

000 CONTINUOUS SYMMETRY DISCOVERY AND EN- 001 002 FORCEMENT FOR IMAGE DATA 003 004

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007 008 ABSTRACT 009

010 Symmetry is an often-desired quality of machine learning models, leading, among
011 other things, to more predictable model generalization. Continuous symmetry
012 detection and enforcement for machine learning are two related topics that have
013 recently been explored using the Lie derivative along vectors fields, which vector
014 field approach has led to improved outcomes. However, though image data is
015 replete with continuous symmetries under which image classifiers are meant to be
016 invariant, the application of the Lie derivative for the detection and enforcement
017 of continuous symmetries for image data remains underexplored. In this work,
018 we derive vector field infinitesimal generators for various continuous symmetries
019 for image data. We then use these generators to enforce continuous symmetries in
020 image classifiers. We also demonstrate vector field symmetry detection in image
021 data, obtaining close similarity with the ground truth symmetry.

024 1 INTRODUCTION

025 Symmetry-informed machine learning is a field which has continued to accumulate interest, with
026 previous work demonstrating its effectiveness in improving model performance (Lyle et al., 2020;
027 Bergman, 2019; Craven et al., 2022; Tahmasebi & Jegelka, 2023; Ko et al., 2024). Within the
028 context of image classification tasks, one is often concerned with learning neural networks which
029 are invariant to pre-specified transformations of images. Although a common approach to building
030 invariant models in practice is to explicitly augment the training data with transformed copies of
031 images, recent work has been conducted with the goal of enforcing either equivariance or invariance
032 in models without the need for data augmentation (Finzi et al., 2020). The general need to enforce
033 symmetry in models without augmentation is due to two primary factors: (1) data augmentation
034 can be computationally expensive (Mumuni & Mumuni, 2022); (2) data augmentation can only be
035 applied if one knows or can discover the symmetry group explicitly. It is the second of these factors
036 that is primarily applicable to image data, as common augmentations, such as rotations or adjusting
037 the brightness, are relatively computationally inexpensive: nevertheless, employing augmentation
038 for image data requires an explicit transformation of the data, while the transformation group may
039 not be known *a priori*.

040 Recent work has shown that the Lie derivative along vector fields can be used to enforce and dis-
041 cover continuous symmetries in models quite generally, which approach can be computationally
042 efficient (Otto et al., 2024; Shaw et al., 2024; 2025). However, it appears that, despite the many con-
043 tinuous symmetries for image data, previous work with regard to image data deals only with certain
044 types of continuous transformations, such as planar affine transformations. In this work, we propose
045 an approach to symmetry discovery and enforcement for image data using the Lie derivative along
046 vector fields, thereby expanding the types of symmetries that can be considered.

047 Continuous symmetry enforcement for models has recently been accomplished using vector fields
048 as a regularization term in the model loss function, similar to the framework of Physics-Informed
049 Neural Networks (Shaw et al., 2025; Otto et al., 2024; Raissi et al., 2019). However, this approach,
050 as currently defined, cannot directly be applied to symmetries for image data. This is because the
051 approach requires that each datapoint reside in \mathbb{R}^d , where d is the number of features: Although
052 an image can be “flattened” so as to be readily interpretable as a vector in \mathbb{R}^d , common continuous
053 image transformations such as those making use of convolutions, may not be recognizable in this
feature space. Even when an image transformation is easily expressible in the flattened feature

space, it is often best practice to use convolutional layers when training on image data, highlighting the benefit of maintaining the integrity of the image structure. Therefore, the extension to image data is non-trivial and, as will be shown, is fruitful. As the vector field approach to model symmetry discovery and enforcement relies on the computation of the gradient of the model function, our extension to image data makes use of the gradient of the model with respect to an input image. The mathematical framework for this extension also relies on the notion of diagonal group actions.

Our main contribution is summarized as follows: we provide a mathematical framework for the extension of continuous symmetry discovery and enforcement for image data. We also provide experimental results which demonstrate the feasibility of implementing our method. Our results suggest that continuous symmetry enforcement in image classification using vector fields is comparable to augmentation and is applicable in cases in which augmentation cannot be used. Our results also suggest that model symmetry can be discovered and quantified using the Lie derivative.

