
Wave Interpolation Neural Operator: Interpolated Prediction of Electric Fields Across Untrained Wavelengths

Joonhyuk Seo
Hanyang University
{yhy258}@hanyang.ac.kr

Chanik Kang
Hanyang University
{chanik}@hanyang.ac.kr

Dongjin Seo
Yale University
{dongjin.seo}@yale.edu

Haejun Chung*
Hanyang University
{haejun}@hanyang.ac.kr

Abstract

Existing surrogate solvers are limited to performing inference at fixed simulation conditions, such as wavelengths, and require retraining for different conditions. To address this, we propose Wave Interpolation Neural Operator (WINO), a novel surrogate solver enabling simulation condition interpolation across a continuous spectrum of broadband wavelengths. WINO introduces the Fourier Group Convolution Shuffling operator and a new conditioning method to efficiently predict electric fields from both trained and untrained wavelength data, achieving significant improvements in parameter efficiency and spectral interpolation performance.²

1 Introduction

Metalens is at the forefront of next-generation optical technologies. Designing metalenses involves conducting electromagnetic simulations grounded in the partial differential equations (PDEs) called Maxwell's equations. Generally, the simulations numerically solve Maxwell's equations using finite-difference frequency-domain (FDFD) Hughes et al. [2019, 2018] or finite-difference time-domain (FDTD) Kunz and Luebbers [1993], Oskooi et al. [2010] solvers. However, the requirement of high computational resources makes large-scale simulations and designs intractable Kang et al. [2024a]. To address this challenge, there is a need for surrogate solvers that operate faster than traditional simulations and provide precise results.

Most previous studies have employed conventional data-driven neural networks to predict plausible results based on a constrained number of design factors Jiang et al. [2019], Kang et al. [2024b]. However, as demonstrated in a prior study Kovachki et al. [2023], these models frequently fail to guarantee discretization invariance. This poses difficulties in learning PDEs with mesh information that changes according to wavelength or resolution. Furthermore, they typically perform well within the fixed simulation settings in which they are trained. Generating a new dataset and retraining the model to predict outcomes for new simulation parameters are necessary.

In this work, we propose a parameter-efficient surrogate solver named Wave Interpolation Neural Operator (WINO). For the first time, WINO enables the interpolation of simulation parameters by training with discrete wavelength simulation data through a novel conditioning method, allowing it to

*Corresponding author.

²The code and dataset are available at https://anonymous.4open.science/r/WINO_interp-4B11.

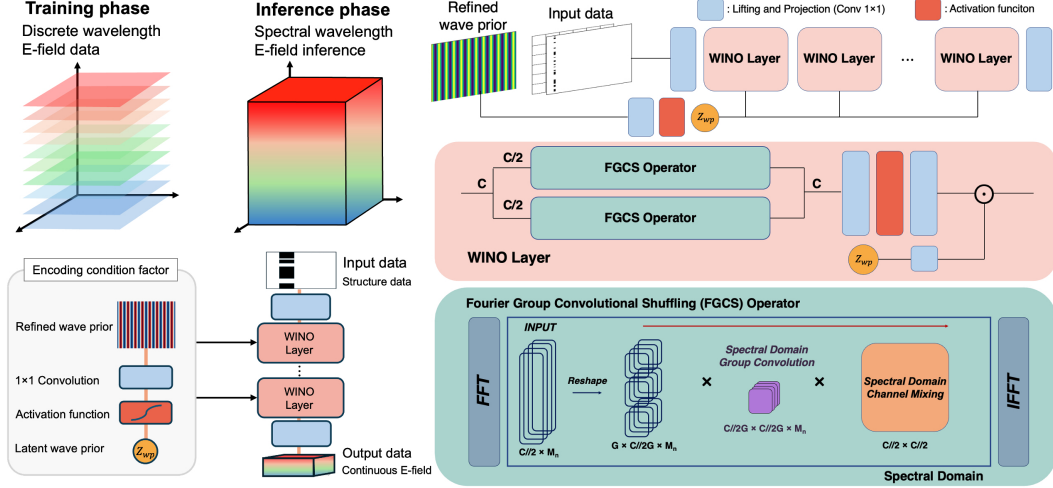


Figure 1: (left) Overall framework of WINO. (right) Schematic illustration of (right top) WINO including (right middle) WINO layer and (right bottom) Fourier Group Convolution Shuffling (FGCS) operator.

infer across a continuous spectrum of broadband wavelengths. The main contributions of our study are as follows:

1. To the best of our knowledge, our work presents the first broadband electromagnetic surrogate solver.
2. Parameter efficiency improves by approximately 74% over the previous state-of-the-art (SOTA) model.
3. Performance improvements: 81% enhancement for untrained wavelengths and 10% enhancement for trained wavelengths compared to the SOTA.

2 Methodology

2.1 Wave Interpolation Neural Operator

We propose Wave Interpolation Neural Operator (WINO), which enables the interpolated prediction of unseen wavelengths from simulation observations of discrete wavelengths (Figure 1). WINO comprises two main components: (1) the Fourier Group Convolution Shuffling (FGCS) operator that is designed to be highly parameter-efficient while maintaining performance, and (2) a novel conditioning method that enables the seamless interpolation of data conditions.

We define the WINO layers as $\mathcal{L}^{(l)}(z^{(l)}) = z^{(l)} + W_2^{(l)} \sigma(W_1^{(l)} \tilde{\mathcal{K}}^{(l)}(z^{(l)} + b_1^{(l)}) + b_2^{(l)})$. We employ a residual connection after two linear transformations with an activation function to enhance the capacity of our model, like Vaswani et al. [2017], Tran et al. [2021]. Detailed information on problem setting is in the Appendix section titled ‘‘Problem Setting.’’

2.2 Kernel Integral Operator $\tilde{\mathcal{K}}$ of WINO

The kernel integral operator in the original FNO is defined as follows:

$$\mathcal{K}^{(l)} = \mathcal{F}^{-1}(R^{(l)} \cdot \mathcal{F}), \quad R \in \mathbb{C}^{M_v \times M_h \times C \times C}$$

where \mathcal{F} is a Fourier transformation, M_v , M_h , and C denote the number of frequency modes of vertical and horizontal components and the number of channels, respectively. Thus, the number of parameters in a 2D problem is $C^2 M_v M_h$. This leads to a significantly high parameter complexity. Thus, as illustrated in Figure 1 (middle), we construct the cross-shaped block Gu et al. [2022] to achieve superior parameter efficiency. We split the inputs into two chunks $z^{(l)} = [z_h^{(l)}, z_v^{(l)}]$ along the channel axis and factorize the Fourier transforms over the horizontal and vertical dimensions $\tilde{\mathcal{K}}^{(l)}(z^{(l)}) = [\mathcal{K}_h^{(l)}(z_h^{(l)}), \mathcal{K}_v^{(l)}(z_v^{(l)})]$. Therefore, the weights $R_h^{(l)} \in \mathbb{C}^{M_h \times \frac{1}{2}C \times \frac{1}{2}C}$ and

$R_v^{(l)} \in \mathbb{C}^{M_v \times \frac{1}{2}C \times \frac{1}{2}C}$ of both two blocks have $\frac{C^2}{4}(M_h + M_v)$ parameters, fewer than the $C^2 M_h M_v$ parameters in the kernel integral operator of an FNO.

