
Neural Operators as Fast Surrogate Models for the Transmission Loss of Parameterized Sonic Crystals

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Abstract

Neural operators serve as efficient, data-driven surrogate models for complex physical and engineering problems. In this work, we demonstrate that neural operators can directly learn the key properties of sonic crystals, a type of acoustic metamaterial consisting of a lattice of parameterized shapes. We predict the transmission loss curve, a critical characteristic in applications, bypassing the expensive meshing and solving steps typical of classical techniques. We evaluate established architectures, DeepONet (DON) and Fourier Neural Operator (FNO), alongside two new ones, Deep Neural Operator (DNO) and Deep Cat Operator (DCO), which demonstrate significant performance improvements. In our experiments, all models achieve high accuracy, while being up to 10^6 times faster than the traditional method, significantly advancing practical real-time metamaterial design.

1 Introduction

Metamaterials are composite materials with behavior not found in nature, whose properties arise from their internal structure rather than their composition [1]. They can be designed to have, e.g., new elastic, electromagnetic, or thermal properties [2, 3, 4], and have numerous applications including advanced sensors, antenna design, energy harvesting, and civil engineering [5, 6, 7, 8], to name a few. In acoustics, they are mostly concerned with sound attenuation [9, 10], wave manipulation [11, 12, 13], or architectural acoustics [7]. Efficient and accurate simulation methods are essential in order to predict the performance of acoustic metamaterials *in silico*. State-of-the-art approaches [14, 15, 16, 17, 18], while sufficient for this purpose, can be computationally expensive, requiring a mesh of the domain and solving optimization problems for a large number of parameters.

Sonic crystals are acoustic metamaterials composed of identical unit cells organized in a lattice, cf. Figure 1. The composition creates so-called band-gaps in the transmission loss graph, where the transmission loss (TL) describes the decrease in power from the incident to the transmitted wave [9, 19]. The graph of the TL over the frequency is a key characteristic, describing the interaction of the crystal with the environment, cf. Figure 1 (right). The primary objective in designing these metamaterials is to predict and achieve a sonic crystal with a specific TL profile.

Neural operators (NOs) [20] and related techniques [21, 22, 23, 24] have been successfully applied to model complex physical systems, particularly to approximate the solution operators of partial differential equations. They are accurate enough to complement or even replace traditional numerical simulators in areas like computational fluid dynamics, weather forecasting, and material modeling, while being orders of magnitude faster [25, 26].

In this work, we train NOs to learn the mapping from the parameterization of sonic crystals to their corresponding *transmission loss graph*, see Section 2. This bypasses the expensive step required in conventional approaches, where one must first compute the pressure solution of the system before

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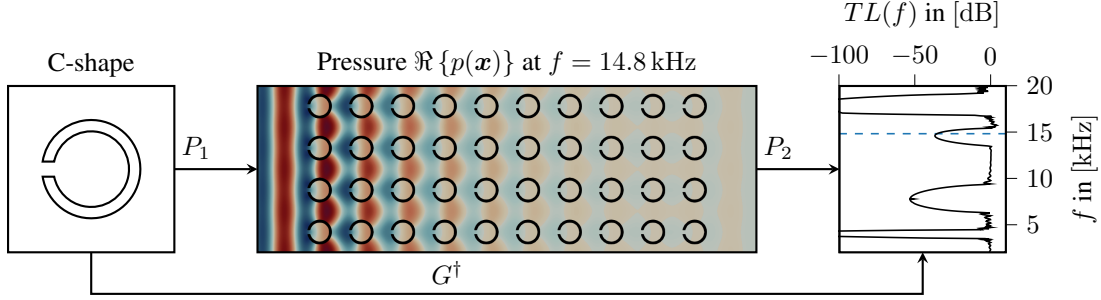


Figure 1: Our sonic crystals consist of parameterized C-shapes (left), 10 in a row in x -direction with periodicity in y -direction. The sound wave is incident from the left-hand side, traveling to the right-hand side. We learn the operator G^\dagger mapping the parameterization to the transmission loss graph of the sonic crystal (right). The plot shows the sonic crystal for $R_1 = 6.5$ mm, $R_2 = 5.0$ mm, and $b = 2.0$ mm (see text), along with the real part of the corresponding pressure field for $f = 14.8$ kHz.

evaluating the transmission loss, cf. Figure 1. We compare standard NO architectures like DeepONet (DON) [22] and FNO [21] with two new architectures: Deep Neural Operator (DNO) and Deep Cat Operator (DCO), see Section 3. All trained NOs show good generalization with slightly differing performance, see Section 4, and offer significant speed-ups over traditional methods.

