we have $\|\mathbf{n}_t\|^2 \ge \frac{3(\nu_0^2 + \sigma_t^2) + (\nu_{\max}^2 + \sigma_t^2)}{4}d = \frac{3\nu_0^2 + \nu_{\max}^2 + 4\sigma_t^2}{4}d$, thus the problem reduces to the two sub-arrays $\mathbf{n}_0, \cdots, \mathbf{n}_{t-1}$ and $\mathbf{n}_t, \cdots, \mathbf{n}_T$, which can be solved by induction.
- Suppose $\delta_t \leq \nu_0^2 + \sigma_t^2$ for all $t \in [T]$. If there exists some $t \in [T]$ such that $\|\mathbf{n}_{t-1}\|^2 > 36(\nu_0^2 + \sigma_{t-1}^2)d \geq 36(\nu_0^2 + \sigma_t^2)d$, by Lemma 5 we know that with probability at least 1 $\exp(-\Omega(d))$, we have $\|\mathbf{n}_t\|^2 \geq (\nu_0^2 + \sigma_t^2)d > \frac{3\nu_0^2 + \nu_{\max}^2 + 4\sigma_t^2}{4}d$, thus the problem similarly reduces to the two sub-arrays $\mathbf{n}_0, \cdots, \mathbf{n}_{t-1}$ and $\mathbf{n}_t, \cdots, \mathbf{n}_T$, which can be solved by induction.

577 578 • Suppose $\delta_t \leq \nu_0^2 + \sigma_t^2$ and $\|\mathbf{n}_{t-1}\|^2 \leq 36(\nu_0^2 + \sigma_{t-1}^2)d$ for all $t \in [T]$. Consider a surrogate sequence $\hat{\mathbf{n}}_t$ such that $\hat{\mathbf{n}}_0 = \mathbf{n}_0$ and for all $t \geq 1$,

$$\hat{\mathbf{n}}_t = \hat{\mathbf{n}}_{t-1} - \frac{\delta_t}{2\nu_0^2 + 2\sigma_t^2} \hat{\mathbf{n}}_{t-1} + \sqrt{\delta_t} \boldsymbol{\epsilon}_t^{(\mathbf{n})}$$

Since $\nu_0 > \nu_i$ and $c_\sigma \ge \sigma_t$ for all $t \in \{0\} \cup [T]$, we have $\frac{\nu_i^2 + c_\sigma^2}{\nu_0^2 + c_\sigma^2} \ge \frac{\nu_i^2 + \sigma_t^2}{\nu_0^2 + \sigma_t^2}$. Notice that for function $f(z) = \log z - \frac{z}{2} + \frac{1}{2z}$, we have $\frac{\mathrm{d}}{\mathrm{d}z}f(z) = \frac{1}{z} - \frac{1}{2} - \frac{1}{2z^2} = -\frac{1}{2}\left(\frac{1}{z} - 1\right)^2 \le 0$. Thus, by the assumption 579 580

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$$\|\boldsymbol{\mu}_{i} - \boldsymbol{\mu}_{0}\|^{2} \leq \frac{\nu_{0}^{2} - \nu_{i}^{2}}{2} \left(\log \left(\frac{\nu_{i}^{2} + c_{\sigma}^{2}}{\nu_{0}^{2} + c_{\sigma}^{2}} \right) - \frac{\nu_{i}^{2} + c_{\sigma}^{2}}{2\nu_{0}^{2} + c_{\sigma}^{2}} + \frac{\nu_{0}^{2} + c_{\sigma}^{2}}{2\nu_{i}^{2} + c_{\sigma}^{2}} \right) d,$$

we have that for all $t \in [T]$, 582

$$\|\boldsymbol{\mu}_{i} - \boldsymbol{\mu}_{0}\|^{2} \leq \frac{\nu_{0}^{2} - \nu_{i}^{2}}{2} \left(\log \left(\frac{\nu_{i}^{2} + \sigma_{t}^{2}}{\nu_{0}^{2} + \sigma_{t}^{2}} \right) - \frac{\nu_{i}^{2} + \sigma_{t}^{2}}{2\nu_{0}^{2} + \sigma_{t}^{2}} + \frac{\nu_{0}^{2} + \sigma_{t}^{2}}{2\nu_{i}^{2} + \sigma_{t}^{2}} \right) d$$

Conditioned on $\|\mathbf{n}_{t-1}\|^2 > \frac{\nu_0^2 + \nu_{\max}^2 + 2\sigma_{t-1}^2}{2}d$ for all $t \in [T]$, by Lemma 6 we have that for $T = \exp(\mathcal{O}(d))$, 583 584

$$\|\hat{\mathbf{n}}_T - \mathbf{n}_T\| < \left(\sqrt{\frac{5\nu_0^2 + 3\nu_{\max}^2 + 8\sigma_T^2}{8}} - \sqrt{\frac{\nu_0^2 + \nu_{\max}^2 + 2\sigma_T^2}{2}}\right)\sqrt{d}$$

By Lemma 8 we have that with probability at least $1 - \exp(-\Omega(d))$, 585

$$\|\hat{\mathbf{n}}_T\|^2 \ge \frac{5\nu_0^2 + 3\nu_{\max}^2 + 8\sigma_T^2}{8}d.$$

Combining the two inequalities implies the desired bound 586

$$\|\mathbf{n}_T\| \ge \|\hat{\mathbf{n}}_T\| - \|\hat{\mathbf{n}}_T - \mathbf{n}_T\| > \sqrt{\frac{\nu_0^2 + \nu_{\max}^2 + 2\sigma_T^2}{2}d}.$$

Hence by induction we obtain $\|\mathbf{n}_t\|^2 > \frac{\nu_0^2 + \nu_{\max}^2 + 2\sigma_t^2}{2} d$ for all $t \in \{0\} \cup [T]$ with probability 587 at least 588

$$1 - (T - 1)\exp(-\Omega(d))) \cdot (1 - \exp(-\Omega(d))) \ge 1 - T\exp(-\Omega(d)).$$

Therefore we complete the proof of Proposition 4. 589

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Finally, combining equation 11 and Proposition 4 finishes the proof of Theorem 2. 590

A.3 Proof of Theorem 3: Langevin Dynamics under Sub-Gaussian Mixtures 591

The proof framework is similar to the proof of Theorem 1. To begin with, we validate Assumption 592 2.v. in the following lemma: 593

Lemma 9. For constants ν_0, ν_i, c_ν, c_L satisfying Assumptions 2.iii. and 2.iv., we have $\frac{(1-c_\nu)\nu_0^2-\nu_i^2}{2(1-c_\nu)} > 0$ and $\log \frac{c_\nu \nu_i^2}{(c_L^2+c_\nu c_L)\nu_0^2} - \frac{\nu_i^2}{2(1-c_\nu)\nu_0^2} + \frac{(1-c_\nu)\nu_0^2}{2\nu_i^2} > 0$ are both positive constants. 594 595

Proof of Lemma 9. From Assumption 2.iv. that $\nu_0^2 > \frac{\nu_{\max}^2}{1-c_\nu} \ge \frac{\nu_i^2}{1-c_\nu}$, we easily obtain $\frac{(1-c_\nu)\nu_0^2 - \nu_i^2}{2(1-c_\nu)} > 0$ is a positive constant. For the second property, let $f(z) := \log \frac{c_\nu \nu_i^2}{(c_L^2 + c_\nu c_L)z} - \frac{\nu_i^2}{2(1-c_\nu)z} + \frac{(1-c_\nu)z}{2\nu_i^2}$. 596 597 For any $z > \frac{\nu_i^2}{1-c_v}$, the derivative of f(z) satisfies 598

$$\frac{\mathrm{d}}{\mathrm{d}z}f(z) = -\frac{1}{z} + \frac{\nu_i^2}{2(1-c_\nu)z^2} + \frac{1-c_\nu}{2\nu_i^2} = \frac{\nu_i^2}{2(1-c_\nu)}\left(\frac{1-c_\nu}{\nu_i^2} - \frac{1}{z}\right)^2 > 0.$$

599 Therefore, when $\frac{4(c_L^2 + c_\nu c_L)}{c_\nu (1 - c_\nu)} \leq 1$, we have

$$f(\nu_0^2) > f\left(\frac{\nu_i^2}{1 - c_\nu}\right) = \log \frac{c_\nu(1 - c_\nu)}{c_L^2 + c_\nu c_L} \ge \log 4 > 0.$$

600 When $\frac{4(c_L^2 + c_\nu c_L)}{c_\nu (1 - c_\nu)} > 1$, we have

$$\begin{split} f(\nu_0^2) > f\left(\frac{4(c_L^2 + c_\nu c_L)}{c_\nu(1 - c_\nu)} \frac{\nu_i^2}{1 - c_\nu}\right) &= 2\log\frac{c_\nu(1 - c_\nu)}{2(c_L^2 + c_\nu c_L)} - \frac{c_\nu(1 - c_\nu)}{8(c_L^2 + c_\nu c_L)} + \frac{2(c_L^2 + c_\nu c_L)}{c_\nu(1 - c_\nu)} \\ &\geq 2 - 2\log2 - \frac{2(c_L^2 + c_\nu c_L)}{c_\nu(1 - c_\nu)} - \frac{c_\nu(1 - c_\nu)}{8(c_L^2 + c_\nu c_L)} + \frac{2(c_L^2 + c_\nu c_L)}{c_\nu(1 - c_\nu)} > 2 - 2\log2 - \frac{1}{2} > 0. \end{split}$$