Furthermore, we propose an efficient Fourier Group Convolution Shuffling (FGCS) operator that is extremely parameter-efficient while maintaining high performance, as shown in Figure 1 (middle and bottom). We group the input in the Fourier domain along the channel dimension using the group parameter G , as inspired by Xie et al. [2017], Guibas et al. [2021], Kim et al. [2024], but we share weights across the divided groups. Let the horizontal and vertical components be represented by $\mathcal{F}(z_h^{(l)}) \in \mathbb{C}^{H \times M_h \times \frac{1}{2}C}$ and $\mathcal{F}(z_v^{(l)}) \in \mathbb{C}^{M_v \times W \times \frac{1}{2}C}$, respectively. By introducing a group parameter G , horizontal component becomes $R_h^{(l)} \in \mathbb{C}^{M_h \times \frac{1}{2G}C \times \frac{1}{2G}C}$, and the vertical component is transformed into $R_v^{(l)} \in \mathbb{C}^{M_v \times \frac{1}{2G}C \times \frac{1}{2G}C}$. Consequently, we construct a kernel integral operator with a parameter complexity of $C^2 M/G^2$. However, the weight-sharing scheme in the group convolution leads to a significant information loss, resulting in a drastic performance decline. Thus, we further perform channel shuffling in the Fourier domain in order to compensate for the lost information, which can be functioned through simple matrix multiplication along the channel dimension with $W_{\text{ch}}^{(l)} \in \mathbb{C}^{C \times C}$.

2.3 Novel Conditioning Method for Interpolating Untrained Wavelengths' Simulations

The changes in the field wave patterns with varying wavelengths, along with resonance and diffraction phenomena during wave-material interactions, are highly nonlinear. This nonlinearity presents challenges in predicting simulation fields for unseen wavelengths during training. To deal with the challenges, first, we propose a conditioning method for interpolation in the parameter space of the broadband spectrum, as shown in Figure 1 (top and middle). We assume that the condition data have high spatial structure similarity with fields and accurately reflect the features of the wave patterns. Under these assumptions, preserving the spatial structure of the condition information when constructing high-dimensional embeddings is necessary to accurately reflect the field changes according to continuous wavelengths. Thus, we create a common embedding, which is then injected into all layers, using a 1×1 convolutional layer with a nonlinear activation function. Then, each layer adaptively transforms the embedding using an additional 1×1 convolutional layer. We highlight that employing only channel-wise operations guarantees the preservation of the spatial structure of conditional information, which is substantially similar to the field. The final conditional embeddings created in each layer are element-wisely multiplied with the output of the feed-forward network before the residual connection in each layer.

Next, we construct condition information that accurately reflects the wave characteristics by employing a wave prior Gu et al. [2022]. The wave prior is an artificial wave pattern derived from the solution of the wave equation and is expressed as $\mathcal{W}_x = e^{j \frac{2\pi\sqrt{\epsilon_r}}{\lambda} x \mathbf{1}^T \Delta l_x}$ and $\mathcal{W}_z = e^{j \frac{2\pi\sqrt{\epsilon_r}}{\lambda} \mathbf{1} z^T \Delta l_z}$. Here, $\epsilon_r \in \mathbb{C}^{H \times W}$ represents the relative permittivity of the structures at each coordinate, λ is the wavelength, Δl_x and Δl_z are the simulation's step length, and x and z are the coordinates. Because the wave prior emulates the solution of the wave equation, it is substantially similar to the wave patterns of the field generated by the simulator, which result from light propagation according to wavelengths, achieving the precluded assumption. However, the wave prior considers only the structures' permittivity values without accounting for the diffraction phenomena, because it is defined by simultaneously reflecting coordinates. The wave prior exhibits sharp and physically implausible features in the regions where the permittivity values change (Figure 8 (top left and right)). In addition, since the simulation's permittivity information is already provided as input data, including it in the wave prior is redundant. Hence, we propose a refined wave prior that excludes the ϵ_r term, defined as $\mathcal{W}_x = e^{j \frac{2\pi}{\lambda} x \mathbf{1}^T \Delta l_x}$ and $\mathcal{W}_z = e^{j \frac{2\pi}{\lambda} \mathbf{1} z^T \Delta l_z}$ (Figure 8 (bottom left)). The refined condition information aims to provide a more physically plausible representation of light behavior in free space.

3 Results and Discussion

In the following, we compare the interpolation performance for unobserved wavelengths of Unet Ho et al. [2020], Gupta and Brandstetter [2022], FNO Li et al. [2020], and F-FNO Tran et al. [2021], commonly used in a surrogate solver study. Moreover, we evaluate WINO against NeurOLight Gu et al. [2022], the state-of-the-art model for single wavelength field prediction in electromagnetic

Model	Param (M)	Test loss		Design region test loss	
		Untrained	Trained	Untrained	Trained
Unet Ho et al. [2020]	11.5999	0.003332	0.001732	0.00619	0.003775
FNO Li et al. [2020]	3.2868	0.0133	0.007719	0.028339	0.025067
F-FNO Tran et al. [2021]	1.887	0.03309	0.02204	0.043035	0.037091
NeurOLight Gu et al. [2022]	1.65	0.01073	0.001973	0.019503	0.004526
WINO	0.426	0.002035	0.001774	0.003812	0.003929

Table 1: Comparisons of field prediction and interpolation. Untrained represents the loss when predicting fields using unobserved wavelength data, which is the loss for the wavelength interpolation.

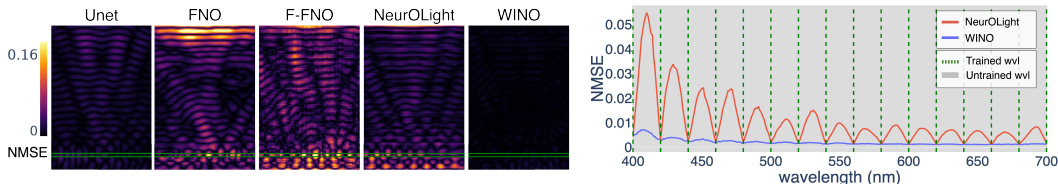


Figure 2: Field interpolation comparisons. (left) Comparison of model errors at 410nm, the unobserved wavelength with the highest error. Darker colors indicate lower errors. Other models show generally high errors in both the design region (green) and the overall area, our model. Nevertheless, WINO demonstrates near-zero errors across the entire area, including the design region. (right) Comparison of entire simulation region errors in the 400-700nm wavelength range.

simulation. Detailed information on model architectures is in the Appendix section titled “Detailed model architectures.”