2 Background and related work

Acoustics is the science that describes the behavior of sound in media, and the Helmholtz equation

$$\Delta p(x) + k^2 p(x) = 0, \quad (1)$$

is the standard modeling approach for the complex-valued sound pressure p , where $k = 2\pi f/c$ is a given *wave number* depending on a frequency f and the speed of sound c [27]. The Helmholtz equation allows for the analysis of how sound behaves under steady-state conditions, essential for understanding resonance, transmission, and the acoustic properties of materials [19].

Sonic crystals are acoustic metamaterials consisting of identical unit cells arranged in a lattice [9]. A common parameterization of sonic crystals is given by C-shapes [10], where the parameters are outer radius R_1 , inner radius R_2 , and opening width b of the C-shaped inclusion, cf. Figure 1 (left). A typical design objective is to prevent specific frequency ranges from propagating through the material, a phenomenon seen as *band-gaps* in the graph of the *transmission loss* [9].

The **transmission loss** [19] describes the decrease from the power W_i of the wave incident on one side of the material to the power W_t of the transmitted wave as it reaches the opposite side, $TL = 10 \log_{10}(W_i/W_t)$. As the transmission loss depends on the frequency f of the wave, the **transmission loss graph** $TL(f)$ is the function that maps frequency to transmission loss, and is a key characteristic of sonic crystals, cf. Figure 1 (right) [9, 10].

Conventional methods like the Finite Element Method (FEM) in frequency-domain, or the Boundary Element Method (BEM) [27] are typically used to compute transmission losses. Both require a mesh of the domain for each C-shape geometry, and compute the complex-valued pressure field by solving (1) for a range of frequencies. Afterward, the resulting pressure fields can be integrated to compute the transmission loss graph [19, 28, 10].

Operator learning uses neural networks (NNs) to approximate maps between infinite dimensional function spaces, in contrast with the usual approximation of functions between Euclidean spaces. Expressing physical quantities as functions in spaces \mathcal{A} and \mathcal{U} , operator learning approximates a physical process described by an operator $G^\dagger : \mathcal{A} \rightarrow \mathcal{U}$ with a parametric map G_θ , such that

$$G^\dagger \approx G_\theta : \mathcal{A} \rightarrow \mathcal{U}, \quad \theta \in \mathbb{R}^l, \quad (2)$$

for some number of parameters $l \in \mathbb{N}$. NOs are a flexible framework well-suited for many real-world problems. Because they are discretization-invariant, they can train and predict at different resolutions, while being much faster than traditional solvers [20].

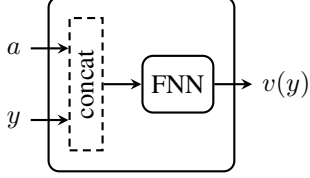


Figure 2: DNO concatenates the input function evaluations a and the evaluation location y before passing them through a single feedforward NN.

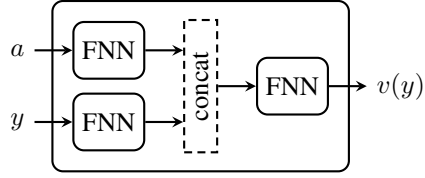


Figure 3: DCO passes the input function evaluations a and the evaluation location y through a branch and a trunk network, respectively, before passing the concatenated outputs into a third feedforward *cat network*.

