# Sparse identification of nonlinear dynamics WITH SHALLOW RECURRENT DECODER NETWORKS

Anonymous authors

Paper under double-blind review

### ABSTRACT

Spatio-temporal modeling of real-world data is a challenging problem as a result of inherent high-dimensionality, measurement noise, and expensive data collection procedures. In this paper, we present Sparse Identification of Nonlinear Dynamics with SHallow Recurrent Decoder networks (SINDy-SHRED) to jointly solve the sensing and model identification problems with simple implementation, efficient computation, and robust performance. SINDy-SHRED utilizes Gated Recurrent Units (GRUs) to model the temporal sequence of sensor measurements along with a shallow decoder network to reconstruct the full spatio-temporal field from the latent state space using only a few available sensors. Our proposed algorithm introduces a SINDy-based regularization. Beginning with an arbitrary latent state space, the dynamics of the latent space progressively converges to a SINDy-class functional, provided the projection remains within the set. We conduct a systematic experimental study including synthetic PDE data, real-world sensor measurements for sea surface temperature, and direct video data. With no explicit encoder, SINDy-SHRED allows for efficient training with minimal hyperparameter tuning and laptop-level computing. SINDy-SHRED demonstrates robust generalization in a variety of applications with minimal to no hyperparameter adjustments. Additionally, the interpretable SINDy model of latent state dynamics enables accurate long-term video predictions, achieving state-of-the-art performance and outperforming all baseline methods considered, including Convolutional LSTM, PredRNN, ResNet, and SimVP.

033

004

010 011

012

013

014

015

016

017

018

019

021

025

026

027

028

### 1 INTRODUCTION

034 Modeling unknown physics is an exceptionally challenging task that is complicated further by the computational burden of high-dimensional state spaces and expensive data collection. Partial dif-035 ferential equations (PDEs) derived from first principles remain the most ubiquitous class of models to describe physical phenomena. However, we frequently find that the simplifying assumptions 037 necessary to construct a PDE model can render it ineffectual for real data where the physics is multi-scale in nature, only partially known, or where first principles models currently do not exist. In such cases, machine learning (ML) method offers an attractive alternative for learning both the 040 physics and coordinates necessary for quantifying observed spatio-temporal phenomenon. Many 041 recent efforts utilizing ML techniques seek to relax the computational burden for PDE simulation 042 by learning surrogate models to forward-simulate or predict spatiotemporal systems. However, this 043 new machine learning paradigm frequently exhibits instabilities during the training process, unstable 044 roll outs when modeling future state predictions, and often yields minimal computational speedups.

Shallow Recurrent Decoder networks (SHRED) (Williams et al.) 2024) are a recently introduced architecture that utilize data from sparse sensors to reconstruct and predict the entire spatiotem poral domain. Similar to Takens' embedding theorem, SHRED models trade spatial information at a single time point for a trajectory of sensor measurements across time. Previous work has shown SHRED can achieve excellent performance in examples ranging from weather forecast ing, atmospheric ozone concentration modeling, and turbulent flow reconstructions. In this paper, we introduce Sparse Identification of Nonlinear Dynamics with SHallow Recurrent Decoder networks (SINDy-SHRED). SINDy-SHRED exploits the latent space of recurrent neural networks for sparse sensor modeling, and enforces interprebility via a SINDy-based functional class. In this way, SINDy-SHRED enables a robust and sample-efficient joint discovery of the governing equation and

078

079 080 081



Figure 1: Illustration of the SINDy-SHRED architecture. SINDy-SHRED transfers the original sparse sensor signal (red) to an interpretable latent representation (purple) that falls into the SINDy-class functional. The shallow decoder performs a reconstruction in the pixel space.

082 083 coordinate system. With the correct governing equation, SINDy-SHRED can perform an accurate long-term prediction in a learned latent space, and in turn allow for long-term forecasting in the 084 pixel space. For practical applications, SINDy-SHRED is a lightweight model which can perform 085 low-rank recovery with only a few (e.g. three) active sensors, which is critical for large-scale scientific data modeling and real-time control. It does not require large amounts of data during training, 087 thereby avoiding a common pitfall in existing ML techniques for accelerating physics simulations. 088 SINDy-SHRED also exhibits remarkable training speed, even when executed on a single laptop. 089 Furthermore, SINDy-SHRED is highly reproducible with minimal effort in hyperparameter tuning. 090 The recommended network structure, hyperparameter, and training setting can generalize to many 091 different datasets. In short, we demonstrate SINDy-SHRED to be very robust and highly applicable in many modern scientific modeling problems. 092 093

Existing work seeking to perform data-driven, long-term forecasts of spatio-temporal phenomena typically suffers from (i) instabilities and (ii) massive computational requirements. We find that SINDy-SHRED ameliorates many of these issues because (i) it is based on a stable equation discovery and (ii) the learned model is an ODE in a learned latent space, rendering simulation more computationally efficient. We further conjecture that the fact SINDy-SHRED does not include a spatial encoder contributes to the architecture's robustness, rendering it more difficult for the model to overfit during training.

100 We perform a wide range of studies to demonstrate the effectiveness of SINDy-SHRED. We first 101 apply the model on the sea surface temperature data, which is a complex real-world problem. We 102 also consider data from a complex simulation of atmospheric chemistry, video data of flow over a 103 cylinder, and video data of a pendulum. The ability of SINDy-SHRED to perform well on video data is an important result for so-called "GoPro physics." With extremely small sample size and 104 105 noisy environments, SINDy-SHRED achieves governing equation identification with stable longterm predictions. Finally, we demonstrate the performance of SINDy-SHRED on a chaotic 2D 106 Kolmogorov flow (in Appendix  $\mathbf{D}$ ), finding a reasonable model even in the presence of chaos. The 107 contribution of our paper is three-fold.