In EM simulation, while predicting the entire E-field is important, accurately predicting the E-field within the design region is even more crucial. For studies involving surrogate solvers to predict E-fields and subsequently optimize designs, error-free predictions in the design region are essential for designing high-performance photonic devices. However, existing models struggle with the nonlinear resonances observed in design regions. These resonances can vary with different source wavelengths, even for the same structure, leading to significant errors in interpolation.

As shown in Figure 2 (left), the overall models, except WINO, exhibit significant errors in the design region. Such resonance patterns can occur in the metalens used in our experiments and other photonic structures designed at the sub-wavelength scale. Therefore, our model can accurately predict resonance patterns during wavelength interpolation in the design region, making it broadly applicable. Furthermore, we aimed to improve the overall error compared to the SOTA and significantly reduce the error in the design region. Our wave interpolation results achieved SOTA performance in both test loss across the entire region and loss in the design region, compared to existing models Li et al. [2020], Tran et al. [2021], Gu et al. [2022] as shown in Figure 2 (right), Table 1 and Figure 3. Specifically, in the design region, which is represented in Table 1 and Figure 9, our model showed an impressive 80.5% improvement over the previous SOTA model, NeurOLight, for untrained wavelengths and an improvement of 13.2% for trained wavelengths. Moreover, we aimed to drastically reduce the number of parameters, achieving a SOTA model with only 0.426 million parameters, a 74.18% reduction compared to NeurOLight and a 96.33% reduction compared to Unet.

4 Conclusion

In this study, we introduce the Wave Interpolation Neural Operator (WINO), a novel surrogate solver designed to address the limitations of existing models in predicting electric fields across a continuous spectrum of wavelengths. By leveraging Fourier Group Convolution Shuffling (FGCS) and a new conditioning method, WINO significantly improves parameter efficiency and interpolation performance. Our experimental results demonstrate that WINO achieves state-of-the-art performance in both test loss and design region test loss compared to existing models such as FNO, F-FNO, and the previous SOTA, NeurOLight. Furthermore, WINO accurately predicts the nonlinear resonance patterns in photonic structures. This capability and SOTA performance position WINO as a powerful tool, paving the way for broadband photonic applications and enabling the design of highly complex photonic structures with a surrogate solver.

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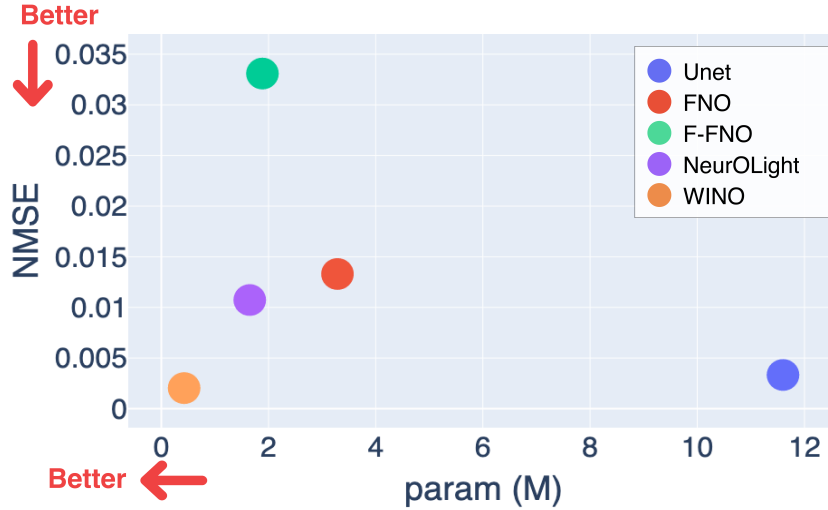


Figure 3: Comparison of WINO with various surrogate solvers in terms of the number of parameters and prediction performance for untrained wavelengths.

A Related Work

A.1 Neural Network-Based Approximations

Light is an electromagnetic wave, and its behavior is governed by PDEs known as Maxwell’s equations. Maxwell’s equations describe the electromagnetic behavior of light such as the mutual induction of electric and magnetic fields, the propagation of light waves, and the interaction of light with matter. Previous efforts in electromagnetic field prediction have focused on training neural networks to approximate field values. Data-driven neural network research Jiang et al. [2021], Trivedi et al. [2019], An et al. [2020] attempted to learn and predict Maxwell’s equations using simulated field data. However, these approaches often struggled to accurately capture the physical characteristics of Maxwell’s equations. To address this issue, Physics-Informed Neural Networks (PINNs) were employed in Physics-Augmented Deep Learning research Chen et al. [2022]. This method combined data loss and Maxwell loss as the total loss, integrating Maxwell’s equations into the field prediction process for given structures. While this approach improved prediction speed compared to traditional FDFD solvers, it still had significant limitations. Specifically, it required datasets for each individual wavelength and could only predict fields for the trained wavelengths. Additionally, the field predictions were confined to small and simple photonic structures.

A.2 Advancements in Operator Learning

Previous research using data-driven approaches has struggled to accurately learn and infer the characteristics of Maxwell’s equations. As a result, researchers have begun exploring neural network-based operators designed to model and predict complex physical systems Azizzadenesheli et al. [2024], Li et al. [2020]. Recent advancements in the field of Operator Learning offer promising solutions that effectively learn PDE characteristics while preserving discretization invariance. The FNO proposed by Li et al. [2020] presents a deep learning architecture designed to learn mappings between infinite-dimensional function spaces. This approach significantly enhances the efficiency and accuracy of solving PDEs. The FNO achieves resolution invariance and supports zero-shot super-resolution, establishing it as a pioneering method in operator learning. After FNO has been published, the Factorized-FNO Tran et al. [2021], named F-FNO, improves the performance of the original FNO by using separable spectral layers and enhanced residual connections, significantly reducing errors in various PDE problems. Additionally, Multi-Grid Tensorized Fourier Neural Operator (MG-TFNO) Kossaifi et al. [2023]. MG-TFNO introduces tensor decomposition techniques that significantly reduce the memory footprint while maintaining model expressivity. Specifically,

tensorization might help optimize our model’s performance when predicting the electric field over a broader range of wavelengths, without compromising the computational efficiency.

In electromagnetic field prediction, the state-of-the-art model NeurOLight Gu et al. [2022] leverages the FNO. Gu et al. [2022] encodes wave prior by combining wavelength information and material permittivity, effectively learning the PDEs of Maxwell’s equations. The model has demonstrated superior parameter efficiency and higher accuracy compared to previous studies. However, it still has limitations as it predicts fields for a limited narrow range of wavelengths and grid steps, making it less suitable for broadband wavelength applications.