DeepONet (DON, [22]) is an NO architecture motivated by the universal approximation theorem for operators. It consists of a *trunk network* T that learns a set of basis functions, and a *branch network* B that learns the corresponding basis coefficients. Given a vector a of evaluations of the input function and a coordinate y where the output function will be evaluated, a dot product linearly combines the trunk network’s basis functions with the branch network’s coefficients:

$$\mathcal{G}_{\text{DON}}(a)(y) = B(a) \cdot T(y). \quad (3)$$

The **Fourier Neural Operator** (FNO, [21]) has been successfully applied to various PDE-based problems with significant speedups [25]. It is a generic NO:

$$\mathcal{G}_{\text{FNO}}(a) = Q \circ (W_L + \mathcal{K}_L) \circ \dots \circ \sigma(W_1 + \mathcal{K}_1) \circ P(a), \quad (4)$$

where P and Q are pointwise lifting and projection operators, and the intermediate layers consist of pointwise operators W_l , integral kernel operators \mathcal{K}_l , and an activation function σ [20]. The integral kernel operators of the FNO are linear transformations in Fourier space, $\mathcal{K}_l(v) = \mathcal{F}^{-1}(R_l \cdot \mathcal{F}(v))$, where R_l are complex-valued matrices, and \mathcal{F} (resp., \mathcal{F}^{-1}) denotes the (inverse) Fourier transform.

3 Methodology

We consider parameterized sonic crystals and use NOs to predict their transmission loss graphs. Figure 1 shows the C-shaped unit cell we consider, and the transmission loss for a specific parameterization $a = [R_1, R_2, b]$ of the cell. We want to approximate the operator

$$G^\dagger(a)(f) = P_2(P_1(a, f)) : a \mapsto \text{TL}_a(f), \quad a \in \mathcal{A}, \quad f \in \mathbb{R}, \quad (5)$$

where TL_a is the transmission loss graph, evaluated at a frequency f , and a is understood as a constant input function. The intermediate maps P_1 and P_2 are the two parts of conventional techniques, where P_1 solves the Helmholtz equation for a given geometry and frequency, and P_2 computes the transmission loss from the resulting pressure field by integration. When training an NO, we can skip P_1 and P_2 and learn the mapping from the geometry parameters to the transmission loss graph directly.

We train a total of four different architectures: DON, FNO and our new DNO, DCO.

Deep Neural Operator (DNO) is inspired by the DeepONet architecture, although branch and trunk are replaced by a single (fully connected) neural network N processing both the input function and the evaluation coordinate simultaneously, cf. Figure 2. Using the same notation as in (3), the DNO is:

$$\mathcal{G}_{\text{DNO}}(a)(y) = N([a, y]), \quad (6)$$

where $[\cdot]$ denotes the concatenation of vectors. The motivation for DNO is to integrate operator information early by connecting all layers of the branch and trunk networks, maximizing flexibility in learning latent representations at the cost of increased computational complexity. Additionally, when the size of the concatenated tensors differ significantly, the model may disproportionately emphasize one over the other, leading to inefficiencies.

Deep Cat Operator (DCO) follows a similar motivation, but first passes the input function and the evaluation locations through branch and trunk networks, B and T , and concatenates those outputs in latent space before passing the result through a third *cat network* C , cf. Figure 3:

$$\mathcal{G}_{\text{DCO}}(a)(y) = C([B(a), T(y)]). \quad (7)$$

Essentially, DCO replaces the dot product in DeepONet by an NN, offering more flexibility for non-linear combinations of the trunk network’s outputs while being more robust to disparate input tensor sizes than DNO.

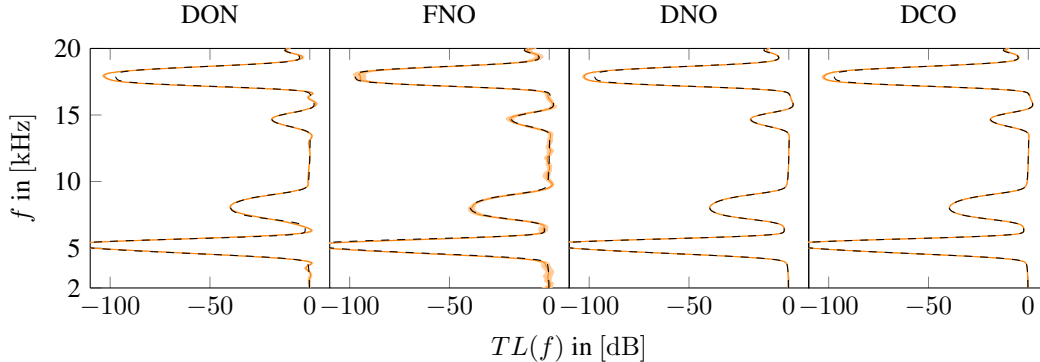


Figure 4: Mean learned transmission loss graph of four neural operators for a sample from the test set with the parameters $a = [5.7 \text{ mm}, 4.5 \text{ mm}, 3.2 \text{ mm}]$. The darker line is the mean performance, the lighter area shows the bootstrapped 95% confidence interval with mean statistic. The dashed line is the ground truth (FEM) solution.