111

112

113

- We propose SINDy-SHRED to incorporate symbolic understanding of the latent space of recurrent models of spatio-temporal dynamics.
- We further analyze the latent space of case studies and propose scientific models for these systems.
  - We systematically study SINDy-SHRED and compare to popular deep learning algorithms in spatio-temporal prediction.
- 114 115 116

117

# 2 RELATED WORKS

118 Traditionally, spatio-temporal physical phenomena are modeled by Partial Differential Equations 119 (PDEs). To accelerate PDE simulations, recent efforts have leveraged neural networks to model physics. By explicitly assuming the underlying PDE, physics-informed neural networks (Raissi 120 et al. 2019) utilize the PDE structure as a constraint for small sample learning. However, assuming 121 the exact form of governing PDE for real data can be a strong limitation. There have been many 122 recent works on learning and predicting PDEs directly using neural networks (Khoo et al., 2021; Li 123 et al., 2020b; Holl et al., 2020; Lu et al., 2021; Lin et al., 2021). Meanwhile, PDE-find (Rudy et al., 124 2017; Messenger & Bortz, 2021; Fasel et al., 2022) offers a data-driven approach to identify PDEs 125 from the spatial-temporal domain. Still, the high-dimensionality and required high data quality can 126 be prohibitive for practical applications. 127

In parallel, previous efforts in the discovery of physical law through dimensionality reduction tech-128 niques (Champion et al., 2019; Lusch et al., 2018; Mars Gao & Nathan Kutz, 2024) provide yet 129 another perspective on the modeling of scientific data. The discovery of physics from a learned 130 latent space has previously been explored by (Fukami et al., 2021; Cheng et al., 2024; Farenga et al., 131 2024; Conti et al., 2023; Wu et al., 2022; Li et al., 2020a), yet none of these methods consider a 132 regularization on the latent space with no explicit encoder. Yu et al. proposed the idea of physics-133 guided learning (Yu & Wang, 2024) which combines physics simulations and neural network ap-134 proximations. Directly modeling physics from video is also the subject of much research in the 135 field of robotics (Finn et al., 2016; Todorov et al., 2012; Sanchez-Gonzalez et al., 2018), computer 136 vision (Xie et al., 2024; Wu et al., 2017) and computer graphics (Kandukuri et al., 2020; Liu et al.; 137 Wu et al., 2015; Mrowca et al., 2018), since many fields of research require better physics models for simulation and control. From the deep learning side, combining the structure of differential 138 equations into neural networks (He et al., 2016; Chen et al., 2018) has been remarkably successful 139 in a wide range of tasks. When spatial-temporal modeling is framed as a video prediction problem, 140 He et al. found (He et al., 2022) that random masking can be an efficient spatio-temporal learner, 141 and deep neural networks can provide very good predictions for the next 10 to 20 frames (Shi et al., 142 2015; Wang et al., 2017; Gao et al., 2022; Guen & Thome, 2020). Generative models have also been 143 found to be useful for scientific data modeling (Mirza, 2014; Song et al., 2021; Cachay et al., 2024).

144 145

146

# 3 Methods

147 The shallow recurrent decoder network (SHRED) is a computational technique that utilizes recur-148 rent neural networks to predict the spatial domain. (Williams et al., 2024). The method functions 149 by trading high-fidelity spatial information for trajectories of sparse sensor measurements at given spatial locations. Mathematically, consider a high-dimensional data series  $\{X_i\}_{i=1}^T \in \mathbb{R}^{(W \times \tilde{H}) \otimes T}$ 150 151 that represents the evolution of a spatio-temporal dynamical system, W, H, and T denote the width, height, and total time steps of the system, respectively. In SHRED, each sensor collects data from a 152 fixed spatial position in a discretized time domain. Denote the subset of sensors as S, the input data 153 of SHRED is  $\{X^{S}\}_{i=1}^{T} \in \mathbb{R}^{\operatorname{card}(S)\otimes T}$ . Provided the underlying PDE allows spatial information to 154 propagate, these spatial effects will appear in the time history of the sensor measurements, enabling 155 the sensing of the entire field using only a few sensors. In vanilla SHRED, a Long Short-Term Mem-156 ory (LSTM) module is used to map the sparse sensor trajectory data into a latent space, followed by 157 a shallow decoder to reconstruct the entire spatio-temporal domain at the current time step. 158

SHRED enables efficient sparse sensing that is widely applicable to many scientific problems (Ebers
 et al., 2024; Kutz et al., 2024; Riva et al., 2024). The advantage of SHRED comes from three
 aspects. First, SHRED only requires minimal sensor measurements. Under practical constraints, collecting full-state measurements for data prediction and control can be prohibitively expensive.

Second, SHRED does not require grid-like data collection, which allows for generalization to more complex data structures. For example, it is easy to apply SHRED to graph data with an unknown underlying structure, such as human motion data on joints, robotic sensor data, and financial market data. Moreover, SHRED is theoretically rooted in PDE modeling methods from the perspective of separation of variables, which has the potential to offer strong theoretical guarantees such as convergence and stability.

168 169

## 3.1 EMPOWERING SHRED WITH REPRESENTATION LEARNING AND PHYSICS DISCOVERY

To achieve a parsimonious representation of physics, it is important to find a representation that effectively captures the underlying dynamics and structure of the system. In SINDy-SHRED (shown in Fig. ]], we extend the advantages of SHRED, and perform a joint discovery of coordinate systems and governing equations. This is accomplished by enforcing that the latent state of the recurrent network follows an ODE in the SINDy class of functions.

175

176 **Finding better representations** SHRED has a natural advantage in modeling the latent govern-177 ing physics due to its small model size. SHRED is based on a shallow decoder with a relatively 178 small recurrent network structure. The relative simplicity of the model allows the latent represen-179 tation to maintain many advantageous properties such as smoothness and Lipschitzness. Exper-180 imentally, we observe that the hidden state space of a SHRED model is generally very smooth. Second, SHRED does not have an explicit encoder, which avoids the potential problem of spec-181 tral bias (Rahaman et al., 2019). Many reduced-order modeling methods that rely on an encoder 182 architecture struggle to learn physics and instead focus only on modeling the low-frequency infor-183 mation (background) (Refinetti & Goldt, 2022; Champion et al., 2019; Mars Gao & Nathan Kutz, 184 2024). Building upon SHRED, we further incorporate SINDy to regularize the learned recurrence 185 with a well-characterized and simple form of governing equation. In other words, we perform a joint discovery of a coordinate system (which transfers the high-dimensional observation into a low-187 dimensional representation) and the governing law (which describes how the summarized latent 188 representation progresses forward with respect to time) of the latent space of a SHRED model. This 189 approach is inspired by the principle in physics that, under an ideal coordinate system, physical phe-190 nomena can be described by a parsimonious dynamical model (Champion et al., 2019; Mars Gao 191 & Nathan Kutz, 2024). When the latent representation and the governing law are well-aligned, this configuration is likely to capture the true underlying physics. This joint discovery results in a la-192 tent space that is both interpretable and physically meaningful, enabling robust and stable future 193 prediction based on the learned dynamics. 194