B Problem Setting

Variable	Meaning	Dimensionality
$R_h^{(l)}$	Horizontal component weights in Fourier domain	$M_h C^2 / 4G^2$
$R_v^{(l)}$	Vertical component weights in Fourier domain	$M_v C^2 / 4G^2$
$K^{(l)}$	Kernel integral operator in the original FNO	$C^2 M_h M_v$
$z^{(l)}$	Input chunk at layer l	Infinite
L	Number of WINO layers	1
\mathcal{F}	Fourier transform	Infinite
\mathcal{F}^{-1}	Inverse Fourier transform	Infinite
σ	Non-linear activation function	Infinite
b	Bias vector	Infinite
$W_{ch}^{(l)}$	Channel shuffling weights	C^2
G	Number of groups in Fourier space	Between 1 and $\frac{1}{2} \min\{s_1, \dots, s_d\}$
λ	Wavelength of light	1
ϵ_r	Relative permittivity	Infinite
$\mathcal{W}_x, \mathcal{W}_z$	Wave prior in the x and z directions	Infinite
$\Delta l_x, \Delta l_z$	Simulation step length in x and z directions	1
a	Input function	Infinite
u	Output function	Infinite
d_a, d_u	Dimension of input/output co-domain	1

Table 2: Table of notation

Let $\Omega \subset \mathbb{R}^d$, $\mathcal{A} = \mathcal{A}(\Omega; \varepsilon^{d_a})$ where $\varepsilon = \{\epsilon_{\text{air}}, \epsilon_{\text{material}}\}$, and $\mathcal{U} = \mathcal{U}(\Omega; \mathbb{C}^{d_u})$ be a bounded open set of the underlying domain, infinite-dimensional spaces of the relative permittivity and field of the simulation. ϵ_{air} and $\epsilon_{\text{material}}$ represent the relative permittivity values of air and material. Additionally, we assume that $\mathcal{W} \subset \mathbb{R}_{>0}$ is the broadband range, such as the visible light spectrum, and $\tilde{\mathcal{W}} \subset \mathcal{W}$ is a discrete collection of wavelengths evenly distributed at a specific interval from \mathcal{W} . Our model G_θ learns an ideal electromagnetic simulator $G^\dagger : \mathcal{A} \rightarrow \mathcal{U}$ for continuous wavelengths \mathcal{W} by mapping between infinite-dimensional function spaces using a finite set of Maxwell PDE input-output pairs $\{(w_j, a_j), u_j\}_{j=1}^N$, where $w \in \mathcal{W}$. Similar to a prior work Kovachki et al. [2023], WINO has an iterative process to map between infinite-dimensional function spaces.

$$u = \mathcal{G}_\theta(a) = (\mathcal{Q} \circ \mathcal{L}^{(L)} \circ \dots \circ \mathcal{L}^{(1)} \circ \mathcal{P})(a)$$

where $\mathcal{L}^{(l)}$ is a l -th operator layer, L indicates the number of layers, and \mathcal{P}, \mathcal{Q} are lifting and projection operators, respectively.

C Experiment Setup

C.1 Comprehending Optical Field Simulation

To collect training and validation datasets, we employ the FDFD method which is one of the numerical methods used to solve Maxwell’s equations. The FDFD method discretizes the spatial domain into coordinates and solves electromagnetic fields in the frequency domain, making it particularly effective for steady-state analysis. This method is especially effective for generating field data for single frequencies as light passes through the metalens. Resonances can occur in subwavelength photonic

structures and significantly affect their performance. However, predicting the resonance of electric fields in subwavelength structures becomes challenging when the wavelength changes due to the nonlinearity of resonance Dorodnyy et al. [2023], Nieto-Vesperinas [2020]. Therefore, surrogate solvers trained on specific wavelengths struggle to predict fields for different wavelengths, even when applied to the same structure. For a detailed explanation of the nonlinearity of resonance, refer to the Appendix section titled “Nonlinearity of Subwavelength Photonic Structure Resonance.”

C.2 FDFD Simulation Setup

In our study, we utilize Ceviche Hughes et al. [2019], an open-source FDFD simulation tool, to generate electromagnetic field data for photonic structures. Specifically, we analyze the field of a metalens structure, simulating only half of the cylindrical symmetry lens to reduce computation time. We implement the FDFD simulation with a grid-based approach at a spatial resolution of 40 points per μm . The simulation domain measures $5.2 \mu m \times 6.85 \mu m$, excluding a $1\text{-}\mu m$ thick perfectly matched layer (PML) at the boundaries. We use SiO_2 , a material commonly employed in the design of metalenses Park et al. [2019]. As demonstrated in Figure 4 (right), we randomly place either SiO_2 or air within the minimum design grid size, a width of 75nm , in the design area. This simulation setup generates data for the E and H fields when visible light passes through the structure.

C.3 Datasets

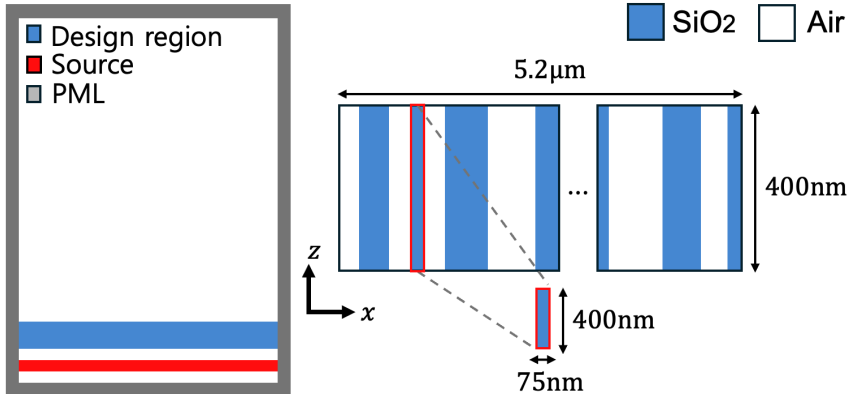


Figure 4: FDFD simulation setup overview. (left) The entire simulation space includes the metalens structure. The blue region represents the design region of the lens, red line represents the visible light source, and the gray region represents the PML used to minimize reflections. (right) The design region of the metalens, shown as the blue area in (left).