601 Thus we obtain Lemma 9.

Without loss of generality, we assume $\mu_0 = \mathbf{0}_d$. Similar to the proof of Theorem 1, we decompose

$$\mathbf{x}_t = \mathbf{R}\mathbf{r}_t + \mathbf{N}\mathbf{n}_t$$
, and $\boldsymbol{\epsilon}_t = \mathbf{R}\boldsymbol{\epsilon}_t^{(\mathbf{r})} + \mathbf{N}\boldsymbol{\epsilon}_t^{(\mathbf{n})}$

where $\mathbf{R} \in \mathbb{R}^{d \times r}$ an orthonormal basis of the vector space $\{\boldsymbol{\mu}_i\}_{i \in [k]}$ and $\mathbf{N} \in \mathbb{R}^{d \times n}$ an orthonormal basis of the null space of $\{\boldsymbol{\mu}_i\}_{i \in [k]}$. To show $\|\mathbf{x}_t - \boldsymbol{\mu}_i\|^2 > \left(\frac{\nu_0^2}{2} + \frac{\nu_{\max}^2}{2(1-c_\nu)}\right) d$, it suffices to prove $\|\mathbf{n}_t\|^2 > \left(\frac{\nu_0^2}{2} + \frac{\nu_{\max}^2}{2(1-c_\nu)}\right) d$. By Proposition 2, if \mathbf{x}_0 is initialized in the distribution $P^{(0)}$, i.e., $\mathbf{x}_0 \sim P^{(0)}$, since $\nu_0^2 > \frac{1}{1-c_\nu}\nu_{\max}^2$, with probability at least $1 - \exp(-\Omega(d))$ we have

$$\|\mathbf{n}_0\|^2 \ge \left(\frac{3\nu_0^2}{4} + \frac{\nu_{\max}^2}{4(1-c_\nu)}\right) d.$$
(12)

Then, conditioned on $\|\mathbf{n}_0\|^2 \geq \left(\frac{3\nu_0^2}{4} + \frac{\nu_{\max}^2}{4(1-c_\nu)}\right) d$, the following proposition shows that $\|\mathbf{n}_t\|$ remains large with high probability.

Proposition 5. Consider a distribution P satisfying Assumption 2. We follow the Langevin dynamics for $T = \exp(\mathcal{O}(d))$ steps. Suppose that the initial sample satisfies $\|\mathbf{n}_0\|^2 \ge \left(\frac{3\nu_0^2}{4} + \frac{\nu_{\max}^2}{4(1-c_\nu)}\right) d_{tot}$ then with probability at least $1 - T \cdot \exp(-\Omega(d))$, we have that $\|\mathbf{n}_t\|^2 > \left(\frac{\nu_0^2}{2} + \frac{\nu_{\max}^2}{2(1-c_\nu)}\right) d_t$ for all $t \in \{0\} \cup [T]$.

Proof of Proposition 5. Firstly, by Lemma 3, if $\delta_t > \nu_0^2$, since $\nu_0^2 > \frac{\nu_{\max}^2}{1-c_{\nu}}$, we similarly have that $\|\mathbf{n}_t\|^2 \ge \left(\frac{3\nu_0^2}{4} + \frac{\nu_{\max}^2}{4(1-c_{\nu})}\right) d$ with probability at least $1 - \exp(-\Omega(d))$ regardless of the previous state \mathbf{x}_{t-1} . We then consider the case when $\delta_t \le \nu_0^2$. Intuitively, we aim to prove that the score function is close to $-\frac{\mathbf{x}}{\nu_0^2}$ when $\|\mathbf{n}\|^2 \ge \left(\frac{\nu_0^2}{2} + \frac{\nu_{\max}^2}{2(1-c_{\nu})}\right) d$. Towards this goal, we first show that $P^{(0)}(\mathbf{x})$ is exponentially larger than $P^{(i)}(\mathbf{x})$ for all $i \in [k]$ in the following lemma:

Lemma 10. Suppose P satisfies Assumption 2. Then for any $\|\mathbf{n}\|^2 \ge \left(\frac{\nu_0^2}{2} + \frac{\nu_{\max}^2}{2(1-c_\nu)}\right) d$, we have $\frac{P^{(i)}(\mathbf{x})}{P^{(0)}(\mathbf{x})} \le \exp(-\Omega(d)) \text{ and } \frac{\|\nabla_{\mathbf{x}} P^{(i)}(\mathbf{x})\|}{P(\mathbf{x})} \le \exp(-\Omega(d)) \text{ for all } i \in [k].$

Proof of Lemma 10. We first give an upper bound on the sub-Gaussian probability density. For any vector $\mathbf{v} \in \mathbb{R}^d$, by considering some vector $\mathbf{m} \in \mathbb{R}^d$, from Markov's inequality and the definition in equation 4 we can bound

$$\mathbb{P}_{\mathbf{z}\sim P^{(i)}}\left(\mathbf{m}^{T}(\mathbf{z}-\boldsymbol{\mu}_{i}) \geq \mathbf{m}^{T}(\mathbf{v}-\boldsymbol{\mu}_{i})\right) \leq \frac{\mathbb{E}_{\mathbf{z}\sim P^{(i)}}\left[\exp\left(\mathbf{m}^{T}(\mathbf{z}-\boldsymbol{\mu}_{i})\right)\right]}{\exp\left(\mathbf{m}^{T}(\mathbf{v}-\boldsymbol{\mu}_{i})\right)} \\ \leq \exp\left(\frac{\nu_{i}^{2} \|\mathbf{m}\|^{2}}{2} - \mathbf{m}^{T}(\mathbf{v}-\boldsymbol{\mu}_{i})\right).$$

⁶²³ Upon optimizing the last term at $\mathbf{m} = \frac{\mathbf{v} - \boldsymbol{\mu}_i}{\nu_i^2}$, we obtain

$$\mathbb{P}_{\mathbf{z}\sim P^{(i)}}\left((\mathbf{v}-\boldsymbol{\mu}_i)^T(\mathbf{v}-\mathbf{z})\leq 0\right)\leq \exp\left(-\frac{\|\mathbf{v}-\boldsymbol{\mu}_i\|^2}{2\nu_i^2}\right).$$
(13)

 $\text{ Denote } \mathbb{B} := \big\{ \mathbf{z} : (\mathbf{v} - \boldsymbol{\mu}_i)^T (\mathbf{v} - \mathbf{z}) \leq 0 \big\}. \text{ To bound } \mathbb{P}_{\mathbf{z} \sim P^{(i)}} (\mathbf{z} \in \mathbb{B}), \text{ we first note that }$

$$\log P^{(i)}(\mathbf{v}) - \log P^{(i)}(\mathbf{z})$$

$$= \int_{0}^{1} \langle \mathbf{v} - \mathbf{z}, \nabla \log P^{(i)}(\mathbf{v} + \lambda(\mathbf{z} - \mathbf{v})) \rangle d\lambda$$

$$= \langle \mathbf{v} - \mathbf{z}, \nabla \log P^{(i)}(\mathbf{v}) \rangle + \int_{0}^{1} \langle \mathbf{v} - \mathbf{z}, \nabla \log P^{(i)}(\mathbf{v} + \lambda(\mathbf{z} - \mathbf{v})) - \nabla \log P^{(i)}(\mathbf{v}) \rangle d\lambda$$

$$\leq \|\mathbf{v} - \mathbf{z}\| \left\| \nabla \log P^{(i)}(\mathbf{v}) \right\| + \int_{0}^{1} \|\mathbf{v} - \mathbf{z}\| \left\| \nabla \log P^{(i)}(\mathbf{v} + \lambda(\mathbf{z} - \mathbf{v})) - \nabla \log P^{(i)}(\mathbf{v}) \right\| d\lambda$$

$$\leq \|\mathbf{v} - \mathbf{z}\| \cdot L_{i} \|\mathbf{v} - \boldsymbol{\mu}_{i}\| + \int_{0}^{1} \|\mathbf{v} - \mathbf{z}\| \cdot L_{i} \|\lambda(\mathbf{z} - \mathbf{v})\| d\lambda$$

$$\leq \frac{L_{i}c_{\nu}}{2c_{L}} \|\mathbf{v} - \boldsymbol{\mu}_{i}\|^{2} + \left(\frac{c_{L} + c_{\nu}}{2c_{\nu}}\right) L_{i} \|\mathbf{v} - \mathbf{z}\|^{2},$$
(14)

where equation 14 follows from Assumption 2.ii. that $\nabla \log P^{(i)}(\mu_i) = \mathbf{0}_d$ and Assumption 2.iii. that the score function $\nabla \log P^{(i)}$ is L_i -Lipschitz. Therefore we obtain

$$\mathbb{P}_{\mathbf{z}\sim P^{(i)}}(\mathbf{z}\in\mathbb{B}) = \int_{\mathbf{z}\in\mathbb{B}} P^{(i)}(\mathbf{z}) \,\mathrm{d}\mathbf{z}$$

$$\geq \int_{\mathbf{z}\in\mathbb{B}} P^{(i)}(\mathbf{v}) \exp\left(-\frac{L_i c_{\nu}}{2c_L} \|\mathbf{v}-\boldsymbol{\mu}_i\|^2 - \frac{c_L + c_{\nu}}{2c_{\nu}} L_i \|\mathbf{v}-\mathbf{z}\|^2\right) \,\mathrm{d}\mathbf{z}$$

$$= P^{(i)}(\mathbf{v}) \exp\left(-\frac{L_i c_{\nu}}{2c_L} \|\mathbf{v}-\boldsymbol{\mu}_i\|^2\right) \int_{\mathbf{z}\in\mathbb{B}} \exp\left(-\frac{c_L + c_{\nu}}{2c_{\nu}} L_i \|\mathbf{v}-\mathbf{z}\|^2\right) \,\mathrm{d}\mathbf{z}.$$
 (15)