2 RELATED WORK AND BACKGROUND

In this section, we discuss both related work and background information which is useful for understanding the context of our contributions. In Sections 2.1 and 2.2, we discuss previous work in symmetry discovery and enforcement, respectively. Section B contains an overview of vector fields and flows, being a somewhat new concept for machine learning, though we note that a similar discussion is contained in existing literature (Otto et al., 2024; Shaw et al., 2024). Section C gives a brief overview of infinitesimal generators of multi-parameter groups, and section 2.3 gives a brief overview of diagonal group actions and their infinitesimal generators.

2.1 RELATED WORK: SYMMETRY DISCOVERY

Early work on symmetry detection in machine learning focused primarily on detecting symmetry in image and video data (Rao & Ruderman, 1999; Sohl-Dickstein et al., 2010), where symmetries described by straight-line and rotational transformations were discovered. Other work has made strides in symmetry discovery by restricting the types of symmetries being sought. In one case, detection was limited to compact Abelian Lie groups (Cohen & Welling, 2014) and used for the purpose of learning disentangled representations. Another case uses meta-learning to discover any *finite* symmetry group (Zhou et al., 2021). Finite groups have also been used in symmetry discovery in representation learning (Anselmi et al., 2019). In physics-based applications of machine learning, a method was developed that can discover any classical Lie group symmetry (Forestano et al., 2023). Symmetry discovery of shapes has also been explored (Je et al., 2024).

Other work has focused on detecting affine transformation symmetries and encoding the discovered symmetries automatically into a model architecture. Some methods identify Lie algebra generators to describe the symmetries. For example, *augerino* (Benton et al., 2020) attempts to learn a distribution over augmentations, subsequently training a model with augmented data. The *Lie algebra convolutional network* (Dehmamy et al., 2021), which generalizes Convolutional Neural Networks in the presence of affine symmetries, uses infinitesimal generators represented as vector fields to describe the symmetry. SymmetryGAN (Desai et al., 2022) has also been used to detect rotational symmetry (Yang et al., 2023).

Another notable contribution to efforts to detect symmetries of data is *LieGAN*. LieGAN is a generative-adversarial network intended to return infinitesimal generators of the continuous symmetry group of a given dataset (Yang et al., 2023). LieGAN has been shown to detect continuous affine symmetries, including transformations from the Lorentz group. It has also been shown to identify discrete symmetries such as rotations by a fixed angle.

While most continuous symmetry detection methods attempt to discover symmetries which are affine transformations, the representation of infinitesimal generators using vector fields has led to the discovery of continuous symmetries which are not affine (Ko et al., 2024; Shaw et al., 2024). In one case, the domains of image data and partial differential equations are examined in particular (Ko et al., 2024).

Continuous symmetry detection is more difficult than discrete symmetry detection (Zhou et al., 2021) since the condition of invariance given as $f \circ S = f$ must hold for all values of the continuous

108 parameter of S . This is corroborated by the increasingly complex methods used to calculate even
 109 simple symmetries such as planar rotations (Benton et al., 2020; Dehmamy et al., 2021; Yang et al.,
 110 2023). Some methods introduce discretization, where multiple parameter values are chosen and
 111 evaluated. LieGAN does this by generating various transformations from the same infinitesimal
 112 generator (Yang et al., 2023). We believe a vector field approach, as has become more recently
 113 common (Otto et al., 2024), addresses the issue of discretization. A vector field-based method
 114 reduces the required model complexity of continuous symmetry detection while offering means to
 115 detect symmetries beyond affine transformations (Finzi et al., 2020; Shaw et al., 2024).

116 2.2 RELATED WORK: SYMMETRY ENFORCEMENT

117 Our work is most closely related to previous work on symmetry enforcement using vector
 118 fields (Finzi et al., 2020; Otto et al., 2024; Shaw et al., 2025), since we are directly adapting the
 119 analogous methods presented by the various authors for image data and common image transforma-
 120 tions. This method is also analogous to Physics-Informed Neural Networks (PINNs) (Raissi et al.,
 121 2019). With PINNs, model training is regularized using differential constraints which represent the
 122 governing equations for a physical system. The method of symmetry enforcement employed here
 123 differs from PINNs since the differential constraints obtained using infinitesimal generators do not
 124 generally have the interpretation of defining governing equations for a physical system.