We aim to predict the E-field for all wavelengths from discrete data in the $400\text{-}700\text{nm}$ (visible light) range. To achieve this, we generate field data for discrete wavelengths as our dataset. Different lens structures are randomly placed in the design region for each data, and the field data resulting from FDFD simulations, along with the permittivity(ϵ) of the materials, wavelength(λ), and structural information, are used as pairs for training. Specifically, the training set consists of 12,000 data sampled randomly at 20nm intervals within the $400\text{-}700\text{nm}$ wavelength range. As mentioned above, our goal is to achieve seamless field prediction across the broadband wavelength. Therefore, to use unseen data for testing, we use 6,020 data each for the test and validation sets, composed of 20 data for each wavelength, sampled at 1nm intervals between 400 and 700nm .

C.4 Training Details

Our model training runs for 200 epochs with 32 batch size. The best model is selected based on the lowest loss in predicting previously unseen wavelengths during the training phase. We employ the GELU (Gaussian Error Linear Unit) activation function and the AdamW optimizer with the following parameters: a learning rate of 0.002, beta values of (0.9, 0.000), epsilon set to 10^{-8} , and a weight decay factor of 0.0001. To dynamically adjust the learning rate, we utilize a cosine annealing learning rate scheduler with a minimum learning rate of 0.00001. Mode is set to (50, 60). We leverage

PyTorch, and the training process is conducted using a single RTX 6000 Ada and an AMD EPYC 7763 CPU.

C.5 Training Objective and Evaluation Metric

Fields typically exhibit different statistics despite the fixed source power. We employ the normalized mean squared error (N-MSE) objective, $L(G_\theta(a), G^\dagger(a)) = (\|G_\theta(a) - G^\dagger(a)\|_2^2 / \|G^\dagger(a)\|_2^2)$, to distribute the optimization effort evenly across several field data. We also use N-MSE for evaluation.

D Ablation Study

D.1 Hyperparameters of FGCS Operator

Weight sharing	Channel shuffling	Groups	Channels	Params	Test loss	
					Untrained	Trained
		1	32	0.423	0.00413	0.00361
		2	48	0.562	0.00278	0.00193
		4	64	0.654	0.00251	0.00165
✓		2	48	0.372	0.00398	0.00342
✓		4	64	0.401	0.00234	0.00199
	✓	2	48	0.576	0.0033	0.00276
	✓	4	64	0.679	0.00211	0.00153
✓	✓	2	48	0.386	0.00366	0.00302
✓	✓	4	64	0.426	0.00204	0.00177
		8	128	1.898	0.00162	0.00092
✓		8	128	1.306	0.00123	0.00091
	✓	8	128	1.996	0.00145	0.00081
✓	✓	8	128	1.405	0.00105	0.00081

Table 3: Ablation of design components of proposed FGCS operator. The ✓ indicates the use of the design components.

We assess the validity of the hyperparameters (groups and channels) and design choices (weight sharing and channel shuffling) of the proposed FGCS operator by conducting an extensive ablation study, as shown in Table 3. We increase the number of groups in proportion to the growth of channels to consider a consistent number of parameters in all cases. Row 9 presents the FGCS operator used in our model. Rows 1-3, which do not utilize weight sharing and channel shuffling in the Fourier domain, follow the same operational form as described in a previous study Kim et al. [2024]. When only weight sharing is employed (rows 4-5), the prediction of fields for both unobserved and observed wavelengths during the training phase demonstrates worse performance compared to that in the previous study Kim et al. [2024] due to a substantially small number of parameters. However, as the number of groups and channels increases, the performance becomes competitive. Notably, the performance in predicting the fields for unobserved wavelengths during training is better. Thus, the weight-sharing scheme significantly aids in generalization for the unobserved spectral parameters. In addition, comparing the previous method Kim et al. [2024] (rows 2-3) with the case where only channel shuffling is utilized (rows 6-7), the field prediction performances for both the untrained and trained wavelengths are better. This suggests that employing the channel shuffling method with the grouping method significantly improves the overall performance. Following the design configuration of the FGCS operator (rows 8-9), we can construct a parameter-efficient model that achieves an overall performance improvement as the number of groups and channels is heightened. While the FGCS operator shows slightly lower prediction performance for trained wavelengths than when only a channel shuffling scheme is employed, it demonstrates superior performance for untrained wavelengths. Therefore, the proposed FGCS operator contributes to improving interpolation performance for unseen broadband spectrum parameters during training. Rows 10-13 present varying performances when increasing the model width. The performance improvement of the FGCS operator becomes increasingly significant with the addition of more channels. In particular, there is a substantial improvement when the number of groups and channels increases to 8 and 128, respectively. We believe that excessively sparse constructions of weights tensor and information compensating with channel shuffling may improve the generalization capacity of large models.

D.2 Conditioning Methods

We conduct a comparative analysis of our proposed method and two other conditioning methods to evaluate our conditioning method’s efficacy.

Conditioning method	Params	Test loss	
		Untrained	Trained
Spectral parameter conditioning	1.086	0.13761	0.002577
Concatenating wave prior	0.326	0.002821	0.002022
Our conditioning method	0.426	0.002035	0.001774

Table 4: Ablation study of several conditioning methods.

Table 4 shows the results of the ablation study of several conditioning methods employed on WINO. As in the previous experiments, The terms “untrained” and “trained” represent the setups in which fields are predicted using untrained and trained wavelength data, respectively. The conditioning methods used for comparison are spectral parameter conditioning Gupta and Brandstetter [2022], Mao et al. [2024] and the concatenating wave prior Gu et al. [2022]. The spectral parameter conditioning method transforms the scalar wavelength into a vector via sinusoidal embedding Vaswani et al. [2017]. It then employs a two-layer feed-forward network to project the embeddings onto a higher-dimensional space of $4 \times$ hidden channels. The embeddings are mapped to the Fourier space using the module named FreqLinear module described in Gupta and Brandstetter [2022]. Finally, the embeddings in the Fourier space are mode-wise multiplied with the input of the Fourier integral kernel to perform parameter conditioning in the spectral domain. The concatenating wave prior injects the condition information by simply concatenating the given refined wave prior to the input data. Although the spectral parameter conditioning method effectively predicts the fields of wavelengths observed during training, it fails to perform interpolation in the parameter space of the broadband spectrum. Thus, the spectral parameter conditioning method is incapable of learning the patterns of fields with continuously varying wavelengths. Our conditioning method thoroughly outperforms the concatenating wave prior conditioning method in both the trained and untrained wavelength cases, achieving a 12.3% improvement in trained wavelengths and a 27.9% improvement in untrained wavelengths. Consequently, the proposed method is effective for injecting conditions that are substantially similar to the fields in terms of the spatial and structural features. Furthermore, we conducted an additional comparative analysis between our conditioning method and other methods using various models. For a detailed explanation of the more ablation studies, refer to the Appendix section titled “Additional Ablation Studies for Conditioning Methods.”

As the influence of simulation parameters like higher relative permittivity increases, the nonlinear optics effects become more pronounced, and the precluded assumption weakens. However, we found that our method still works reasonably well even in these cases. For a detailed experiments, refer to the Appendix section titled “Additional Experiments for the Weakened Precluded Assumption.”