4 Experiments

For our experiments, we employ FEniCS [29] to produce the data set and train the NO architectures implemented in `continuiti`². The setup matches the implementation in the COMSOL application gallery [28], and our dataset consists of C-shaped sonic crystals with outer radius $R_O \leq 8.5 \text{ mm}$, inner radius $R_I \leq 8.5 \text{ mm}$, and gap width $b \leq 8.5 \text{ mm}$ embedded in a unit cell with dimensions $d = 22 \text{ mm}$, while ensuring the configurations are geometrically valid. The meshes employ at least 15 elements per wavelength at 20 kHz, and we provide 673 different parameterizations with 256 TL evaluations each for training. We train the NOs for 500 epochs using Adam to minimize the MSE, and compare the final solution to a test set with 128 parameterizations, again with 256 TL evaluations. The reported results are averaged over 10 seeds, where the runtime baseline is computed using the COMSOL implementation [28]. More details on the implementation and some additional plots can be found in Appendix A.

The plot in Figure 4 shows the transmission loss for a representative sample from the test set for all four architectures. All models show good approximation properties, with DNO and DCO yielding the best generalization properties, see Table 1.

Table 1: Performance metrics of the DON, FNO, DNO, and DCO on the test dataset obtained from $n = 10$ different initial seeds. The relative mean squared error (RMSE), the bootstrapped 95% confidence interval (CI) for the model RMSE with mean statistic, number of parameters and speedup of all reported neural operators in comparison to our baselines.

Method	RMSE	95% CI	Num. Param.	Speedup
DON	0.469 %	[0.450 %, 0.492 %]	101 152	6.0×10^5
FNO	2.18 %	[1.99 %, 2.39 %]	99 553	5.7×10^6
DNO	0.199 %	[0.196 %, 0.202 %]	102 929	5.1×10^5
DCO	0.208 %	[0.203 %, 0.212 %]	100 773	1.4×10^5

5 Conclusion and future work

Our work demonstrates that neural operators can effectively learn complex relationships that require extensive computation with classical methods. The results emphasize the significance of selecting the appropriate architecture, with our new DNO and DCO models outperforming standard methods. We acknowledge the limited scope of this paper, as we focus on a single problem with a specific dataset and do not compare our approach to classical surrogate models. Thus, further experimentation is warranted, particularly to evaluate the performance of our new architectures, DNO and DCO, across a broader range of operator learning tasks. Additionally, future research could focus on directly learning the inverse problem, proposing metamaterials that achieve specific transmission loss profiles.

²Available on <https://github.com/aai-institute/continuiti>.

6 Acknowledgements

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A Implementation details

For the ground truth (FEM) solution, we simulate a domain consisting of 10 C-shaped structures aligned longitudinally in the x -direction, with a setup that matches the implementation in the COMSOL application gallery [28]. The domain is truncated in the x -direction using adiabatic layers [30] to ensure anechoic transmission for the crystal. It is truncated in the y -direction using hard sound boundaries, which is identical to periodic boundary conditions, due to the symmetry of the sonic crystal along the x -axis. We model the C-shapes with hard reflection properties, and choose the length of 10 unit cells in order to fully observe band gaps. All C-shapes have their opening pointing in the negative x -direction. We use a scattered field formulation [27] for the sound pressure to model a plane wave field incident normal to the crystal (traveling in x -direction).

For both the training and the test dataset, we uniformly sample the parameter space spanned by interval limits $R_O \leq 8.5$ mm, $R_I \leq 8.5$ mm, $b = 8.5$ mm, and a tolerance $\delta = 0.1$ mm. The unit cell has a constant width and height $d = 22$ mm. We use the tolerance to ensure that all sampled sonic crystals are valid, requiring $R_O \geq R_I + \delta$ and $R_I \geq b + \delta$. The parameter space is sampled 673 times for the training (90% train, 10% validation) and the test dataset 128 times.