By observing that $g: \mathbb{B} \to \{\mathbf{z} : (\mathbf{v} - \boldsymbol{\mu}_i)^T (\mathbf{v} - \mathbf{z}) \ge 0\}$ with $g(\mathbf{z}) = 2\mathbf{v} - \mathbf{z}$ is a bijection such that $\|\mathbf{v} - \mathbf{z}\| = \|\mathbf{v} - g(\mathbf{z})\|$ for any $\mathbf{z} \in \mathbb{B}$, we have

$$\int_{\mathbf{z}\in\mathbb{B}} \exp\left(-\frac{c_L + c_\nu}{2c_\nu} L_i \|\mathbf{v} - \mathbf{z}\|^2\right) \, \mathrm{d}\mathbf{z} = \frac{1}{2} \int_{\mathbf{z}\in\mathbb{R}^d} \exp\left(-\frac{c_L + c_\nu}{2c_\nu} L_i \|\mathbf{v} - \mathbf{z}\|^2\right) \, \mathrm{d}\mathbf{z}$$
$$= \frac{1}{2} \left(\frac{2\pi c_\nu}{(c_L + c_\nu)L_i}\right)^{\frac{d}{2}}.$$
(16)

Hence, by combining equation 13, equation 15, and equation 16, we obtain

$$\exp\left(-\frac{\|\mathbf{v}-\boldsymbol{\mu}_i\|^2}{2\nu_i^2}\right) \ge \mathbb{P}_{\mathbf{z}\sim P^{(i)}}\left((\mathbf{v}-\boldsymbol{\mu}_i)^T(\mathbf{v}-\mathbf{z})\le 0\right)$$
$$\ge P^{(i)}(\mathbf{v})\exp\left(-\frac{L_i c_\nu}{2c_L}\|\mathbf{v}-\boldsymbol{\mu}_i\|^2\right) \cdot \frac{1}{2}\left(\frac{2\pi c_\nu}{(c_L+c_\nu)L_i}\right)^{\frac{d}{2}}.$$

By Assumption 2.iii. that $L_i \leq \frac{c_L}{\nu_i^2}$ we obtain the following bound on the probability density:

$$P^{(i)}(\mathbf{v}) \le 2 \left(\frac{2\pi c_{\nu} \nu_i^2}{(c_L + c_{\nu}) c_L} \right)^{-\frac{d}{2}} \exp\left(-\frac{1 - c_{\nu}}{2\nu_i^2} \left\| \mathbf{v} - \boldsymbol{\mu}_i \right\|^2 \right).$$
(17)

Then we can bound the ratio of $P^{(i)}$ and $P^{(0)}$. For all $i \in [k]$, define $\rho_i(\mathbf{x}) := \frac{P^{(i)}(\mathbf{x})}{P^{(0)}(\mathbf{x})}$, then we have

$$\begin{split} \rho_{i}(\mathbf{x}) &= \frac{P^{(i)}(\mathbf{x})}{P^{(0)}(\mathbf{x})} \leq \frac{2(2\pi c_{\nu}\nu_{i}^{2}/(c_{L}^{2}+c_{\nu}c_{L}))^{-d/2}\exp\left(-(1-c_{\nu})\|\mathbf{x}-\boldsymbol{\mu}_{i}\|^{2}/2\nu_{i}^{2}\right)}{(2\pi\nu_{0}^{2})^{-d/2}\exp\left(-\|\mathbf{x}\|^{2}/2\nu_{0}^{2}\right)} \\ &= 2\left(\frac{(c_{L}^{2}+c_{\nu}c_{L})\nu_{0}^{2}}{c_{\nu}\nu_{i}^{2}}\right)^{\frac{d}{2}}\exp\left(\frac{\|\mathbf{x}\|^{2}}{2\nu_{0}^{2}}-\frac{(1-c_{\nu})\|\mathbf{x}-\boldsymbol{\mu}_{i}\|^{2}}{2\nu_{i}^{2}}\right) \\ &= 2\left(\frac{(c_{L}^{2}+c_{\nu}c_{L})\nu_{0}^{2}}{c_{\nu}\nu_{i}^{2}}\right)^{\frac{d}{2}}\exp\left(\left(\frac{1}{2\nu_{0}^{2}}-\frac{1-c_{\nu}}{2\nu_{i}^{2}}\right)\|\mathbf{N}\mathbf{n}\|^{2}+\left(\frac{\|\mathbf{R}\mathbf{r}\|^{2}}{2\nu_{0}^{2}}-\frac{(1-c_{\nu})\|\mathbf{R}\mathbf{r}-\boldsymbol{\mu}_{i}\|^{2}}{2\nu_{i}^{2}}\right)\right) \\ &= 2\left(\frac{(c_{L}^{2}+c_{\nu}c_{L})\nu_{0}^{2}}{c_{\nu}\nu_{i}^{2}}\right)^{\frac{d}{2}}\exp\left(\left(\frac{1}{2\nu_{0}^{2}}-\frac{1-c_{\nu}}{2\nu_{i}^{2}}\right)\|\mathbf{n}\|^{2}+\left(\frac{\|\mathbf{r}\|^{2}}{2\nu_{0}^{2}}-\frac{(1-c_{\nu})\|\mathbf{r}-\mathbf{R}^{T}\boldsymbol{\mu}_{i}\|^{2}}{2\nu_{i}^{2}}\right)\right), \end{split}$$

where the last step follows from the definition that $\mathbf{R} \in \mathbb{R}^{d \times r}$ an orthogonal basis of the vector space $\{\boldsymbol{\mu}_i\}_{i \in [k]} \text{ and } \mathbf{N}^T \mathbf{N} = \boldsymbol{I}_n$. Since $\nu_i^2 < (1 - c_\nu)\nu_0^2$, the quadratic term $\frac{\|\mathbf{r}\|^2}{2\nu_0^2} - \frac{(1 - c_\nu)\|\mathbf{r} - \mathbf{R}^T \boldsymbol{\mu}_i\|^2}{2\nu_i^2}$ is maximized at $\mathbf{r} = \frac{(1 - c_\nu)\nu_0^2 \mathbf{R}^T \boldsymbol{\mu}_i}{(1 - c_\nu)\nu_0^2 - \nu_i^2}$. Therefore, we obtain

$$\frac{\|\mathbf{r}\|^2}{2\nu_0^2} - \frac{(1-c_\nu) \|\mathbf{r} - \mathbf{R}^T \boldsymbol{\mu}_i\|^2}{2\nu_i^2} \le \frac{(1-c_\nu) \|\boldsymbol{\mu}_i\|^2}{2((1-c_\nu)\nu_0^2 - \nu_i^2)}$$

635 Hence, for $\|\boldsymbol{\mu}_{i} - \boldsymbol{\mu}_{0}\|^{2} \leq \frac{(1-c_{\nu})\nu_{0}^{2}-\nu_{i}^{2}}{2(1-c_{\nu})} \left(\log \frac{c_{\nu}\nu_{i}^{2}}{(c_{L}^{2}+c_{\nu}c_{L})\nu_{0}^{2}} - \frac{\nu_{i}^{2}}{2(1-c_{\nu})\nu_{0}^{2}} + \frac{(1-c_{\nu})\nu_{0}^{2}}{2\nu_{i}^{2}} \right) d$ and $\|\mathbf{n}\|^{2} \geq \frac{(1-c_{\nu})\nu_{0}^{2}}{2} + \frac{\nu_{\max}^{2}}{2(1-c_{\nu})} d$, we have

$$\begin{split} \rho_{i}(\mathbf{x}) &\leq 2 \left(\frac{(c_{L}^{2} + c_{\nu}c_{L})\nu_{0}^{2}}{c_{\nu}\nu_{i}^{2}} \right)^{\frac{d}{2}} \exp\left(\left(\frac{1}{2\nu_{0}^{2}} - \frac{1 - c_{\nu}}{2\nu_{i}^{2}} \right) \|\mathbf{n}\|^{2} + \frac{(1 - c_{\nu})\|\boldsymbol{\mu}_{i}\|^{2}}{2((1 - c_{\nu})\nu_{0}^{2} - \nu_{i}^{2})} \right) \\ &\leq 2 \left(\frac{(c_{L}^{2} + c_{\nu}c_{L})\nu_{0}^{2}}{c_{\nu}\nu_{i}^{2}} \right)^{\frac{d}{2}} \exp\left(\left(\frac{1}{2\nu_{0}^{2}} - \frac{1 - c_{\nu}}{2\nu_{i}^{2}} \right) \left(\frac{\nu_{0}^{2}}{2} + \frac{\nu_{i}^{2}}{2(1 - c_{\nu})} \right) d + \frac{(1 - c_{\nu})\|\boldsymbol{\mu}_{i}\|^{2}}{2((1 - c_{\nu})\nu_{0}^{2} - \nu_{i}^{2})} \right) \\ &= 2 \exp\left(- \left(\log \frac{c_{\nu}\nu_{i}^{2}}{(c_{L}^{2} + c_{\nu}c_{L})\nu_{0}^{2}} - \frac{\nu_{i}^{2}}{2(1 - c_{\nu})\nu_{0}^{2}} + \frac{(1 - c_{\nu})\nu_{0}^{2}}{2\nu_{i}^{2}} \right) \frac{d}{2} + \frac{(1 - c_{\nu})\|\boldsymbol{\mu}_{i}\|^{2}}{2((1 - c_{\nu})\nu_{0}^{2} - \nu_{i}^{2})} \right) \\ &\leq 2 \exp\left(- \left(\log \frac{c_{\nu}\nu_{i}^{2}}{(c_{L}^{2} + c_{\nu}c_{L})\nu_{0}^{2}} - \frac{\nu_{i}^{2}}{2(1 - c_{\nu})\nu_{0}^{2}} + \frac{(1 - c_{\nu})\nu_{0}^{2}}{2\nu_{i}^{2}} \right) \frac{d}{4} \right). \end{split}$$

From Lemma 9, we obtain $\rho_i(\mathbf{x}) \leq \exp(-\Omega(d))$.