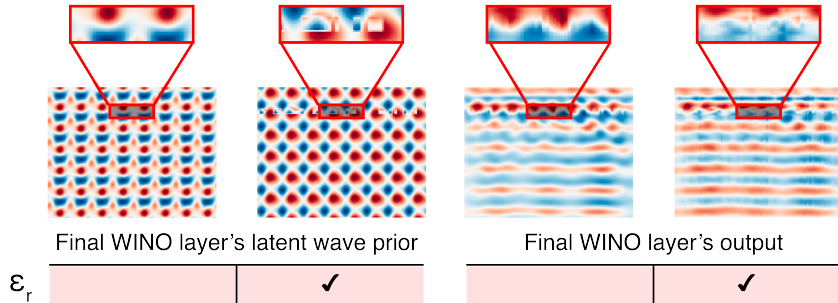


Figure 5: Visualization of the features of condition information and outputs from the final WINO layers for a specific part of the entire region for 400nm wavelength data. (left) The latent wave prior in the final WINO layer, excluding and including the ϵ_r term. (right) The output of the final WINO layer without and with considering the ϵ_r term. The magnified sections are part of the region where structures exist.

D.3 Permittivity in the Wave Prior

ϵ_r term	Test loss	
	Untrained	Trained
✓	0.002372	0.002148
	0.002035	0.001774

Table 5: Ablation study for ϵ_r term in wave prior.

As shown in Table 5, we demonstrate performance enhancement resulting from the exclusion of the ϵ_r term from the wave prior through an ablation study. To identify the reason for the improvement, we investigate the latent wave prior and the output of the last WINO layer. Figure 5 (left) shows the latent wave prior before element-wise multiplication in the last WINO layer. Sharp features emerge in regions with structures when the wave prior reflects the ϵ_r term because our conditioning method utilizes shallow 1×1 convolutional layers to ensure the discretization invariance property and maintain the spatial structure of the refined wave prior. Figures 5 (right) show the outputs of the final WINO layer. The result in Figure 5 (right), where the ϵ_r term is included in the wave prior, has bizarre features present throughout the entire region compared to when the ϵ_r term is not included. Therefore, we conclude that the ϵ_r term in the wave prior introduces unintended bias.

E Detailed Model Architectures

E.1 Unet

We use a 4-level modern convolutional Unet architecture with an initial channel size of 16. Each level consists of two residual blocks with GELU activation function and Group Normalization. We concatenate the refined wave prior with the input for the conditioning method.

E.2 FNO

We set the number of Fourier layers and the number of channels in each Fourier layer to 5 and 32, respectively. The frequency modes for the z and x axes are set to [32, 10]. In addition, the lifting operator is a linear operator, and the projection operator is a 2-layer feed-forward network with GELU activation and a dimension of 128. We concatenate the refined wave prior with the input to inject condition information.

E.3 F-FNO

We employ 12 Fourier layers, each consisting of 64 channels. The frequency modes for the z and x axes are set to [50, 60]. The lifting operator is a linear operator, and the final projection operator is a 2-layer feed-forward network with GELU activation function and a dimension of 256. To incorporate condition information, we concatenate the refined wave prior with the input.

E.4 NeurOLight

We use 12 Fourier layers with 64 channels each. The frequency modes for the z and x axes are set to [50, 60]. The channel expansion factor for the convolutional modules in the NeurOLight layers is configured to be 2. To inject condition information, we concatenate the wave prior with the input. During training, we implement stochastic network depth with a rate of 0.1 to the residual NeurOLight layers to alleviate overfitting.

E.5 WINO

We use 12 WINO layers with 64 channels each and specify the number of groups as 4. The frequency modes for the z and x axes are configured as [50, 60]. We implement a channel expansion scheme with a factor of 2 for the 2-layer feed-forward network in the WINO layers, inspired by Xie et al. [2021], Gu et al. [2022]. The lifting operator is linear, and the final projection operator is a 2-layer feed-forward network with GELU activation and a dimension of 256.

Model	Param (M)	Test loss		Design region test loss	
		Untrained	Trained	Untrained	Trained
Unet-Bottleneck	11.6371	0.11786	0.03303	0.059753	0.05574
Unet-AdaGN	11.779	0.09342	0.03548	0.047393	0.034908
Unet	11.5999	0.003332	0.001732	0.00619	0.003775
FNO	3.2868	0.0133	0.007719	0.028339	0.025067
FNO-Ours	3.2974	0.009966	0.008369	0.027739	0.028831
F-FNO	1.887	0.03309	0.02204	0.043035	0.037091
F-FNO-Ours	1.92	0.007606	0.006705	0.014702	0.014724

Table 6: Comparisons of field prediction and interpolation with various conditioning methods and models. Unet, FNO, and F-FNO refer to concatenating the wave prior with the input data. ‘‘Ours’’ represents our conditioning method.

Model	Param (M)	Test loss		Design region test loss	
		Untrained	Trained	Untrained	Trained
Unet	11.5999	0.017165	0.017799	0.027481	0.02874
FNO	3.2868	0.120722	0.120426	0.333563	0.33381
FNO-Ours	3.2974	0.120246	0.122924	0.351337	0.348711
F-FNO	1.887	0.100605	0.106164	0.192996	0.192103
F-FNO-Ours	1.92	0.067897	0.067897	0.140872	0.140648
NeurOLight	1.65	0.017074	0.011558	0.033887	0.026682
WINO	0.426	0.007932	0.008131	0.018456	0.017944

Table 7: Comparisons of field prediction and interpolation with high relative permittivity.

F Additional Ablation Studies for Conditioning Methods

To verify the effectiveness of our conditioning method, we conduct an additional ablation study by applying the proposed conditioning method to other models and comparing it with other conditioning methods. ‘‘Unet-Bottleneck’’ maps the scalar wavelength value into a vector using a simple three-layers feed-forward network with GELU and then adds the induced condition vector to the input of the bottleneck in Unet. ‘‘Unet-AdaGN’’ reflects the wavelength condition in the model using Adaptive Group Normalization (AdaGN) Dhariwal and Nichol [2021]. The given wavelength value is transformed using sinusoidal embedding Vaswani et al. [2017]. Subsequently, a two-layer feed-forward network is used to project the embeddings into a higher-dimensional space of $4 \times$ hidden channels. The condition embeddings are incorporated into the Unet through AdaGN. ‘‘Unet, FNO, and F-FNO’’ refers to the method of concatenating the wave prior with the input data. Finally, ‘‘Ours’’ refers to our proposed conditioning method. According to Table 6, our method significantly improves the interpolation performance for unobserved wavelengths compared with the various conditioning methods of other models. When our conditioning method is applied to the FNO model, the field prediction performance for the wavelengths used during training is marginally reduced compared to when the method is not applied. However, there is a substantial performance improvement throughout the entire region for the unobserved wavelength data. In particular, the performance enhancement is dramatic in the case of F-FNO, which has conspicuously fewer parameters than Unet and FNO, factorizes the Fourier transforms over the problem dimensions, and is composed of many layers. Applying our conditioning method to F-FNO leads to a 77% increase in performance for unobserved wavelengths and a 67.6% increase for observed wavelength data across the entire simulation region compared with not using the conditioning method. In addition, in the design region, there is a 65.8% performance improvement for unobserved wavelengths and 60% improvement for observed wavelengths. We highlight that the error in the design region is consistent for both observed and unobserved wavelengths. Thus, we verify that our conditioning method substantially contributes to the interpolation capability.