195 196

#### 3.2 SINDY-SHRED: LATENT SPACE REGULARIZATION VIA SINDY

197 As a compressive sensing procedure, there exist infinitely many equally valid solutions for the latent 198 representation. Therefore, it is not necessary for the latent representation induced by SHRED to 199 follow a well-structured differential equation. For instance, even if the exhibited dynamics are fun-200 damentally linear, the latent representation may exhibit completely unexplainable dynamics, making 201 the model challenging to interpret and extrapolate. Therefore, in SINDy-SHRED, our goal is to fur-202 ther constrain the latent representations to lie within the SINDy-class functional. This regularization promotes models that are fundamentally explainable by a SINDy-based ODE, allowing us to iden-203 tify a parsimonious governing equation. The SINDy class of functions typically consists of a library 204 of commonly used functions, which includes polynomials, and Fourier series. Although they may 205 seem simple, these functions possess surprisingly strong expressive power, enabling the model to 206 capture very complex dynamical systems. 207

SINDy as a Recurrent Neural network We first reformulate SINDy using a neural network form, simplifying its incorporation into a SHRED model. ResNet (He et al., 2016) and Neural ODE (Chen et al., 2018) utilize skip connections to model residual and temporal derivatives. Similarly, this could also be done via a Recurrent Neural Network (RNN) which has a general form of

$$z_{t+1} = z_t + f(x_t), (1)$$

where  $f(\cdot)$  is some function of the input. From the Euler method, the ODE forward simulation via SINDy effectively falls into the category of Recurrent Neural Networks (RNNs) which has the form

$$z_{t+1} = z_t + f_{\Theta}(x_t, \Xi, \Delta t), \tag{2}$$

216 where  $f_{\Theta}(x_t, \Xi, \Delta t) = \Theta(x_t) \Xi \Delta t$  is a nonlinear func-217 tion. Notice that this  $f_{\Theta}(\cdot)$  has exactly the same formula-218 tion as in SINDy (Brunton et al., 2016). The application 219 of function libraries with sparsity constraints is a manner 220 of automatic neural architecture search (NAS) (Zoph & Le, 2016). Compared to all prior works (Champion et al., 221 2019; Fukami et al., 2021; Conti et al., 2023), this imple-222 mentation of the SINDy unit fits better in the framework of neural network training and gradient descent. We utilize 224 trajectory data  $\{\mathbf{z}_i\}_{i=1}^T$  and forward simulate the SINDy-225 based ODE using a trainable parameter  $\Xi$ . To achieve bet-226 ter stability and accuracy for forward integration, we use 227 Euler integration with k mini-steps (with time step  $\frac{\Delta t}{k}$ ) to 228 obtain  $\mathbf{z}_{t+1}$ . In summary, defining  $h = \frac{\Delta t}{k}$ , we optimize 229  $\Xi$  with the following: 230



Figure 2: Diagram of the RNN form of SINDy.

$$\Xi = \arg\min\left\|\mathbf{z}_{t+1} - \left(\mathbf{z}_t + \sum_{i=0}^{k-1} \Theta(\mathbf{z}_{t+ih}) \Xi h\right)\right\|_2^2, \quad \mathbf{z}_{t+ih} = \mathbf{z}_t + \Theta(\mathbf{z}_{t+(i-1)h}) \Xi h, \quad \min\|\Xi\|_0.$$
(3)

To achieve  $\ell_0$  optimization, we perform pruning with  $\ell_2$  which approximates  $\ell_0$  regularization under regularity conditions (Zheng et al., 2014; Gao et al., 2023; Blalock et al., 2020). Applying SINDy unit has the following benefits: (a) The SINDy-function library contains frequently used functions in physics modeling (e.g. polynomials and Fourier series). (b) With sparse system identification, the neural network is more likely to identify governing physics, which is fundamentally important for extrapolation and long-term stability.

**Latent space regularization via ensemble SINDy** We first note that we deviate from the original 242 SHRED architecture by using a GRU as opposed to an LSTM. This choice was made because we 243 generally found that GRU provides a smoother latent space. Now, recall that our goal is to find a 244 SHRED model with a latent state that is within the SINDy-class functional. However, the initial 245 latent representation from SHRED does not follow the SINDy-based ODE structure at all. On the 246 one hand, if we naively apply SINDy to the initial latent representation, the discovery is unlikely 247 to fit the latent representation trajectory. On the other hand, if we directly replace the GRU unit to 248 SINDy and force the latent space to follow the discovered SINDy model, it might lose information 249 that is important to reconstruction the entire spatial domain. Therefore, it is important to let the two 250 latent spaces align progressively.

251 In Algorithm I, we describe our training procedure that allows the two trajectories to progressively 252 align with each other. To further ensure a gradual adaptation and avoid over-regularization, we intro-253 duce ensemble SINDy units with varying levels of sparsity constraints, which ranges the effect from 254 promoting a full model (all terms in the library are active) to a null model (where no dynamics are represented). From the initial latent representation  $z_{1:t}^{iter 0}$  from SHRED, the SINDy model first pro-255 256 vides an initial estimate of ensemble SINDy coefficients  $\{\hat{\Xi}_0^i\}_{i=b}^B$ . Then, the parameters of SHRED 257 will be updated towards the dynamics simulated by  $\{\hat{\Xi}_0\}_{i=b}^B$ , which generates a new latent repre-258 sentation trajectory  $\mathbf{z}_{1:t}^{\text{iter 1}}$ . We iterate this procedure and jointly optimize the following loss function 259 to let the SHRED latent representation trajectory approximate the SINDy generated trajectory: 260

$$\mathcal{L} = \left\| \mathbf{X}_{t} - f_{\theta_{D}}(f_{\theta_{\text{GRU}}}(\mathbf{X}_{t-L:t}^{S})) \right\|_{2}^{2} + \sum_{i=1}^{B} \left\| \mathbf{Z}_{t}^{\text{GRU}} - \left( \mathbf{Z}_{t-1}^{\text{GRU}} + \sum_{j=0}^{k-1} \Theta(\mathbf{Z}_{t-1+jh}) \Xi^{(i)}h \right) \right\|_{2}^{2} + \lambda \left\| \Xi \right\|_{0}$$
(4)