To show $\frac{\|\nabla_{\mathbf{x}} P^{(i)}(\mathbf{x})\|}{P(\mathbf{x})} \leq \exp(-\Omega(d))$, from Assumptions 2.ii. and 2.iii. we have

$$\begin{aligned} \left\| \frac{\nabla_{\mathbf{x}} P^{(i)}(\mathbf{x})}{P^{(i)}(\mathbf{x})} \right\| &= \left\| \frac{\nabla_{\mathbf{x}} P^{(i)}(\mathbf{x})}{P^{(i)}(\mathbf{x})} - \frac{\nabla_{\mathbf{x}} P^{(i)}(\boldsymbol{\mu}_i)}{P^{(i)}(\boldsymbol{\mu}_i)} \right\| = \left\| \nabla_{\mathbf{x}} \log P^{(i)}(\mathbf{x}) - \nabla_{\mathbf{x}} \log P^{(i)}(\boldsymbol{\mu}_i) \right\| \\ &\leq L_i \left\| \mathbf{x} - \boldsymbol{\mu}_i \right\| \leq \frac{c_L}{\nu_i^2} \left\| \mathbf{x} - \boldsymbol{\mu}_i \right\|. \end{aligned}$$

639 Therefore, we can bound $\frac{\|\nabla_{\mathbf{x}} P^{(i)}(\mathbf{x})\|}{P(\mathbf{x})} \leq \frac{c_L}{\nu_i^2} \rho_i(\mathbf{x}) \|\mathbf{x} - \boldsymbol{\mu}_i\|$. When $\|\mathbf{x} - \boldsymbol{\mu}_i\| = \exp(o(d))$ is 640 small, by $\rho_i(\mathbf{x}) \leq \exp(-\Omega(d))$ we directly have $\frac{\|\nabla_{\mathbf{x}} P^{(i)}(\mathbf{x})\|}{P(\mathbf{x})} \leq \exp(-\Omega(d))$. When $\|\mathbf{x} - \boldsymbol{\mu}_i\| = \exp(\Omega(d))$ is exceedingly large, from equation 17 we have

$$\frac{\left\|\nabla_{\mathbf{x}}P^{(i)}(\mathbf{x})\right\|}{P(\mathbf{x})} \leq \frac{2c_L}{\nu_i^2} \left(\frac{(c_L^2 + c_\nu c_L)\nu_0^2}{c_\nu \nu_i^2}\right)^{\frac{d}{2}} \exp\left(\frac{\left\|\mathbf{x}\right\|^2}{2\nu_0^2} - \frac{(1 - c_\nu)\left\|\mathbf{x} - \boldsymbol{\mu}_i\right\|^2}{2\nu_i^2}\right) \left\|\mathbf{x} - \boldsymbol{\mu}_i\right\|.$$

642 Since $\nu_0^2 > \frac{\nu_i^2}{1-c_\nu}$, when $\|\mathbf{x} - \boldsymbol{\mu}_i\| = \exp(\Omega(d)) \gg \|\boldsymbol{\mu}_i\|$ we have

$$\exp\left(\frac{\|\mathbf{x}\|^{2}}{2\nu_{0}^{2}} - \frac{(1-c_{\nu})\|\mathbf{x}-\boldsymbol{\mu}_{i}\|^{2}}{2\nu_{i}^{2}}\right) = \exp(-\Omega(\|\mathbf{x}-\boldsymbol{\mu}_{i}\|^{2})).$$

643 Therefore $\frac{\|\nabla_{\mathbf{x}} P^{(i)}(\mathbf{x})\|}{P(\mathbf{x})} \le \exp(-\Omega(d))$. Thus we complete the proof of Lemma 10.

Similar to Lemma 5, the following lemma proves that when the previous state n_{t-1} is far from a mode, a single step of Langevin dynamics with bounded step size is not enough to find the mode.

Lemma 11. Suppose $\delta_t \leq \nu_0^2$ and $\|\mathbf{n}_{t-1}\|^2 > 36\nu_0^2 d$, then we have $\|\mathbf{n}_t\|^2 \geq \nu_0^2 d$ with probability at *least* $1 - \exp(-\Omega(d))$.

Proof of Lemma 11. For simplicity, denote $\mathbf{v} := \mathbf{n}_{t-1} + \frac{\delta_t}{2} \mathbf{N}^T \nabla_{\mathbf{x}} \log P(\mathbf{x}_{t-1})$. Since $P = \sum_{i=0}^k w_i P^{(i)}$ and $P^{(0)} = \mathcal{N}(\boldsymbol{\mu}_0, \nu_0^2 \boldsymbol{I}_d)$, the score function can be written as

$$\nabla_{\mathbf{x}} \log P(\mathbf{x}) = \frac{\nabla_{\mathbf{x}} P(\mathbf{x})}{P(\mathbf{x})} = \frac{\nabla_{\mathbf{x}} w_0 P^{(0)}(\mathbf{x})}{P(\mathbf{x})} + \sum_{i \in [k]} \frac{\nabla_{\mathbf{x}} w_i P^{(i)}(\mathbf{x})}{P(\mathbf{x})}$$
$$= -\frac{w_0 P^{(0)}(\mathbf{x})}{P(\mathbf{x})} \cdot \frac{\mathbf{x}}{\nu_0^2} + \sum_{i \in [k]} \frac{w_i \nabla_{\mathbf{x}} P^{(i)}(\mathbf{x})}{P(\mathbf{x})}$$
$$= -\frac{\mathbf{x}}{\nu_0^2} + \sum_{i \in [k]} \frac{w_i P^{(i)}(\mathbf{x})}{P(\mathbf{x})} \cdot \frac{\mathbf{x}}{\nu_0^2} + \sum_{i \in [k]} \frac{w_i \nabla_{\mathbf{x}} P^{(i)}(\mathbf{x})}{P(\mathbf{x})}.$$
(18)

For $\|\mathbf{n}_{t-1}\|^2 > 36\nu_0^2 d$ by Lemma 10 we have $\frac{\|\nabla_{\mathbf{x}} P^{(i)}(\mathbf{x}_{t-1})\|}{P(\mathbf{x}_{t-1})} \le \exp(-\Omega(d))$. Since $\delta_t \le \nu_0^2$, we can bound the norm of \mathbf{v} by

$$\begin{aligned} \|\mathbf{v}\| &= \left\| \mathbf{n}_{t-1} + \frac{\delta_{t}}{2} \mathbf{N}^{T} \nabla_{\mathbf{x}} \log P(\mathbf{x}_{t-1}) \right\| \\ &= \left\| \mathbf{n}_{t-1} - \frac{\delta_{t}}{2\nu_{0}^{2}} \mathbf{n}_{t-1} + \sum_{i \in [k]} \frac{w_{i} \delta_{t}}{2\nu_{0}^{2}} \frac{P^{(i)}(\mathbf{x}_{t-1})}{P(\mathbf{x}_{t-1})} \mathbf{n}_{t-1} + \sum_{i \in [k]} \frac{w_{i} \delta_{t}}{2} \frac{\mathbf{N}^{T} \nabla_{\mathbf{x}} P^{(i)}(\mathbf{x}_{t-1})}{P(\mathbf{x}_{t-1})} \right\| \\ &\geq \left\| \left(1 - \frac{\delta_{t}}{2\nu_{0}^{2}} + \sum_{i \in [k]} \frac{w_{i} \delta_{t}}{2\nu_{0}^{2}} \frac{P^{(i)}(\mathbf{x}_{t-1})}{P(\mathbf{x}_{t-1})} \right) \mathbf{n}_{t-1} \right\| - \sum_{i \in [k]} \frac{w_{i} \delta_{t}}{2} \frac{\left\| \nabla_{\mathbf{x}} P^{(i)}(\mathbf{x}_{t-1}) \right\|}{P(\mathbf{x}_{t-1})} \\ &\geq \frac{1}{2} \left\| \mathbf{n}_{t-1} \right\| - \sum_{i \in [k]} \frac{w_{i} \delta_{t}}{2} \exp(-\Omega(d)) \\ &> 2\nu_{0} \sqrt{d}. \end{aligned}$$

652 On the other hand, from $\boldsymbol{\epsilon}_t^{(\mathbf{n})} \sim \mathcal{N}(\mathbf{0}_n, \boldsymbol{I}_n)$ we know $\frac{\langle \mathbf{v}, \boldsymbol{\epsilon}_t^{(\mathbf{n})} \rangle}{\|\mathbf{v}\|} \sim \mathcal{N}(0, 1)$ for any fixed $\mathbf{v} \neq \mathbf{0}_n$, hence 653 by Lemma 2 we have

$$\mathbb{P}\left(\frac{\langle \mathbf{v}, \boldsymbol{\epsilon}_t^{(\mathbf{n})} \rangle}{\|\mathbf{v}\|} \ge \frac{\sqrt{d}}{4}\right) = \mathbb{P}\left(\frac{\langle \mathbf{v}, \boldsymbol{\epsilon}_t^{(\mathbf{n})} \rangle}{\|\mathbf{v}\|} \le -\frac{\sqrt{d}}{4}\right) \le \frac{4}{\sqrt{2\pi d}} \exp\left(-\frac{d}{32}\right)$$

654 Combining the above inequalities gives

$$\|\mathbf{n}_t\|^2 = \left\|\mathbf{v} + \sqrt{\delta_t} \boldsymbol{\epsilon}_t^{(\mathbf{n})}\right\|^2 \ge \|\mathbf{v}\|^2 - 2\nu_0 |\langle \mathbf{v}, \boldsymbol{\epsilon}_t^{(\mathbf{n})} \rangle| \ge \|\mathbf{v}\|^2 - \frac{\nu_0 \sqrt{d}}{2} \|\mathbf{v}\| > \nu_0^2 d$$

with probability at least $1 - \frac{8}{\sqrt{2\pi d}} \exp\left(-\frac{d}{32}\right) = 1 - \exp(-\Omega(d))$. This proves Lemma 11.