G Additional Experiments for the Weakened Precluded Assumption

We conducted additional experiments to evaluate the performance of our proposed method when the precluded assumption (that the field and condition data have high similarity) becomes weak. We

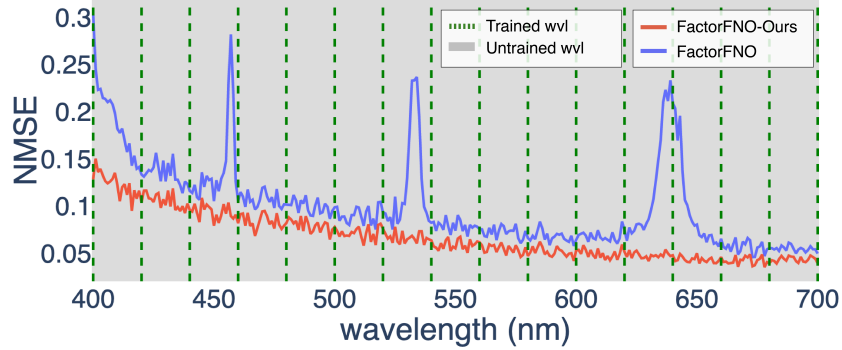


Figure 6: Field interpolation comparisons of F-FNO with simply concatenating conditioning and our conditioning method. Comparison of entire simulation region errors in the 400-700nm wavelength range.

Model	Param (M)	Test loss		Design region test loss	
		Untrained	Trained	Untrained	Trained
WINO	0.426	0.007932	0.008131	0.018456	0.017944
WINO ($l=1$)	0.417	0.007626	0.007741	0.017677	0.01782
WINO ($l=2$)	0.409	0.007534	0.007552	0.017484	0.017566
WINO ($l=3$)	0.401	0.008208	0.008391	0.019047	0.019464

Table 8: Field prediction and interpolation comparisons between injecting the condition into all but the last l WINO layers when relative permittivity is high.

set up the simulation with a more complex field using a material with higher relative permittivity (lithium niobate, refractive index $n \approx 2.3$)³.

FNO is inefficient in terms of the number of parameters. Despite having fewer layers and widths and lacking additional modules like fully connected layers, it has the highest number of parameters among the neural operators compared. As shown in Table 7⁴, FNO struggles to learn complex optical properties even when our conditioning method is applied due to the limited learning capacity caused by the given problems. In contrast, F-FNO has sufficient capacity because of its superior parameter efficiency, which allows for many layers, widths, and additional MLPs. F-FNO demonstrates inferior performance despite its superior learning capacity if concatenating the wave prior with input data for conditioning. In particular, it exhibited peaks in the loss at specific wavelengths (Figure 6). On the other hand, applying our conditioning method to F-FNO significantly improves the field prediction performance and resolves the peak problem. This demonstrates that our method effectively helps to learn complex optical phenomena for a narrow range of wavelengths when sufficient learning capacity is available. While NeuOLight exhibits superior performance in field inference, it struggles to predict E-field across the broadband spectrum. However, WINO still demonstrates the best performance even when the precluded assumption becomes weak.

G.1 Weakening the Precluded Assumption in WINO

WINO element-wisely multiplies latent condition data at every layer. While the conditioning method serves as a powerful regularization means for interpolation on a wide range of conditions, it can have

³Permittivity is the square of the refractive index, denoted by n , which indicates how much light is refracted when it passes through a medium. Therefore, materials with higher permittivity have higher refractive indices. Consequently, when light passes through a material with high permittivity, it slows down more than in a material with lower permittivity, resulting in a more complex field pattern.

⁴The inference performance for untrained wavelengths often appears better than for trained wavelengths. This is because nonlinear optics phenomena (diffraction, resonance) caused by the high epsilon occur strongly, resulting in the loss exhibiting a slightly exponential change with respect to the wavelength. Consequently, if spectrum (condition) interpolation is successful, it can result in lower loss values for untrained wavelengths.

adverse effects when the precluded assumption is weakened due to more complex fields simulated from materials with high relative permittivity.

To assess the effects of weakening the assumption in WINO, we conduct an ablation study, injecting information into layers except for the last l layers.

As shown in Table 8, injecting the condition into all but the last 1 and 2 WINO layers ($l = 1, 2$) reflects weakened assumption well and leads to performance improvement. However, injecting the condition into all but the final 3 WINO layers ($l = 3$) results in a decrease in both field inference performance and interpolation performance. This demonstrates the trade-off between the powerful regularization of our conditioning method and reflecting the weakened assumption.

H Nonlinearity of Subwavelength Photonic Structure Resonance

Mie scattering Dorodnyy et al. [2023], Nieto-Vesperinas [2020], a fundamental mechanism describing the interaction of electromagnetic waves with particles smaller than the wavelength of light, is pivotal in understanding resonance phenomena in subwavelength photonic structures. The mathematical formulation of Mie scattering is expressed through an infinite series of spherical harmonics, where each term represents a different mode of scattering:

$$\sigma(\lambda, r) = \frac{2\pi}{k^2} \sum_{n=1}^{\infty} (2n+1) (|a_n|^2 + |b_n|^2)$$

where λ is the wavelength, r is the radius of the particle, $k = \frac{2\pi}{\lambda}$ is the wave number, and a_n, b_n are the scattering coefficients for electric and magnetic modes, respectively. The nonlinearity arises from the complex interaction between these modes, which can be highly sensitive to structural and material parameters, making accurate predictions difficult.

Furthermore, high Q-factors are indicative of sharp resonance peaks and narrow bandwidths in the response of photonic structures, which implies that the system has a high selectivity in frequency response. While beneficial for many applications, high Q-factors also mean that the resonance is extremely sensitive to slight deviations in system parameters, such as changes in material properties or geometric alterations. This sensitivity leads to significant challenges in predicting the behavior of the system under slightly altered conditions:

$$Q = \frac{\omega_0}{\Delta\omega}$$

where ω_0 is the resonance frequency and $\Delta\omega$ is the bandwidth of the resonance.