We use Gmsh [31] to build meshes from the parameterizations. The meshes have at least 15 elements per wavelength at 20 kHz. We compute the FEM solution for each observation on 256 equidistant frequencies $f \in [2 \text{ kHz}, 20 \text{ kHz}]$ using the DOLFINx computational environment of the FEniCSx framework [29]. Details for this procedure can be found on Github³.

Our neural operators are implemented in our framework `continuiti`⁴. We aim to fix the numbers of all neural operators to approximately 10^5 parameters, and optimize the most important hyperparameter for our experiments, the batch size, using Optuna [32]. The internal structure of all our neural operators are fully connected NNs with residual connections, layer normalization, and hyperbolic tangent activation functions.

In order to allow the FNO to handle the parameter inputs, we use random Fourier features [33, 34] where we use Gaussian encoding with an embedding size 256 to encode the parameters of the sonic crystal. The FNO has a width of 8, depth 6, and we train it using a batch size of 4. The DON has a branch and trunk network with a width of 44, and a depth of 24 each. We use 64 basis functions, and a batch size of 4. The DNO has a width of 56, depth of 32 layers, and we train it using a batch size of 20. The DCO has a branch and trunk network with a width of 42, and a depth of 8 layers each, and the cat network is 32 layers deep with a width of 46. We train the DCO with a batch size of 20.

Each architecture is trained 10 times for the best parameters using different random seeds (but identical across different architectures). We employ Adam optimizer to minimize the mean squared error between prediction and ground truth values for 500 epochs. We reduce the learning rate from $\eta_0 = 1 \times 10^{-3}$ to $\eta_N = 1 \times 10^{-5}$ over the entire training process using cosine annealing. All samples undergo low-pass-filtering and normalization prior to being processed by the operators. The final operator is the best performing model with respect to the validation dataset. For greater detail, refer to `Neural-Operator-TL`⁵ on Github. Training for 500 epochs on our machine takes on average 526 s for the DCO, 420 s for the DNO, 1130 s for the FNO, and 1934 s for the DON.

To evaluate the error metrics, we calculate the relative mean squared error across all models, using 10 different seeds. From this we employ bootstrapped statistics to get the 95% confidence intervals for the model performance.

In terms of hardware, we use an RTX 3060 Ti, an AMD Ryzen 7 2700x CPU and 16 GB of memory. To compare the speed of our approach to conventional methods, we use a COMSOL implementation [28], which takes $t = 21.6$ s for 256 samples of the domain for one specific parameterization. The COMSOL implementation is faster than our FEniCSx code providing a reasonable state-of-the-art baseline, but in its current form it is not capable of solving many different parameterized geometries.

Figure 6 and Figure 7 show the transmission loss graph of two different configurations with both unfiltered and low-pass-filtered graphs. Figure 5 and Figure 8 display four frequency samples each for

³<https://github.com/JakobEliasWagner/Helmholtz-Sonic-Crystals>

⁴<https://github.com/aai-institute/continuiti>

⁵<https://github.com/JakobEliasWagner/Neural-Operator-TL>

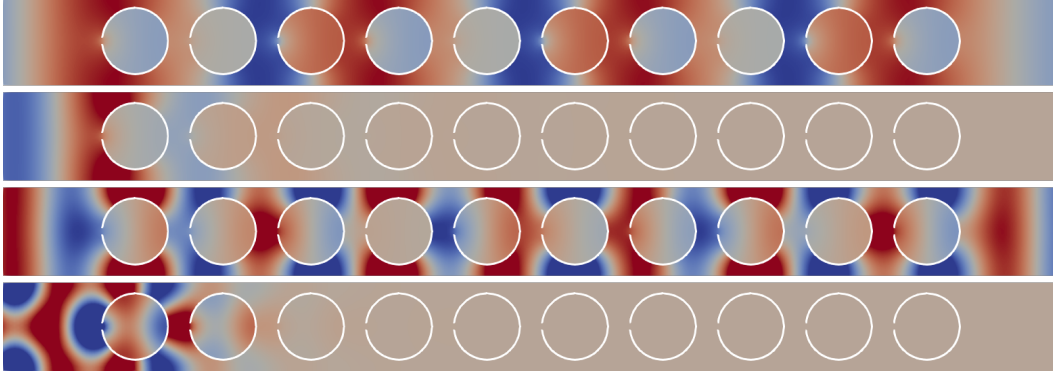


Figure 5: Real-valued pressure field for frequencies $f = \{4964 \text{ Hz}, 7505 \text{ Hz}, 11\,105 \text{ Hz}, 17\,600 \text{ Hz}\}$ (top to bottom) for the configuration $a = [8.5 \text{ mm}, 8 \text{ mm}, 1 \text{ mm}]$.