125 Continuous symmetry enforcement in images is far from new, and we note some recently-developed
 126 methods. Some methods seek to enforce symmetry by augmenting the training dataset according to
 127 known symmetries (Bergman, 2019). *Augerino* attempts to enforce symmetry using augmented data,
 128 though the symmetries are discovered from the data rather than given *a priori*. Another established
 129 method of enforcing symmetry is feature averaging, which is thought to be generally more effective
 130 than data augmentation (Lyle et al., 2020).

131 A growing number of methods seek to use symmetry to construct invariant or equivariant models
 132 without the need for augmented training data, and there is previous work accomplishing this using
 133 infinitesimal generators (Dehmamy et al., 2021; Yang et al., 2023; Finzi et al., 2020; Otto et al.,
 134 2024). Some previous work addresses specific cases: the special case of compact groups is stud-
 135 ied (Bloem-Reddy & Teh, 2020), and the case of equivariant CNNs on Homogeneous spaces is
 136 studied (Cohen et al., 2019). Other work speaks to the universality of invariant architectures (Maron
 137 et al., 2019; Keriven & Peyré, 2019; Yarotsky, 2018).

138 2.3 DIAGONAL GROUP ACTIONS AND THEIR INFINITESIMAL GENERATORS

139 Fundamental to our approach is the concept of multiparameter groups and their infinitesimal gener-
 140 ators. As these topics are discussed in existing literature Shaw et al. (2024; 2025), we provide an
 141 overview only in Appendices B and C. Of fundamental importance for the specialization to image
 142 data is the subject of infinitesimal generators of diagonal group actions, which we now discuss.

143 Let G be a group acting on a set Y . The action of G on Y induces an action on the set Y^m , where
 144 a group element $g \in G$ acts on $(y_1, y_2, \dots, y_m) \in Y^m$ by acting on each component y_i separately
 145 and in the manner in which g acts on $y \in Y$. It is said that the action of G on Y^m is a diagonal
 146 action.

147 If $\Psi(t, \mathbf{x})$ is a flow on \mathbb{R}^n , the induced diagonal action on $\mathbb{R}^n \times \mathbb{R}^n \times \dots \times \mathbb{R}^n = (\mathbb{R}^n)^m$ is defined:

$$148 \tilde{\Psi}(t, (\mathbf{x}_1, \mathbf{x}_2, \dots, \mathbf{x}_m)) = (\Psi(t, \mathbf{x}_1), \Psi(t, \mathbf{x}_2), \dots, \Psi(t, \mathbf{x}_m)).$$

149 If X is the vector field infinitesimal generator for $\Psi(t, \mathbf{x})$, we denote the infinitesimal generator for
 150 $\Psi(t, \mathbf{x}_i)$ by X_i , so that the infinitesimal generator \tilde{X} for $\tilde{\Psi}$ is given as

$$151 \tilde{X} = X_1 + X_2 + \dots + X_m.$$

152 For example, let $\Psi(t, (x, y)) = (x + t, ty)$, so that $X = \partial_x + y\partial_y$, the level curves for which can be
 153 written as $y = ce^x$. The infinitesimal generator for the diagonal action of Ψ on $\mathbb{R}^2 \times \mathbb{R}^2 \times \mathbb{R}^2$ is

$$154 \tilde{X} = \partial_{x_1} + y_1\partial_{y_1} + \partial_{x_2} + y_2\partial_{y_2} + \partial_{x_3} + y_3\partial_{y_3}.$$