656 When $\|\mathbf{n}_{t-1}\|^2 \leq 36\nu_0^2 d$, similar to Theorem 1, we consider a surrogate recursion $\hat{\mathbf{n}}_t$ such that 657 $\hat{\mathbf{n}}_0 = \mathbf{n}_0$ and for all $t \geq 1$,

$$\hat{\mathbf{n}}_t = \hat{\mathbf{n}}_{t-1} - \frac{\delta_t}{2\nu_0^2} \hat{\mathbf{n}}_{t-1} + \sqrt{\delta_t} \boldsymbol{\epsilon}_t^{(\mathbf{n})}.$$
(19)

The following Lemma shows that $\hat{\mathbf{n}}_t$ is sufficiently close to the original recursion \mathbf{n}_t .

Lemma 12. For any $t \ge 1$, given that for all $j \in [t]$, $\delta_j \le \nu_0^2$ and $\left(\frac{\nu_0^2}{2} + \frac{\nu_{\max}^2}{2(1-c_\nu)}\right) d \le \|\mathbf{n}_{j-1}\|^2 \le 36\nu_0^2 d$, if $\boldsymbol{\mu}_i$ satisfies Assumption 2.v. for all $i \in [k]$, we have $\|\hat{\mathbf{n}}_t - \mathbf{n}_t\| \le \frac{t}{\exp(\Omega(d))}\sqrt{d}$.

⁶⁶¹ *Proof of Lemma 12.* By equation 18 we have that for all $j \in [t]$,

$$\begin{aligned} \|\hat{\mathbf{n}}_{j} - \mathbf{n}_{j}\| &= \left\|\hat{\mathbf{n}}_{j-1} - \mathbf{n}_{j-1} - \frac{\delta_{j}}{2\nu_{0}^{2}}\hat{\mathbf{n}}_{j-1} - \frac{\delta_{j}}{2}\mathbf{N}^{T}\nabla_{\mathbf{x}}\log P(\mathbf{x}_{j-1})\right\| \\ &= \left\|\hat{\mathbf{n}}_{j-1} - \mathbf{n}_{j-1} - \sum_{i\in[k]}\frac{w_{i}P^{(i)}(\mathbf{x}_{j-1})}{\nu_{0}^{2}P(\mathbf{x}_{j-1})}\mathbf{n}_{j-1} - \sum_{i\in[k]}\frac{w_{i}\mathbf{N}^{T}\nabla_{\mathbf{x}}P^{(i)}(\mathbf{x}_{j-1})}{P(\mathbf{x}_{j-1})}\right\| \\ &\leq \|\hat{\mathbf{n}}_{j-1} - \mathbf{n}_{j-1}\| + \sum_{i\in[k]}\frac{w_{i}P^{(i)}(\mathbf{x}_{j-1})}{\nu_{0}^{2}P(\mathbf{x}_{j-1})}\|\mathbf{n}_{j-1}\| + \sum_{i\in[k]}\frac{w_{i}\|\nabla_{\mathbf{x}}P^{(i)}(\mathbf{x}_{j-1})\|}{P(\mathbf{x}_{j-1})}.\end{aligned}$$

By Lemma 10, we have $\frac{P^{(i)}(\mathbf{x}_{j-1})}{P^{(0)}(\mathbf{x}_{j-1})} \leq \exp(-\Omega(d))$ and $\frac{\|\nabla_{\mathbf{x}}P^{(i)}(\mathbf{x}_{j-1})\|}{P(\mathbf{x}_{j-1})} \leq \exp(-\Omega(d))$ for all i $\in [k]$, hence from $\|\mathbf{n}_{j-1}\| \leq 6\nu_0\sqrt{d}$ we obtain a recursive bound

$$\|\hat{\mathbf{n}}_j - \mathbf{n}_j\| \le \|\hat{\mathbf{n}}_{j-1} - \mathbf{n}_{j-1}\| + \frac{1}{\exp(\Omega(d))}\sqrt{d}.$$

664 Finally, by $\hat{\mathbf{n}}_0 = \mathbf{n}_0$, we have

$$\|\hat{\mathbf{n}}_t - \mathbf{n}_t\| = \sum_{j \in [t]} (\|\hat{\mathbf{n}}_j - \mathbf{n}_j\| - \|\hat{\mathbf{n}}_{j-1} - \mathbf{n}_{j-1}\|) \le \frac{t}{\exp(\Omega(d))} \sqrt{d}.$$

665 Hence we obtain Lemma 12.

Armed with the above lemmas, we are now ready to establish Proposition 5 by induction. Please note that we also apply some lemmas from the proof of Theorem 1 by substituting ν_{\max}^2 with $\frac{\nu_{\max}^2}{1-c_{\nu}}$. Suppose the theorem holds for all T values of $1, \dots, T-1$. We consider the following 3 cases:

• If there exists some $t \in [T]$ such that $\delta_t > \nu_0^2$, by Lemma 3 we know that with probability at least $1 - \exp(-\Omega(d))$, we have $\|\mathbf{n}_t\|^2 \ge \left(\frac{3\nu_0^2}{4} + \frac{\nu_{\max}^2}{4(1-c_\nu)}\right) d$, thus the problem reduces to the two sub-arrays $\mathbf{n}_0, \cdots, \mathbf{n}_{t-1}$ and $\mathbf{n}_t, \cdots, \mathbf{n}_T$, which can be solved by induction.

• Suppose $\delta_t \leq \nu_0^2$ for all $t \in [T]$. If there exists some $t \in [T]$ such that $\|\mathbf{n}_{t-1}\|^2 > 36\nu_0^2 d$, by Lemma 11 we know that with probability at least $1 - \exp(-\Omega(d))$, we have $\|\mathbf{n}_t\|^2 \geq \nu_0^2 d > \frac{3\nu_0^2}{4} + \frac{\nu_{\max}^2}{4(1-c_{\nu})} d$, thus the problem similarly reduces to the two sub-arrays $\mathbf{n}_0, \cdots, \mathbf{n}_{t-1}$ and $\mathbf{n}_t, \cdots, \mathbf{n}_T$, which can be solved by induction.

• Suppose
$$\delta_t \leq \nu_0^2$$
 and $\|\mathbf{n}_{t-1}\|^2 \leq 36\nu_0^2 d$ for all $t \in [T]$. Conditioned on $\|\mathbf{n}_{t-1}\|^2 > \begin{pmatrix} \frac{\nu_0^2}{2} + \frac{\nu_{\max}^2}{2(1-c_\nu)} \end{pmatrix} d$ for all $t \in [T]$, by Lemma 12 we have that for $T = \exp(\mathcal{O}(d))$,

$$\|\hat{\mathbf{n}}_T - \mathbf{n}_T\| < \left(\sqrt{\frac{5\nu_0^2}{8} + \frac{3\nu_{\max}^2}{8(1 - c_\nu)}} - \sqrt{\frac{\nu_0^2}{2} + \frac{\nu_{\max}^2}{2(1 - c_\nu)}}\right)\sqrt{d}$$

678

By Lemma 8 we have that with probability at least $1 - \exp(-\Omega(d))$,

$$\|\hat{\mathbf{n}}_T\|^2 \ge \left(\frac{5\nu_0^2}{8} + \frac{3\nu_{\max}^2}{8(1-c_\nu)}\right) d.$$

Combining the two inequalities implies the desired bound

$$\|\mathbf{n}_T\| \ge \|\hat{\mathbf{n}}_T\| - \|\hat{\mathbf{n}}_T - \mathbf{n}_T\| > \sqrt{\left(\frac{\nu_0^2}{2} + \frac{\nu_{\max}^2}{2(1 - c_{\nu})}\right)d}.$$

Hence by induction we obtain $\|\mathbf{n}_t\|^2 > \left(\frac{\nu_0^2}{2} + \frac{\nu_{\max}^2}{2(1-c_\nu)}\right) d$ for all $t \in [T]$ with probability at least

$$(1 - (T - 1)\exp(-\Omega(d))) \cdot (1 - \exp(-\Omega(d))) \ge 1 - T\exp(-\Omega(d)).$$

⁶⁸² Therefore we complete the proof of Proposition 5.

⁶⁸³ Finally, combining equation 12 and Proposition 5 finishes the proof of Theorem 3.

684 A.4 Proof of Theorem 4: Annealed Langevin Dynamics under Sub-Gaussian Mixtures

Assumption 3. Consider a data distribution $P := \sum_{i=0}^{k} w_i P^{(i)}$ as a mixture of sub-Gaussian distributions, where $1 \le k = o(d)$ and $w_i > 0$ is a positive constant such that $\sum_{i=0}^{k} w_i = 1$. Suppose that $P^{(0)} = \mathcal{N}(\boldsymbol{\mu}_0, \nu_0^2 \boldsymbol{I}_d)$ is Gaussian and for all $i \in [k]$, $P^{(i)}$ satisfies

- 688 *i.* $P^{(i)}$ is a sub-Gaussian distribution of mean μ_i with parameter ν_i^2 ,
- 689 *ii.* $P^{(i)}$ is differentiable and $\nabla P^{(i)}_{\sigma_t}(\boldsymbol{\mu}_i) = \mathbf{0}_d$ for all $t \in \{0\} \cup [T]$,
- 690 *iii.* for all $t \in \{0\} \cup [T]$, the score function of $P_{\sigma_t}^{(i)}$ is $L_{i,t}$ -Lipschitz such that $L_{i,t} \leq \frac{c_L}{\nu_i^2 + \sigma_t^2}$ for 691 some constant $c_L > 0$,

692 *iv.*
$$\nu_0^2 > \max\left\{1, \frac{4(c_L^2 + c_\nu c_L)}{c_\nu(1 - c_\nu)}\right\} \frac{\nu_{\max}^2 + c_\sigma^2}{1 - c_\nu} - c_\sigma^2 \text{ for constant } c_\nu \in (0, 1), \text{ where } \nu_{\max} := \max_{i \in [k]} \nu_i,$$