267

268

269

261 262

241

where  $Z_{t-1+ih} = Z_t + \Theta(Z_{t-1+(i-1)h}) \Xi h$ ,  $Z_{t-1} = Z_{t-1}^{\text{GRU}}$ , and  $h = \frac{\Delta t}{k}$ .

4 EXPERIMENT

In the following, we perform case studies across a range of scientific and engineering problems.

289

290

291

292

293

295

296

297

298

299

300

301

302

303 304

305 306 307

311

315

316 317

270 Algorithm 1 Latent state space regularization via SINDy 271 **Input:** input  $X_{t-L:t+1}^{S}$ ,  $X_t$ , SINDy library  $\Theta(\cdot)$ , timestep  $\Delta t$ . 272 1: function LATENTSPACESINDY( $X_{t-L:t+1}^{S}, X_{t+1}, \Delta t$ ) 273 for i in  $0, 1, \dots, n-1$ : do 2: 274  $\begin{aligned} \mathbf{Z}_{t}, \ \mathbf{Z}_{t+1} &= f_{\theta_{\text{GRU}}}(\mathbf{X}_{t-L:t}^{\mathcal{S}}), \ f_{\theta_{\text{GRU}}}(\mathbf{X}_{t-L+1:t+1}^{\mathcal{S}}); \\ \text{for j in } (0, 1, \frac{\Delta t}{k}): \text{do} \\ \mathbf{Z}_{t+\frac{j+1}{k}\Delta t}^{\text{SINDy}} &= \mathbf{Z}_{t+\frac{j}{k}\Delta t}^{\text{SINDy}} + \Theta(\mathbf{Z}_{t+\frac{j}{k}\Delta t}^{\text{SINDy}}) \Xi \Delta t \\ \text{end for} \end{aligned}$ 3: 275 4: ▷ SINDy forward simulation 276 5: 277 278 end for 6:  $\hat{\boldsymbol{X}}_{t+1} = f_{\theta_D}(\boldsymbol{Z}_{t+1})$ 279 7: ▷ SHRED reconstruction  $\theta_{\text{GRU}}, \Xi, \theta_D = \arg\min_{\theta_{\text{GRU}},\Xi,\theta_D} \left\| \mathbf{X}_{t+1} - \hat{\mathbf{X}}_{t+1} \right\|_2^2 + \left\| \mathbf{Z}_{t+1}^{\text{GRU}} - \mathbf{Z}_{t+1}^{\text{SINDy}} \right\|_2^2 + \lambda \left\| \Xi \right\|_0$ 8: 281 9: **if**  $i \mod 100 = 0$  **then** 282  $\Xi[|\Xi| < \text{threshold}] = 0$ 10: 283 end if 11: 284 12: end for ▷ Train until converges 13: end function

**Sea-surface temperature** The first example we consider is that of global sea-surface temperature. The SST data contains 1,400 weekly snapshots of the weekly mean sea surface temperature from 1992 to 2019 reported by NOAA (Reynolds et al., 2002). The data is represented by a  $180 \times 360$ grid, of which 44,219 grid points correspond to sea-surface locations. We standardize the data with its own min and max, which transforms the sensor measurements to within the numerical range of (0, 1). We randomly select 250 sensors from the possible 44,219 locations and set the lag parameter to 52 weeks. The inclusion of 250 sensors is a substantial deviation from previous work with SHRED in which far fewer sensors were used Williams et al. (2024). However, we found greater robustness in the application of E-SINDy to the learned latent state when more sensors were utilized. Thus, for each input-output pair, the input consists of the 52-week trajectories of the selected sensors, while the output is a single temperature field across all 44,219 spatial locations. SINDy-SHRED aims to reconstruct the entire sea surface temperature locations from these randomly selected sparse sensor trajectories. We include the details of the experimental settings of SINDy-SHRED in the Appendix C.1. From the discovered coordinate system, we define the representation of physics - the latent hidden state space - to be  $(z_1, z_2, z_3)$ . The dynamics progresses forward via the following set of equations:

$$\dot{z}_1 = 4.68z_2 - 2.37z_3, 
\dot{z}_2 = -3.10z_1 + 3.25z_3, 
\dot{z}_3 = 2.72z_1 - 5.55z_2.$$
(5)



Figure 3: Extrapolation of latent representation in SINDy-SHRED from the discovered dynamical system for SST. Colored: true latent representation. Grey: SINDy extrapolation.

318 The discovery of a linear system describing the evolution of the latent state is in line with prior 319 work on SST data De Bézenac et al. (2019) in which it as assumed that the underlying physics is an 320 advection-diffusion PDE. In Fig. 3 (a) we further present the accuracy of this discovered system by 321 forward simulating the system from an initial condition for a total of 27 years (c.f. Fig. 15). It is observable how the discovered law is close to the true evolution of latent hidden states and, critically, 322 there appears to be minimal phase slipping. Extrapolating the latent state space via forward inte-323 gration, we can apply the shallow decoder to return forecasts of the high-dimensional data. Doing

6

Real Data SINDy-SHRED Prediction 0 weeks 25 weeks 50 weeks 75 weeks 100 weeks

Figure 4: Long-term global sea-surface temperature prediction via SINDy-SHRED from week 0 to week 100. We crop the global temperature map for better visualization.

so, we find an averaged MSE error of  $0.57 \pm 0.10^{\circ}C$  for all prediction lengths in the test dataset. In Fig. 4, we show SINDy-SHRED produces stable long-term predictions for SST data. We further include Fig. 16 to demonstrate the extrapolation of each sensor. The sensor level prediction is based on the global prediction of the future frame, and we visualize the signal trajectory of specific sensor locations. We find SINDy-SHRED is robust for out-of-distribution sensors, though its extrapolation may not accurately capture anomalous events.