At subwavelength scales, interactions between light and matter involve complex phenomena like electric and magnetic Mie resonances. Therefore, surrogate solvers trained at a single wavelength struggle to predict nonlinear resonances in regions with matter across broadband wavelengths. Our WINO model overcomes this challenge, accurately predicting resonances arising from structures.

We employ a wave prior based on the solution of the wave equation as the condition data ($\mathcal{W}_x = e^{j\frac{2\pi\sqrt{\epsilon_r}}{\lambda}x} \mathbf{1}^T \Delta l_x$ and $\mathcal{W}_z = e^{j\frac{2\pi\sqrt{\epsilon_r}}{\lambda}z} \mathbf{1}_z^T \Delta l_z$). However, when the relative permittivity, ϵ_r , reflects the materials used in the simulation, physically implausible features appear in the regions where the material is present (Figure 8). To resolve this, we set ϵ_r to 1, assuming the entire materials are air. This approach allows us to obtain physically plausible features in free space.

I Limitations and Future Work

Our weight grouping and sharing method reduces the number of MAC (Multiply-ACcumulate) in the matrix multiplication within the Kernel Integral Operator by a factor of $1/G$ compared to conventional FNOs and decreases the number of parameters by a factor of $1/G^2$. However, WINO requires more computation time (Table 9). This increase in time complexity is attributed to the model’s employment of a greater number of layers than the vanilla FNO, as well as the reshaping and concatenation of tensors within each layer in the complex domain to construct the cross-shaped architecture and FCGS operation. On the other hand, despite the additional tensor reshaping processes in the complex domain, WINO still achieves faster inference speeds compared to NeurOLight, which uses a nonlinear Lifting operator and complicated feed-forward networks. Furthermore, our proposed

Model	Runtime (s)	Param (M)
Unet	0.018	11.5999
FNO	0.006	3.2868
FNO-Ours	0.008	3.2974
F-FNO	0.035	1.887
F-FNO-Ours	0.041	1.92
NeurOLight	0.101	1.65
WINO	0.08	0.426

Table 9: Comparisons of the number of parameters and runtime. Runtime is averaged over ten trials with 16 batch sizes.

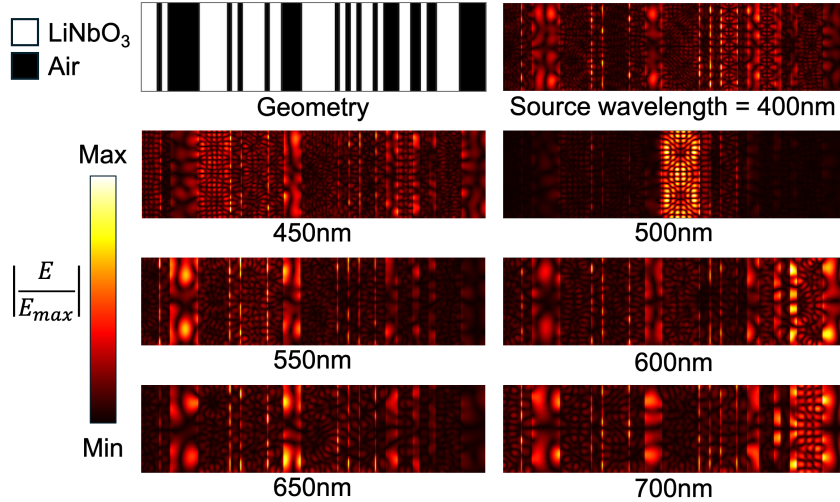


Figure 7: Using the FDTD simulation tool Meep, the E-field at the same time step is visualized for a design region composed of LiNbO₃ and Air. This visualization confirms that resonance occurs nonlinearly with respect to wavelength.

conditional method does not significantly add to the computation time, which is one of the critical factors for improving overall field prediction performance.

We leave the task of optimizing the architecture to deal with the time complexity issues for future work. For future work, we plan to extend optical simulations to a centimeter-scale framework and expand the wavelength range from the visible spectrum to the infrared (IR) bands. Additionally, beyond the metalens-based simulations conducted in the present study, we plan to explore other photonic structures, such as mode converters and waveguides.

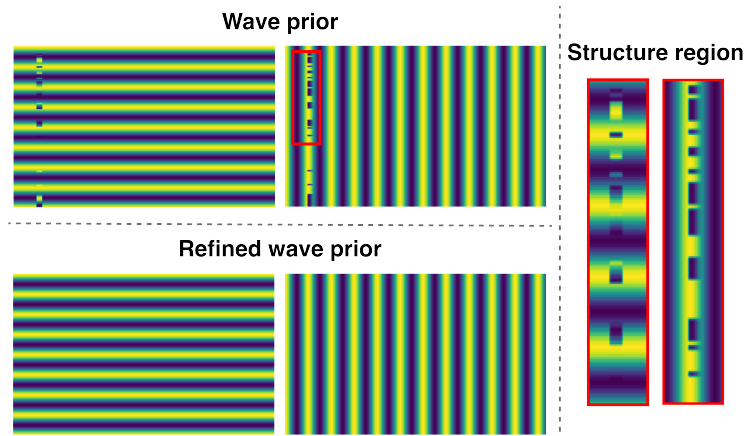


Figure 8: Visualization of (top left) wave prior that includes relative permittivity of materials, (right) physically implausible features of wave prior including relative permittivity of materials, and (bottom left) wave prior that excludes relative permittivity of materials.

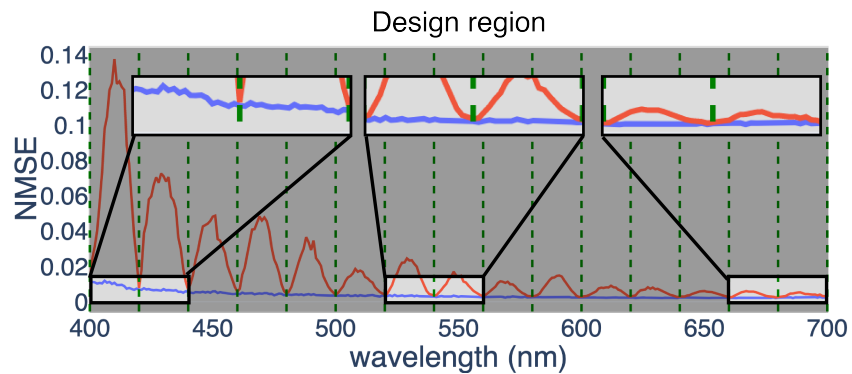


Figure 9: Comparison of design simulation region errors in the 400-700nm wavelength range. The field prediction loss of WINO, particularly in the design region, exhibits a linear change as the wavelengths vary.

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