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$$v. \|\boldsymbol{\mu}_{i} - \boldsymbol{\mu}_{0}\|^{2} \leq \frac{(1-c_{\nu})\nu_{0}^{2} - \nu_{i}^{2} - c_{\nu}c_{\sigma}^{2}}{2(1-c_{\nu})} \left(\log \frac{c_{\nu}(\nu_{i}^{2} + c_{\sigma}^{2})}{(c_{L}^{2} + c_{\nu}c_{L})(\nu_{0}^{2} + c_{\sigma}^{2})} - \frac{(\nu_{i}^{2} + c_{\sigma}^{2})}{2(1-c_{\nu})(\nu_{0}^{2} + c_{\sigma}^{2})} + \frac{(1-c_{\nu})(\nu_{0}^{2} + c_{\sigma}^{2})}{2(\nu_{i}^{2} + c_{\sigma}^{2})} \right) d.$$

The feasibility of Assumption 3.v. can be validated by substituting ν^2 in Lemma 9 with $\nu^2 + c_{\sigma}^2$. To establish Theorem 4, we first note from Proposition 1 that for a sub-Gaussian mixture $P = \sum_{i=0}^{k} w_i P^{(i)}$, the perturbed distribution of noise level σ is $P_{\sigma} = \sum_{i=0}^{k} w_i P^{(i)}_{\sigma}$, where $P^{(0)} = \mathcal{N}(\mu_0, (\nu_i^2 + \sigma^2)\mathbf{I}_d)$ and $P^{(i)}$ is a sub-Gaussian distribution with mean μ_i and sub-Gaussian parameter $(\nu_i^2 + \sigma^2)$. Similar to the proof of Theorem 1, we decompose

$$\mathbf{x}_t = \mathbf{R}\mathbf{r}_t + \mathbf{N}\mathbf{n}_t$$
, and $\boldsymbol{\epsilon}_t = \mathbf{R}\boldsymbol{\epsilon}_t^{(\mathbf{r})} + \mathbf{N}\boldsymbol{\epsilon}_t^{(\mathbf{n})}$,

where $\mathbf{R} \in \mathbb{R}^{d \times r}$ an orthonormal basis of the vector space $\{\boldsymbol{\mu}_i\}_{i \in [k]}$ and $\mathbf{N} \in \mathbb{R}^{d \times n}$ an orthonormal basis of the null space of $\{\boldsymbol{\mu}_i\}_{i \in [k]}$. Now, we prove Theorem 4 by applying the techniques developed in Appendix A.1 and A.3 via substituting ν^2 and $\frac{\nu^2}{1-c_{\nu}}$ with $\frac{\nu^2+\sigma_t^2}{1-c_{\nu}}$ at time step t. Note that for all $t \in \{0\} \cup [T]$, Assumption 3.iv. implies $\nu_0^2 + \sigma_t^2 > \max\left\{1, \frac{4(c_L^2+c_{\nu}c_L)}{c_{\nu}(1-c_{\nu})}\right\} \frac{\nu_{\max}^2 + \sigma_t^2}{1-c_{\nu}}$ because $c_{\sigma} \ge \sigma_t$.

First, by Proposition 2, suppose that the sample is initialized in the distribution $P_{\sigma_0}^{(0)}$, then with probability at least $1 - \exp(-\Omega(d))$, we have

$$\|\mathbf{n}_0\|^2 \ge \left(\frac{3(\nu_0^2 + \sigma_0^2)}{4} + \frac{\nu_{\max}^2 + \sigma_0^2}{4(1 - c_\nu)}\right) d.$$
⁽²⁰⁾

Then, with the assumption that the initialization satisfies $\|\mathbf{n}_0\|^2 \ge \left(\frac{3(\nu_0^2 + \sigma_0^2)}{4} + \frac{\nu_{\max}^2 + \sigma_0^2}{4(1 - c_\nu)}\right) d$, the following proposition similar to Proposition 5 shows that $\|\mathbf{n}_t\|$ remains large with high probability.

Proposition 6. Consider a distribution P satisfying Assumption 3. We follow annealed Langevin dynamics for $T = \exp(\mathcal{O}(d))$ steps with noise level $c_{\sigma} \ge \sigma_0 \ge \sigma_1 \ge \cdots \ge \sigma_T \ge 0$ for some constant $c_{\sigma} > 0$. Suppose that the initial sample satisfies $\|\mathbf{n}_0\|^2 \ge \left(\frac{3(\nu_0^2 + \sigma_0^2)}{4} + \frac{\nu_{\max}^2 + \sigma_0^2}{4(1-c_{\nu})}\right) d$, then with probability at least $1 - T \cdot \exp(-\Omega(d))$, we have that $\|\mathbf{n}_t\|^2 > \left(\frac{\nu_0^2 + \sigma_t^2}{2} + \frac{\nu_{\max}^2 + \sigma_t^2}{2(1-c_{\nu})}\right) d$ for all $t \in \{0\} \cup [T]$.

Proof of Proposition 6. We prove Proposition 6 by induction. Suppose the theorem holds for all T713 values of $1, \dots, T-1$. We consider the following 3 cases: 714

- If there exists some $t \in [T]$ such that $\delta_t > \nu_0^2 + \sigma_t^2$, by Lemma 3 we know that with 715 probability at least $1 - \exp(-\Omega(d))$, we have $\|\mathbf{n}_t\|^2 \ge \left(\frac{3(\nu_0^2 + \sigma_t^2)}{4} + \frac{\nu_{\max}^2 + \sigma_t^2}{4(1 - c_\nu)}\right) d$, thus the problem reduces to the two sub-arrays $\mathbf{n}_0, \dots, \mathbf{n}_{t-1}$ and $\mathbf{n}_t, \dots, \mathbf{n}_T$, which can be solved 716 717 718 by induction.
- Suppose $\delta_t \leq \nu_0^2 + \sigma_t^2$ for all $t \in [T]$. If there exists some $t \in [T]$ such that $\|\mathbf{n}_{t-1}\|^2 > 36(\nu_0^2 + \sigma_{t-1}^2)d \geq 36(\nu_0^2 + \sigma_t^2)d$, by Lemma 11 we know that with probability at least $1 \exp(-\Omega(d))$, we have $\|\mathbf{n}_t\|^2 \geq (\nu_0^2 + \sigma_t^2)d > \left(\frac{3(\nu_0^2 + \sigma_t^2)}{4} + \frac{\nu_{\max}^2 + \sigma_t^2}{4(1-c_\nu)}\right)d$, thus the problem similarly reduces to the two sub-arrays $\mathbf{n}_0, \cdots, \mathbf{n}_{t-1}$ and $\mathbf{n}_t, \cdots, \mathbf{n}_T$, which can be sub-advised by the sub-array \mathbf{n}_t and $\mathbf{n}_t, \cdots, \mathbf{n}_T$. 719 720 721 722
- be solved by induction 723
- Suppose $\delta_t \leq \nu_0^2 + \sigma_t^2$ and $\|\mathbf{n}_{t-1}\|^2 \leq 36(\nu_0^2 + \sigma_{t-1}^2)d$ for all $t \in [T]$. Consider a surrogate sequence $\hat{\mathbf{n}}_t$ such that $\hat{\mathbf{n}}_0 = \mathbf{n}_0$ and for all $t \geq 1$, 724 725

$$\hat{\mathbf{n}}_t = \hat{\mathbf{n}}_{t-1} - \frac{\delta_t}{2\nu_0^2 + 2\sigma_t^2} \hat{\mathbf{n}}_{t-1} + \sqrt{\delta_t} \boldsymbol{\epsilon}_t^{(\mathbf{n})}$$

Since $\nu_0 > \nu_i$ and $c_\sigma \ge \sigma_t$ for all $t \in \{0\} \cup [T]$, we have $\frac{\nu_i^2 + c_\sigma^2}{\nu_0^2 + c_\sigma^2} > \frac{\nu_i^2 + \sigma_t^2}{\nu_0^2 + \sigma_t^2}$. Notice that for 726 function $f(z) = \log z - \frac{z}{2} + \frac{1}{2z}$, we have $\frac{d}{dz}f(z) = \frac{1}{z} - \frac{1}{2} - \frac{1}{2z^2} = -\frac{1}{2}\left(\frac{1}{z} - 1\right)^2 \le 0$. 727

Thus, by Assumption 3.v. we have that for all
$$t \in [T]$$
,

$$\begin{split} \|\boldsymbol{\mu}_{i} - \boldsymbol{\mu}_{0}\|^{2} &\leq \frac{(1 - c_{\nu})\nu_{0}^{2} - \nu_{i}^{2} - c_{\nu}c_{\sigma}^{2}}{2(1 - c_{\nu})} \left(\log \frac{c_{\nu}(\nu_{i}^{2} + c_{\sigma}^{2})}{(c_{L}^{2} + c_{\nu}c_{L})(\nu_{0}^{2} + c_{\sigma}^{2})} - \frac{(\nu_{i}^{2} + c_{\sigma}^{2})}{2(1 - c_{\nu})(\nu_{0}^{2} + c_{\sigma}^{2})} + \frac{(1 - c_{\nu})(\nu_{0}^{2} + c_{\sigma}^{2})}{2(\nu_{i}^{2} + c_{\sigma}^{2})} \right) d \\ &\leq \frac{(1 - c_{\nu})\nu_{0}^{2} - \nu_{i}^{2} - c_{\nu}\sigma_{t}^{2}}{2(1 - c_{\nu})} \left(\log \frac{c_{\nu}(\nu_{i}^{2} + \sigma_{t}^{2})}{(c_{L}^{2} + c_{\nu}c_{L})(\nu_{0}^{2} + \sigma_{t}^{2})} - \frac{(\nu_{i}^{2} + \sigma_{t}^{2})}{2(1 - c_{\nu})(\nu_{0}^{2} + \sigma_{t}^{2})} + \frac{(1 - c_{\nu})(\nu_{0}^{2} + \sigma_{t}^{2})}{2(\nu_{i}^{2} + \sigma_{t}^{2})} \right) d \end{split}$$