350 **3D** Atmospheric ozone concentration The atmospheric ozone concentration dataset (Bey et al., 2001) contains a one-year simulation of the evolution of an ensemble of interacting chemical species 352 through a transport operator using GEOS-Chem. The simulation contains 1,456 temporal samples 353 with a timestep of 6 hours over one year for 99,360 (46 by 72 by 30) spatial locations (latitude, 354 longitude, elevation). The data presented in this work has compressed by performing an SVD and 355 retaining only the first 50 POD modes. As with the SST data, we standardize the data within the 356 range of (0, 1) and randomly select and fix 3 sensors out of 99,360 spatial locations (0.5%). We include the details of the experimental settings of SINDy-SHRED in the Appendix C.2. The converged 357 latent representation presents the following SINDy model: 358

$$\begin{cases} \dot{z}_1 = -0.002 - 0.013z_2 + 0.007z_3, \\ \dot{z}_2 = -0.001z_1 + 0.004z_2 - 0.008z_3, \\ \dot{z}_3 = 0.002 + 0.012z_2 - 0.005z_3. \end{cases}$$
(6)

The identified governing physics is close to a linear system with constant terms for damping. Unlike 363 traditional architectures for similar problems, which may include expensive 3D convolution, SINDy-364 SHRED provides an efficient way of training, taking about half an hour. Although the quantity 365 of data is insufficient to perform long term-predictions, SINDy-SHRED still exhibits interesting 366 behavior for a longer-term extrapolation which converges to the fixed point at 0 (as shown in Fig. 17). 367 From the extrapolation of the latent state space, the shallow decoder prediction has an averaged  $\overline{MSE}$ 368 error of  $1.5e^{-2}$ . In Fig. 6, we visualize the shallow decoder prediction up to 14 weeks. In Fig. 18, 369 we reconstruct the sensor-level predictions which demonstrate the details of the signal prediction. 370 The observations are much noiser than the SST data, but SINDy-SHRED provides a smoothed 371 extrapolation for the governing trends.

372

324

330 331

339 340

341

342 343 344

345

346

347

348

349

351

373 **GoPro physics data: flow over a cylinder** In this subsection, we demonstrate the performance 374 of SINDy-SHRED on an example of so-called "GoPro physics modeling." The considered data 375 is collected from a dyed water channel to visualize a flow over a cylinder (Albright, 2017). The 376 Reynolds number is 171 in the experiment. The dataset contains 11 seconds of video taken at 30 frames per second (FPS). We manually perform data augmentation and repeat the latter part of the 377 video once to increase the number of available training samples . We transfer the original RGB



Figure 5: Extrapolation of latent representation in SINDy-SHRED from the discovered dynamical system for Ozone data. Colored: true latent representation. Grey: SINDy extrapolation.

Real DataImage: Prediction of the second second

Figure 6: Long-term global Ozone data prediction via SINDy-SHRED with elevation 0 from week 0 to week 14.

404 405

386

387

388 389

397

channel to gray scale and remove the background by subtracting the mean of all frames. After the prior processing step, the video data has only one channel (gray) within the range (0, 1) with a height of 400 pixels and a width of 1,000 pixels. We randomly select and fix 200 pixels as sensor measurements from the entire 400,000 space, which is equivalent to only 0.05% of the data. We set the lag parameter to 60 frames. We include the details of the experimental settings of SINDy-SHRED in the Appendix C.4. We define the representation of the hidden latent state space as  $(z_1, z_2, z_3, z_4)$ . We discover the following dynamical system:

$$\begin{cases} \dot{z}_1 = -0.69z_2 + 0.98z_3 - 0.40z_4, \\ \dot{z}_2 = 1.00z_1 - 0.78z_31 - 0.31z_2z_3^2, \\ \dot{z}_3 = -1.029z_1 + 0.59z_2 + 0.41z_4, \\ \dot{z}_4 = -0.26z_1^2 - 0.29z_2^2z_3 - 0.39z_3^3. \end{cases}$$
(7)

420 Compared to the systems discovered in all previous examples, the flow over a cylinder model is much 421 more complex with significant nonlinear interactions. In Eqn. 7 we find that  $z_1$  and  $z_3$  behave like 422 a governing mode of the turbulence swing;  $z_2$  and  $z_4$  further depict more detailed nonlinear effects. 423 We also present the result of extrapolating this learned representation. We generate the trajectory 424 from the initial condition at time point 0 and perform forward integration for extrapolation. As 425 shown in Fig. 7 the learned ODE closely follows the dynamics of  $z_1$  and  $z_3$  up to 7 seconds (210 426 timesteps);  $z_2$  and  $z_4$  also have close extrapolation up to 3 seconds.

This learned representation also nicely predicts the future frames in pixel space. The shallow decoder
prediction has an averaged MSE error of 0.030 (equivalently 3%) over the entire available trajectory.
In Fig. 8 we observe that the autoregressively generated prediction frames closely follow the true data, and further in Fig. 21 we find that the predictions are still stable after 1,000 frames, which is out of the size of the original dataset. The sensor-level prediction in Fig. 20 further demonstrates the accuracy of reconstruction in detail.



Figure 7: Extrapolation of latent representation in SINDy-SHRED from the discovered dynamical system for flow over a cylinder data. Colored: true latent representation. Grey: SINDy extrapolation.



Figure 8: Long-term pixel space video prediction via SINDy-SHRED. We demonstrate the forward prediction outcome up to 180 frames.