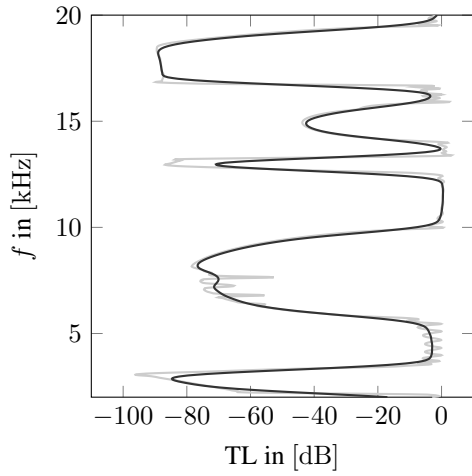


Figure 6: Unprocessed (light gray) and preprocessed (black, low-pass-filter) transmission loss graph for $a = [8.5 \text{ mm}, 8 \text{ mm}, 1 \text{ mm}]$.

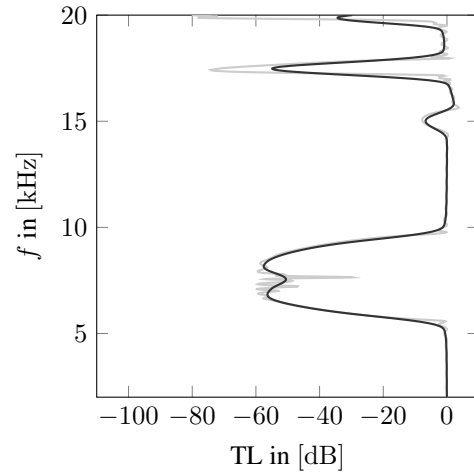


Figure 7: Unprocessed (light gray) and preprocessed (black, low-pass-filter) transmission loss graph for $a = [4 \text{ mm}, 3.7 \text{ mm}, 3 \text{ mm}]$.

both configurations. The band gaps in C-shaped sonic crystals stem from two different phenomena: Bragg-scattering (crystalline structure) and local resonance of the inner chamber of the C-shapes [9].

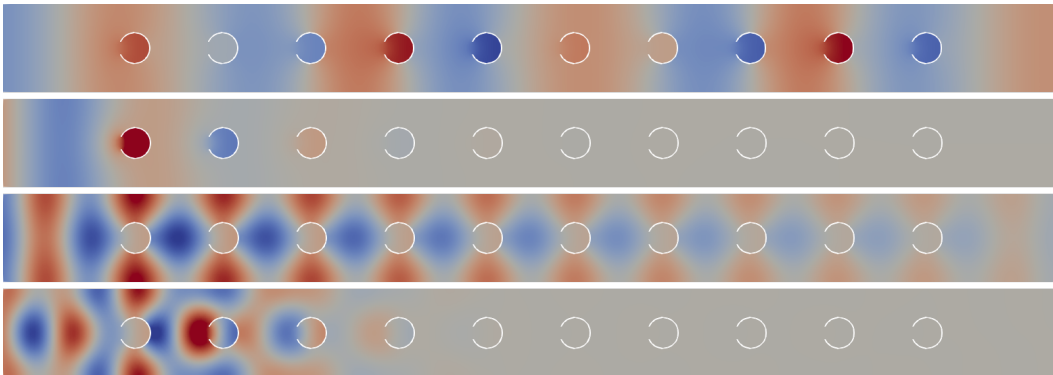


Figure 8: Real-valued pressure field for frequencies $f = \{5105 \text{ Hz}, 7505 \text{ Hz}, 15\,058 \text{ Hz}, 20\,000 \text{ Hz}\}$ (top to bottom) for the configuration $a = [4 \text{ mm}, 3.7 \text{ mm}, 3 \text{ mm}]$.

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