Conditioned on $\|\mathbf{n}_{t-1}\|^2 > \left(\frac{\nu_0^2 + \sigma_{t-1}^2}{2} + \frac{\nu_{\max}^2 + \sigma_{t-1}^2}{2(1-c_\nu)}\right) d$ for all $t \in [T]$, by Lemma 12 we 729 have that for $T = \exp(\mathcal{O}(d))$, 730

$$\|\hat{\mathbf{n}}_T - \mathbf{n}_T\| < \left(\sqrt{\frac{5(\nu_0^2 + \sigma_T^2)}{8} + \frac{3(\nu_{\max}^2 + \sigma_T^2)}{8(1 - c_\nu)}} - \sqrt{\frac{\nu_0^2 + \sigma_T^2}{2} + \frac{\nu_{\max}^2 + \sigma_T^2}{2(1 - c_\nu)}}\right)\sqrt{d}.$$

By Lemma 8 we have that with probability at least $1 - \exp(-\Omega(d))$, 731

$$\|\hat{\mathbf{n}}_T\|^2 \ge \left(\frac{5(\nu_0^2 + \sigma_T^2)}{8} + \frac{3(\nu_{\max}^2 + \sigma_T^2)}{8(1 - c_\nu)}\right) d$$

Combining the two inequalities implies the desired bound 732

$$\|\mathbf{n}_T\| \ge \|\hat{\mathbf{n}}_T\| - \|\hat{\mathbf{n}}_T - \mathbf{n}_T\| > \sqrt{\left(\frac{\nu_0^2 + \sigma_T^2}{2} + \frac{\nu_{\max}^2 + \sigma_T^2}{2(1 - c_\nu)}\right)d}$$

Hence by induction we obtain $\|\mathbf{n}_t\|^2 > \left(\frac{\nu_0^2 + \sigma_T^2}{2} + \frac{\nu_{\max}^2 + \sigma_T^2}{2(1 - c_\nu)}\right) d$ for all $t \in [T]$ with proba-733 bility at least 734

$$(1 - (T - 1)\exp(-\Omega(d))) \cdot (1 - \exp(-\Omega(d))) \ge 1 - T\exp(-\Omega(d)).$$

Therefore we complete the proof of Proposition 6. 735

Finally, combining equation 20 and Proposition 6 finishes the proof of Theorem 4. 736

737 A.5 Proof of Theorem 5: Convergence Analysis of Chained Langevin Dynamics

For simplicity, denote $\mathbf{x}^{[q]} = \{\mathbf{x}^{(1)}, \cdots, \mathbf{x}^{(q)}\}$. By the definition of total variation distance, for all $q \in [d/Q]$ we have

$$\begin{aligned} & \operatorname{TV}\left(\hat{P}\left(\mathbf{x}^{[q]}\right), P\left(\mathbf{x}^{[q]}\right)\right) \\ &= \frac{1}{2} \int \left|\hat{P}\left(\mathbf{x}^{[q]}\right) - P\left(\mathbf{x}^{[q]}\right)\right| \, \mathrm{d}\mathbf{x}^{[q]} \\ &= \frac{1}{2} \int \left|\hat{P}\left(\mathbf{x}^{(q)} \mid \mathbf{x}^{[q-1]}\right) \hat{P}\left(\mathbf{x}^{[q-1]}\right) - P\left(\mathbf{x}^{(q)} \mid \mathbf{x}^{[q-1]}\right) P\left(\mathbf{x}^{[q-1]}\right)\right| \, \mathrm{d}\mathbf{x}^{[q]} \\ &\leq \frac{1}{2} \int \left|\hat{P}\left(\mathbf{x}^{(q)} \mid \mathbf{x}^{[q-1]}\right) \hat{P}\left(\mathbf{x}^{[q-1]}\right) - \hat{P}\left(\mathbf{x}^{(q)} \mid \mathbf{x}^{[q-1]}\right) P\left(\mathbf{x}^{[q-1]}\right)\right| \, \mathrm{d}\mathbf{x}^{[q]} \\ &\quad + \frac{1}{2} \int \left|\hat{P}\left(\mathbf{x}^{(q)} \mid \mathbf{x}^{[q-1]}\right) P\left(\mathbf{x}^{[q-1]}\right) - P\left(\mathbf{x}^{(q)} \mid \mathbf{x}^{[q-1]}\right) P\left(\mathbf{x}^{[q-1]}\right)\right| \, \mathrm{d}\mathbf{x}^{[q]} \\ &= \frac{1}{2} \int \hat{P}\left(\mathbf{x}^{(q)} \mid \mathbf{x}^{[q-1]}\right) \, \mathrm{d}\mathbf{x}^{(q)} \int \left|\hat{P}\left(\mathbf{x}^{[q-1]}\right) - P\left(\mathbf{x}^{[q-1]}\right)\right| \, \mathrm{d}\mathbf{x}^{[q-1]} \\ &\quad + \frac{1}{2} \int \left|\hat{P}\left(\mathbf{x}^{(q)} \mid \mathbf{x}^{[q-1]}\right) - P\left(\mathbf{x}^{(q)} \mid \mathbf{x}^{[q-1]}\right)\right| \, \mathrm{d}\mathbf{x}^{[q-1]} \\ &= \mathrm{TV}\left(\hat{P}\left(\mathbf{x}^{[q-1]}\right), P\left(\mathbf{x}^{[q-1]}\right)\right) + \mathrm{TV}\left(\hat{P}\left(\mathbf{x}^{(q)} \mid \mathbf{x}^{[q-1]}\right), P\left(\mathbf{x}^{[q-1]}\right)\right) \\ &\leq \mathrm{TV}\left(\hat{P}\left(\mathbf{x}^{[q-1]}\right), P\left(\mathbf{x}^{[q-1]}\right)\right) + \varepsilon \cdot \frac{Q}{d}. \end{aligned}$$

⁷⁴⁰ Upon summing up the above inequality for all $q \in [d/Q]$, we obtain

$$\begin{aligned} \mathsf{TV}\left(\hat{P}(\mathbf{x}), P(\mathbf{x})\right) &= \sum_{q=1}^{d/Q} \left(\mathsf{TV}\left(\hat{P}\left(\mathbf{x}^{[q]}\right), P\left(\mathbf{x}^{[q]}\right)\right) - \mathsf{TV}\left(\hat{P}\left(\mathbf{x}^{[q-1]}\right), P\left(\mathbf{x}^{[q-1]}\right)\right) \right) \\ &\leq \sum_{q=1}^{d/Q} \varepsilon \cdot \frac{Q}{d} = \varepsilon \end{aligned}$$

Thus we finish the proof of Theorem 5.

742 **B** Additional Experiments

Algorithm Setup: Our choices of algorithm hyperparameters are based on Song and Ermon (2019). We consider L = 10 different standard deviations such that $\{\lambda_i\}_{i \in [L]}$ is a geometric sequence with $\lambda_1 = 1$ and $\lambda_{10} = 0.01$. For annealed Langevin dynamics with T iterations, we choose the noise levels $\{\sigma_t\}_{t \in [T]}$ by repeating every element of $\{\lambda_i\}_{i \in [L]}$ for T/L times and we set the step size as $\delta_t = 2 \times 10^{-5} \cdot \sigma_t^2 / \sigma_T^2$ for every $t \in [T]$. For vanilla Langevin dynamics with T iterations, we use the same step size as annealed Langevin dynamics. For chained Langevin dynamics with T iterations, the patch size Q is chosen depending on different tasks. For every patch of chained Langevin dynamics, we choose the noise levels $\{\sigma_t\}_{t \in [TQ/d]}$ by repeating every element of $\{\lambda_i\}_{i \in [L]}$ for TQ/dL times and we set the step size as $\delta_t = 2 \times 10^{-5} \cdot \sigma_t^2 / \sigma_{TQ/d}^2$ for every $t \in [TQ/d]$.

752 B.1 Synthetic Gaussian Mixture Model

⁷⁵³ We choose the data distribution P as a mixture of three Gaussian components in dimension d = 100:

$$P = 0.2P^{(0)} + 0.4P^{(1)} + 0.4P^{(2)} = 0.2\mathcal{N}(\mathbf{0}_d, 3\mathbf{I}_d) + 0.4\mathcal{N}(\mathbf{1}_d, \mathbf{I}_d) + 0.4\mathcal{N}(-\mathbf{1}_d, \mathbf{I}_d)$$

Since the distribution is given, we assume that the sampling algorithms have access to the ground-truth score function. We set the batch size as 1000 and patch size Q = 10 for chained Langevin dynamics. We use $T \in \{10^3, 10^4, 10^5, 10^6\}$ iterations for vanilla, annealed, and chained Langevin dynamics. The initial samples are i.i.d. chosen from $P^{(0)}$, $P^{(1)}$, or $P^{(2)}$, and the results are presented in Figures

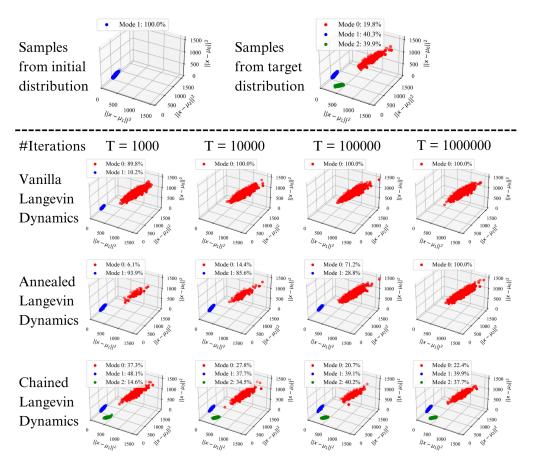


Figure 4: Samples from a mixture of three Gaussian modes generated by vanilla, annealed, and chained Langevin dynamics. Three axes are ℓ_2 distance from samples to the mean of the three modes. The samples are initialized in mode 1.