**Baseline study: prediction of single shot real pendulum recording** In this subsection, we compare the performance of SINDy-SHRED to other popular existing learning algorithms. We perform the baseline study particularly on video data of a pendulum since many deep learning algorithms are hard to scale up to deal with large scientific data. In the following, we demonstrate the result of video prediction on the pendulum data using ResNet (He et al., 2016), convolutional LSTM (con-vLSTM) (Shi et al., 2015), and PredRNN (Wang et al., 2017), and SimVP (Gao et al., 2022). The pendulum in our experiment is not ideal and includes complex damping effects. We use a nail on the wall and place the rod (with a hole) on the nail. This creates complex friction, which slows the rod more when passing the lowest point due to the increased pressure caused by gravity. The full model we discovered from the video (as shown in Fig. ) includes four terms:

$$\ddot{z} = 0.17\dot{z}^2 - 0.06\dot{z}^3 - 10.87\sin(z) + 0.48\sin(\dot{z}).$$
(8)

As shown in Table I. SINDy-SHRED outperforms all baseline methods for total error and long-term predictions. Generally, all baseline deep learning methods perform well for short-term forecasting, but the error quickly accumulates for longer-term predictions. This is also observable from the pre-diction in the pixel space as shown in Fig. 10. SINDy-SHRED is the only method that does not produce collapsed longer-term predictions. In Fig. 22, the sensor level prediction also demonstrates the robustness of the SINDy-SHRED prediction. PredRNN is the second best method as measured by the total error. However, PredRNN is expensive in computation which includes a complex for-ward pass with an increased number of parameters. It is also notable that the prediction of PredRNN collapses after 120 frames, after which only an averaged frame over the entire trajectory is pre-dicted. ConvLSTM has a relatively better result in terms of generation, but the long-term prediction is still inferior compared to SINDy-SHRED. Additionally, we note that 2D convolution is much more computationally expensive. For larger spatiotemporal domains (e.g. the SST example and 3D ozone data), the computational complexity of convolution will scale up very quickly, which makes the algorithm impractical to execute. Similar computational issues will occur for diffusion models and generative models, which is likely to be impractical to compute, and unstable for longer-term predictions. In summary, we observe that SINDy-SHRED is not only a more accurate long-term model, but is also faster to execute and smaller in size.



Figure 10: The pendulum video generation outcome from ResNet, SimVP, ConvLSTM, PredRNN, and SINDy-SHRED from frame 20 to frame 245.

Models	Params #	Training time	T = [0, 100]	T = [100, 200]	T = [200, 275]	Total
ResNet (He et al., 2016)	2.7M	24 mins	$2.08 \times 10^{-2}$	$1.88 \times 10^{-2}$	$2.05 \times 10^{-2}$	$2.00 \times 10^{-2}$
SimVP (Gao et al., 2022)	460K	30 mins	$2.29 \times 10^{-2}$	$2.47 \times 10^{-2}$	$2.83 \times 10^{-2}$	$2.53 \times 10^{-2}$
PredRNN (Wang et al., 2017)	444K	178 mins	$1.02 \times 10^{-2}$	$1.79 \times 10^{-2}$	$1.69 \times 10^{-2}$	$1.48 \times 10^{-2}$
ConvLSTM (Shi et al., 2015)	260K	100 mins	$9.24\times\mathbf{10^{-3}}$	$1.86 \times 10^{-2}$	$1.99 \times 10^{-2}$	$1.55 \times 10^{-2}$
SINDy-SHRED*	44K	17 mins	$1.70 \times 10^{-2}$	$9.36\times\mathbf{10^{-3}}$	$5.31\times\mathbf{10^{-3}}$	$1.05\times 10^{-2}$

Table 1: Comparison table of SINDy-SHRED to baseline methods for parameter size, training time, and mean-squared error over different prediction horizon.

# 5 CONCLUSION

In this paper, we present SINDy-SHRED, which jointly performs the discovery of coordinate systems and governing equations with low computational cost and strong predictive power. Through experiments, we show that our method can produce robust and accurate long-term predictions for a variety of complex problems, including global sea-surface temperature, 3D atmospheric ozone concentration, flow over a cylinder, and a moving pendulum. SINDy-SHRED achieves state-of-the-art performance in long-term autoregressive video prediction, outperforming ConvLSTM, PredRNN, ResNet, and SimVP with the lowest computational cost and training time.

# 540 REFERENCES

549

550

551

556

559

560

561

562

566

567

568

569

573

585

592

542 Ja	cob Albright. Flow	visualization in	a water channel, 20	017. URL	https://	www.youtube.
543	com/watch?v=30	_aADFVL9M.	YouTube video.			

- Andrew R Barron. Universal approximation bounds for superpositions of a sigmoidal function. *IEEE Transactions on Information theory*, 39(3):930–945, 1993.
- Peter L Bartlett and Shahar Mendelson. Rademacher and gaussian complexities: Risk bounds and
   structural results. *Journal of Machine Learning Research*, 3(Nov):463–482, 2002.
  - Peter L Bartlett, Sanjeev R Kulkarni, and S Eli Posner. Covering numbers for real-valued function classes. *IEEE transactions on information theory*, 43(5):1721–1724, 1997.
- Isabelle Bey, Daniel J Jacob, Robert M Yantosca, Jennifer A Logan, Brendan D Field, Arlene M Fiore, Qinbin Li, Honguy Y Liu, Loretta J Mickley, and Martin G Schultz. Global modeling of tropospheric chemistry with assimilated meteorology: Model description and evaluation. *Journal of Geophysical Research: Atmospheres*, 106(D19):23073–23095, 2001.
- 557 Davis Blalock, Jose Javier Gonzalez Ortiz, Jonathan Frankle, and John Guttag. What is the state of 558 neural network pruning? *Proceedings of machine learning and systems*, 2:129–146, 2020.
  - Steven L Brunton, Joshua L Proctor, and J Nathan Kutz. Discovering governing equations from data by sparse identification of nonlinear dynamical systems. *Proceedings of the national academy of sciences*, 113(15):3932–3937, 2016.
- Salva Rühling Cachay, Brian Henn, Oliver Watt-Meyer, Christopher S Bretherton, and Rose Yu.
   Probabilistic emulation of a global climate model with spherical dyffusion. *arXiv preprint arXiv:2406.14798*, 2024.
  - Claudio Canuto, M Yousuff Hussaini, Alfio Quarteroni, and Thomas A Zang. *Spectral methods: evolution to complex geometries and applications to fluid dynamics*. Springer Science & Business Media, 2007.
- Kathleen Champion, Bethany Lusch, J Nathan Kutz, and Steven L Brunton. Data-driven discovery of coordinates and governing equations. *Proceedings of the National Academy of Sciences*, 116 (45):22445–22451, 2019.
- Ricky TQ Chen, Yulia Rubanova, Jesse Bettencourt, and David K Duvenaud. Neural ordinary differential equations. *Advances in neural information processing systems*, 31, 2018.
- Xiangning Chen, Chen Liang, Da Huang, Esteban Real, Kaiyuan Wang, Hieu Pham, Xuanyi Dong,
   Thang Luong, Cho-Jui Hsieh, Yifeng Lu, et al. Symbolic discovery of optimization algorithms.
   Advances in neural information processing systems, 36, 2024.
- Sheng Cheng, Deqian Kong, Jianwen Xie, Kookjin Lee, Ying Nian Wu, and Yezhou Yang. Latent space energy-based neural odes. *arXiv preprint arXiv:2409.03845*, 2024.
- Paolo Conti, Giorgio Gobat, Stefania Fresca, Andrea Manzoni, and Attilio Frangi. Reduced order
   modeling of parametrized systems through autoencoders and sindy approach: continuation of
   periodic solutions. *Computer Methods in Applied Mechanics and Engineering*, 411:116072, 2023.
- Emmanuel De Bézenac, Arthur Pajot, and Patrick Gallinari. Deep learning for physical processes:
   Incorporating prior scientific knowledge. *Journal of Statistical Mechanics: Theory and Experiment*, 2019(12):124009, 2019.
- Megan R Ebers, Jan P Williams, Katherine M Steele, and J Nathan Kutz. Leveraging arbitrary
   mobile sensor trajectories with shallow recurrent decoder networks for full-state reconstruction.
   *IEEE Access*, 2024.
- 593 Nicola Farenga, Stefania Fresca, Simone Brivio, and Andrea Manzoni. On latent dynamics learning in nonlinear reduced order modeling. *arXiv preprint arXiv:2408.15183*, 2024.