1, 4, and 5 respectively. The two subfigures above the dashed line illustrate the samples from the initial distribution and target distribution, and the subfigures below the dashed line are the samples generated by different algorithms. A sample **x** is clustered in mode 1 if it satisfies $\|\mathbf{x} - \boldsymbol{\mu}_1\|^2 \le 5d$ and $\|\mathbf{x} - \boldsymbol{\mu}_1\|^2 \le \|\mathbf{x} - \boldsymbol{\mu}_2\|^2$; in mode 2 if $\|\mathbf{x} - \boldsymbol{\mu}_2\|^2 \le 5d$ and $\|\mathbf{x} - \boldsymbol{\mu}_1\|^2 > \|\mathbf{x} - \boldsymbol{\mu}_2\|^2$; and in mode 0 otherwise. The experiments were run on an Intel Xeon CPU with 2.90GHz.

763 B.2 Image Datasets

Our implementation and hyperparameter selection are based on Song and Ermon (2019). During training, we i.i.d. randomly flip an image with probability 0.5 to construct the two modes (i.e., original and flipped images). All models are optimized by Adam with learning rate 0.001 and batch size 128 for a total of 200000 training steps, and we use the model at the last iteration to generate the samples. We perform experiments on MNIST (LeCun, 1998) (CC BY-SA 3.0 License) and Fashion-MNIST (Xiao et al., 2017) (MIT License) datasets and we set the patch size as Q = 14.

For the score networks of vanilla and annealed Langevin dynamics, following from Song and Ermon 770 (2019), we use the 4-cascaded RefineNet (Lin et al., 2017), a modern variant of U-Net (Ronneberger 771 et al., 2015) with residual design. For the score networks of chained Langevin dynamics, we use the 772 official PyTorch implementation of an LSTM network (Sak et al., 2014) followed by a linear layer. 773 For MNIST and Fashion-MNIST datasets, we set the input size of the LSTM as Q = 14, the number 774 of features in the hidden state as 1024, and the number of recurrent layers as 2. The inputs of LSTM 775 include inputting tensor, hidden state, and cell state, and the outputs of LSTM include the next hidden 776 state and cell state, which can be fed to the next input. To estimate the noisy score function, we first 777 input the noise level σ (repeated for Q times to match the input size of LSTM) and all-0 hidden and 778

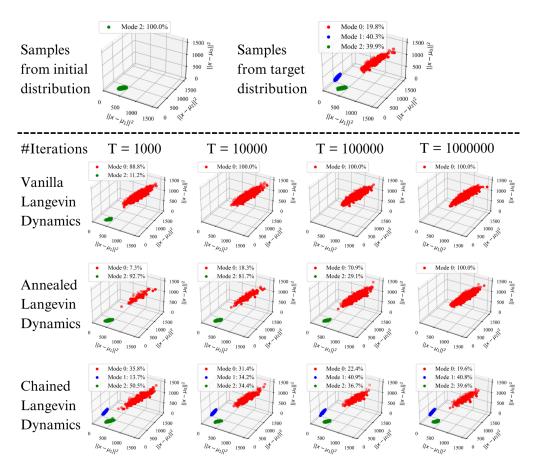


Figure 5: Samples from a mixture of three Gaussian modes generated by vanilla, annealed, and chained Langevin dynamics. Three axes are ℓ_2 distance from samples to the mean of the three modes. The samples are initialized in mode 2.

cell states to obtain an initialization of the hidden and cell states. Then, we divide a sample into d/Qpatches and input the sequence of patches to the LSTM. For every output hidden state corresponding to one patch, we apply a linear layer of size $1024 \times Q$ to estimate the noisy score function of the patch.

To generate samples, we use $T \in \{3000, 10000, 30000, 100000\}$ iterations for vanilla, annealed, and chained Langevin dynamics. The initial samples are chosen as either original or flipped images from the dataset, and the results for MNIST and Fashion-MNIST datasets are presented in Figures 2, 6, 3, and 7 respectively. The two subfigures above the dashed line illustrate the samples from the initial distribution and target distribution, and the subfigures below the dashed line are the samples generated by different algorithms. High-quality figures generated by annealed and chained Langevin dynamics for T = 100000 iterations are presented in Figures 8 and 9.

All experiments were run with one RTX3090 GPU. It is worth noting that the training and inference 790 time of chained Langevin dynamics using LSTM is considerably faster than vanilla/annealed Langevin 791 dynamics using RefineNet. For a course of 200000 training steps on MNIST/Fashion-MNIST, due 792 to the different network architectures, LSTM takes around 2.3 hours while RefineNet takes around 793 9.2 hours. Concerning image generation, chained Langevin dynamics is significantly faster than 794 vanilla/annealed Langevin dynamics since every iteration of chained Langevin dynamics only updates 795 a patch of constant size, while every iteration of vanilla/annealed Langevin dynamics requires 796 computing all coordinates of the sample. One iteration of chained Langevin dynamics using LSTM 797 takes around 1.97 ms, while one iteration of vanilla/annealed Langevin dynamics using RefineNet 798 takes around 43.7 ms. 799

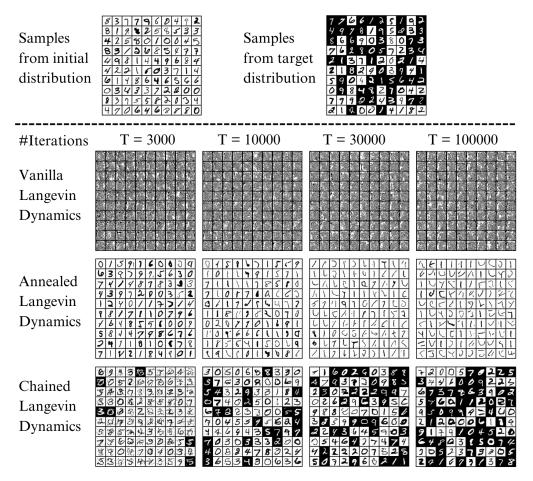


Figure 6: Samples from a mixture distribution of the original and flipped images from the MNIST dataset generated by vanilla, annealed, and chained Langevin dynamics. The samples are initialized as flipped images from MNIST.

800 C Boarder Impacts

This paper presents work whose goal is to advance the field of machine learning. No potential societal consequence of this work needs to be highlighted here.

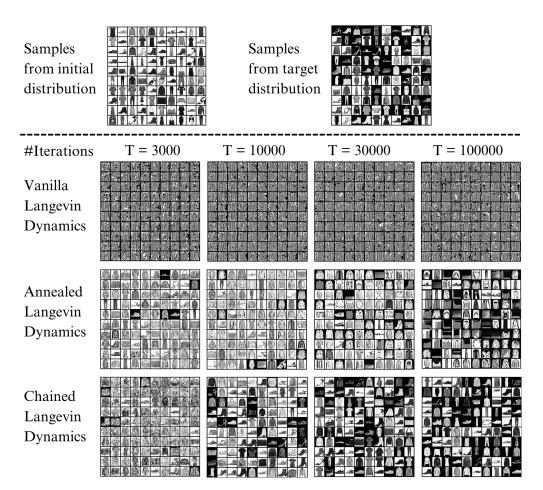


Figure 7: Samples from a mixture distribution of the original and flipped images from the Fashion-MNIST dataset generated by vanilla, annealed, and chained Langevin dynamics. The samples are initialized as flipped images from Fashion-MNIST.

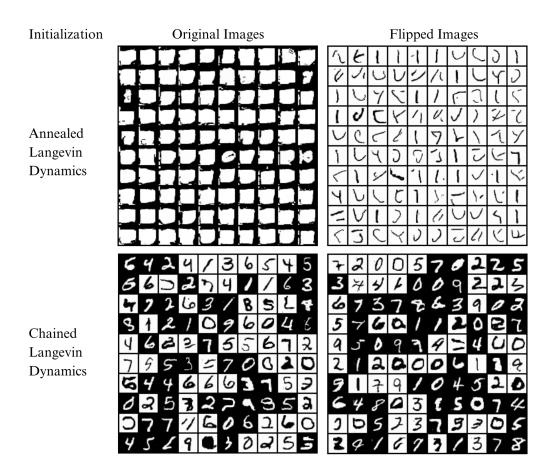


Figure 8: Samples from a mixture distribution of the original and flipped images from the MNIST dataset generated by annealed and chained Langevin dynamics for T = 100000 iterations. The samples are initialized as the original or flipped images from MNIST.

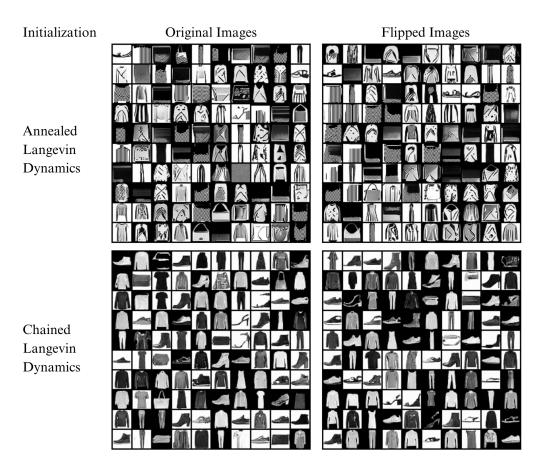


Figure 9: Samples from a mixture distribution of the original and flipped images from the Fashion-MNIST dataset generated by annealed and chained Langevin dynamics for T = 100000 iterations. The samples are initialized as the original or flipped images from Fashion-MNIST.

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