- 594 Urban Fasel, J Nathan Kutz, Bingni W Brunton, and Steven L Brunton. Ensemble-sindy: Robust 595 sparse model discovery in the low-data, high-noise limit, with active learning and control. Pro-596 ceedings of the Royal Society A, 478(2260):20210904, 2022. 597 Chelsea Finn, Ian Goodfellow, and Sergey Levine. Unsupervised learning for physical interaction 598 through video prediction. Advances in neural information processing systems, 29, 2016. 600 Kai Fukami, Takaaki Murata, Kai Zhang, and Koji Fukagata. Sparse identification of nonlinear 601 dynamics with low-dimensionalized flow representations. Journal of Fluid Mechanics, 926:A10, 602 2021. 603 L Gao, Urban Fasel, Steven L Brunton, and J Nathan Kutz. Convergence of uncertainty estimates in 604 ensemble and bayesian sparse model discovery. arXiv preprint arXiv:2301.12649, 2023. 605 606 Zhangyang Gao, Cheng Tan, Lirong Wu, and Stan Z Li. Simvp: Simpler yet better video prediction. 607 In Proceedings of the IEEE/CVF conference on computer vision and pattern recognition, pp. 608 3170-3180, 2022. 609 Vincent Le Guen and Nicolas Thome. Disentangling physical dynamics from unknown factors for 610 unsupervised video prediction. In Proceedings of the IEEE/CVF conference on computer vision 611 and pattern recognition, pp. 11474–11484, 2020. 612 613 Kaiming He, Xiangyu Zhang, Shaoqing Ren, and Jian Sun. Deep residual learning for image recognition. In Proceedings of the IEEE conference on computer vision and pattern recognition, pp. 614 770–778, 2016. 615 616 Kaiming He, Xinlei Chen, Saining Xie, Yanghao Li, Piotr Dollár, and Ross Girshick. Masked au-617 to encoders are scalable vision learners. In Proceedings of the IEEE/CVF Conference on Computer 618 Vision and Pattern Recognition, pp. 16000–16009, 2022. 619 Philipp Holl, Vladlen Koltun, and Nils Thuerey. Learning to control pdes with differentiable physics. 620 arXiv preprint arXiv:2001.07457, 2020. 621 622 Rama Kandukuri, Jan Achterhold, Michael Moeller, and Joerg Stueckler. Learning to identify phys-623 ical parameters from video using differentiable physics. In DAGM German conference on pattern 624 recognition, pp. 44-57. Springer, 2020. 625 Yuehaw Khoo, Jianfeng Lu, and Lexing Ying. Solving parametric pde problems with artificial neural 626 networks. European Journal of Applied Mathematics, 32(3):421–435, 2021. 627 628 J Nathan Kutz, Maryam Reza, Farbod Faraji, and Aaron Knoll. Shallow recurrent decoder for 629 reduced order modeling of plasma dynamics. arXiv preprint arXiv:2405.11955, 2024. 630 Yunzhu Li, Toru Lin, Kexin Yi, Daniel Bear, Daniel Yamins, Jiajun Wu, Joshua Tenenbaum, and 631 Antonio Torralba. Visual grounding of learned physical models. In International conference on 632 machine learning, pp. 5927-5936. PMLR, 2020a. 633 634 Zongyi Li, Nikola Kovachki, Kamyar Azizzadenesheli, Burigede Liu, Kaushik Bhattacharya, An-635 drew Stuart, and Anima Anandkumar. Neural operator: Graph kernel network for partial differential equations. arXiv preprint arXiv:2003.03485, 2020b. 636 637 Guang Lin, Christian Moya, and Zecheng Zhang. Accelerated replica exchange stochastic gradient 638 langevin diffusion enhanced bayesian deeponet for solving noisy parametric pdes. arXiv preprint 639 arXiv:2111.02484, 2021. 640 Shaowei Liu, Zhongzheng Ren, Saurabh Gupta, and Shenlong Wang. Physgen: Rigid-body physics-641 grounded image-to-video generation. 642 643 Lu Lu, Pengzhan Jin, Guofei Pang, Zhongqiang Zhang, and George Em Karniadakis. Learning 644 nonlinear operators via deeponet based on the universal approximation theorem of operators. 645 *Nature machine intelligence*, 3(3):218–229, 2021. 646 Bethany Lusch, J Nathan Kutz, and Steven L Brunton. Deep learning for universal linear embeddings
- 647 Bethany Lusch, J Nathan Kutz, and Steven L Brunton. Deep learning for universal linear embeddings of nonlinear dynamics. *Nature communications*, 9(1):1–10, 2018.

648 649 650	L Mars Gao and J Nathan Kutz. Bayesian autoencoders for data-driven discovery of coordinates, governing equations and fundamental constants. <i>Proceedings of the Royal Society A</i> , 480(2286): 20230506, 2024.
651	
652	Daniel A Messenger and David M Bortz. Weak sindy for partial differential equations. <i>Journal of</i>
653	<i>Computational Physics</i> , 445:110525, 2021.
654 655	Mehdi Mirza. Conditional generative adversarial nets. arXiv preprint arXiv:1411.1784, 2014.
656	Damian Mrowca, Chengxu Zhuang, Elias Wang, Nick Haber, Li F Fei-Fei, Josh Tenenbaum, and
657	Daniel L Yamins. Flexible neural representation for physics prediction. Advances in neural
000	information processing systems, 51, 2018.
660	Nasim Rahaman, Aristide Baratin, Devansh Arpit, Felix Draxler, Min Lin, Fred Hamprecht, Yoshua
661	Bengio, and Aaron Courville. On the spectral bias of neural networks. In <i>International conference</i> on machine learning, pp. 5301–5310. PMLR, 2019.
662	
664 665	Maziar Raissi, Paris Perdikaris, and George E Karniadakis. Physics-informed neural networks: A deep learning framework for solving forward and inverse problems involving nonlinear partial differential equations. <i>Journal of Computational physics</i> , 378:686–707, 2019.
666	
667	autoencoders In International Conference on Machine Learning pp 18/00–18510 DMI P
668	2022
669	
670	Richard W Reynolds, Nick A Rayner, Thomas M Smith, Diane C Stokes, and Wanqiu Wang. An
671	improved in situ and satellite sst analysis for climate. <i>Journal of climate</i> , 15(13):1609–1625,
672	2002.
673	Stefano Riva, Carolina Introini, Antonio Cammi, and J Nathan Kutz. Robust state estimation from
675	partial out-core measurements with shallow recurrent decoder for nuclear reactors. arXiv preprint
676	arXiv:2409.12550, 2024.
677	Constitute a Constitute to the tensor attraction Rectification of the
678	partial differential equations. <i>Science advances</i> , 3(4):e1602614, 2017.
679 680 681 682	Alvaro Sanchez-Gonzalez, Nicolas Heess, Jost Tobias Springenberg, Josh Merel, Martin Riedmiller, Raia Hadsell, and Peter Battaglia. Graph networks as learnable physics engines for inference and control. In <i>International conference on machine learning</i> , pp. 4470–4479. PMLR, 2018.
683	Xingjian Shi, Zhourong Chen, Hao Wang, Dit-Yan Yeung, Wai-Kin Wong, and Wang-chun Woo
684 685	Convolutional lstm network: A machine learning approach for precipitation nowcasting. Ad- vances in neural information processing systems 28, 2015
686	vances in neural information processing systems, 26, 2015.
687	Yang Song, Conor Durkan, Iain Murray, and Stefano Ermon. Maximum likelihood training of
688 689	score-based diffusion models. Advances in neural information processing systems, 34:1415–1428, 2021.
690	
691	Emanuel Iodorov, Iom Erez, and Yuval Iassa. Mujoco: A physics engine for model-based control.
692	IEEE 2012
693	
694	Yunbo Wang, Mingsheng Long, Jianmin Wang, Zhifeng Gao, and Philip S Yu. Predrnn: Recurrent
695 696	neural networks for predictive learning using spatiotemporal lstms. <i>Advances in neural informa-</i> <i>tion processing systems</i> , 30, 2017.
697	Ion D. Williama Olivia John and I. Nothen Kutz, Sancing with shallow as suggested as a start of the
698 699	Proceedings of the Royal Society A, 480(2298):20240054, 2024.
700 701	Jiajun Wu, Ilker Yildirim, Joseph J Lim, Bill Freeman, and Josh Tenenbaum. Galileo: Perceiving physical object properties by integrating a physics engine with deep learning. <i>Advances in neural information processing systems</i> , 28, 2015.

702 703 704	Jiajun Wu, Erika Lu, Pushmeet Kohli, Bill Freeman, and Josh Tenenbaum. Learning to see physics via visual de-animation. <i>Advances in neural information processing systems</i> , 30, 2017.
705 706 707	Tailin Wu, Takashi Maruyama, and Jure Leskovec. Learning to accelerate partial differential equa- tions via latent global evolution. <i>Advances in Neural Information Processing Systems</i> , 35:2240– 2253, 2022.
708 709 710	Tianyi Xie, Zeshun Zong, Yuxing Qiu, Xuan Li, Yutao Feng, Yin Yang, and Chenfanfu Jiang. Physgaussian: Physics-integrated 3d gaussians for generative dynamics. In <i>Proceedings of the</i> <i>IEEE/CVF Conference on Computer Vision and Pattern Recognition</i> , pp. 4389–4398, 2024.
711 712 713	Rose Yu and Rui Wang. Learning dynamical systems from data: An introduction to physics-guided deep learning. <i>Proceedings of the National Academy of Sciences</i> , 121(27):e2311808121, 2024.
714 715 716	Zemin Zheng, Yingying Fan, and Jinchi Lv. High dimensional thresholded regression and shrinkage effect. <i>Journal of the Royal Statistical Society Series B: Statistical Methodology</i> , 76(3):627–649, 2014.
717 718 719	Barret Zoph and Quoc V Le. Neural architecture search with reinforcement learning. <i>arXiv preprint arXiv:1611.01578</i> , 2016.
720	
721	
723	
724	
725	
726	
727	
728	
729	
730	
731	
732	
733	
734	
735	
736	
737	
738	
739	
740	
741	
742	
743	
744	
746	
747	
748	
749	
750	
751	
752	
753	
754	
755	