155 Our model assumption is that an image with m channels is an element of $(\mathbb{R}^n)^m$, with n being the
 156 number of pixels. This merely formalizes the concept of a group “acting on each image channel

separately.” However, another fruitful perspective outside the context of image data deals with samples, where m represents the number of data points in a given dataset and where n represents the number of features. It is within this context that the sample mean can be said to be *equivariant* under certain affine transformations: a group acts on each point $\{\mathbf{x}_i\}_{i=1}^m$ separately, where $\mathbf{x}_i \in \mathbb{R}^n$, so that the sample mean of $\tilde{\Psi}(t, (\mathbf{x}_1, \mathbf{x}_2, \dots, \mathbf{x}_m))$ coincides with $\Psi(t, \cdot)$ applied to the sample mean of $\{\mathbf{x}_i\}_{i=1}^m$.¹

3 METHODS

We model grayscale images as matrices: an image I is in $\mathbb{R}^{p \times q}$, where $p \times q$ is the number of pixels. We model multi-channel images as tensors—that is, matrices of size $p \times q$ for each channel. A diagonal action for an RGB image thus admits multiple interpretations, as it can be (1) a diagonal action on $p \times q$ copies of \mathbb{R}^3 , or (2) a diagonal action on three copies of $\mathbb{R}^{p \times q}$.

Another common model of group actions on images involves modeling images as sections of a vector bundle on \mathbb{R}^2 : that is, as mappings $I : \mathbb{R}^2 \rightarrow \mathbb{R}^c$ (for a c -channel image). Rotations and translations, among other planar transformations, are naturally considered by defining group actions in the plane, and this is the more common approach to symmetry detection and enforcement for image data. In this paper, however, we do not consider group actions on the base space, separating our work from that of previous work.

3.1 FRAMEWORK FOR CONTINUOUS SYMMETRY DISCOVERY

Previous work has dealt with the discovery of vector field infinitesimal generators (Shaw et al., 2024; 2025). The approach previously employed can be summarized as the following multi-step framework:

1. Utilize a parametric assumption about the form of the generators $\{X_j\}$ of the multi-parameter group action
2. Obtain a smooth function f to examine for symmetry
3. Optimize the parameters of the generator(s) sought using $X_j(f) = 0$

For image data, using our model of images as tensors, steps 1 and 2 require modification for two reasons: (1) image data are high-dimensional; (2) not all transformations respect the spatial structure of images.

3.1.1 ADAPTING PARAMETRIC ASSUMPTIONS FOR IMAGE DATA

The high-dimensional problem is no small matter. To detect affine symmetry for the MNIST (Deng, 2012) dataset, for example, one would require at least $28^4 > 600,000$ parameters to be present in the parametric form of the generator, which problem is duplicated for matrix representations, the sizes of which would be $(28 \times 28) \times (28 \times 28)$. For context, this is an order of magnitude larger than the number of training samples available, and closer to two orders of magnitude larger than the number of parameters needed to train a successful classifier on the test dataset. The problem is much worse for ImageNet (Deng et al., 2009), where the sizes of the images are 224×224 : generators with over 2.5 **billion** parameters would be needed, not even counting the need to duplicate across three separate channels. And as is the case for MNIST, this large number of parameters is at least one order of magnitude (though more often three orders of magnitude) larger than the maximum of the number of training images and the number of parameters needed to train a classifier with reasonable performance.

To address this issue, we make one or more parametric assumption(s) about the generators which (1) involve exponentially fewer parameters than described above; (2) are likely appropriate specifically for image data. The first of these assumptions is that the unknown group action acts diagonally: either the same group action on \mathbb{R}^c at each of the $p \times q$ pixels or else the same group action on $\mathbb{R}^{p \times q}$

¹We note that equivariance under flows, like invariance, can be characterized in terms of the Lie derivative (Otto et al., 2024), though the subject of the current work is invariance and not equivariance.

216 across each c channels. In the first case, symmetry discovery takes place in \mathbb{R}^c rather than $\mathbb{R}^{c \times (p \times q)}$:
 217 note that $p \times q >> c$, generally. In the second case, symmetry discovery takes place in $\mathbb{R}^{p \times q}$ instead
 218 of $\mathbb{R}^{c \times (p \times q)}$: a reduction only of a small amount. The first case is appropriate for many types of
 219 images transformations, including the power law transform as well as adjusting brightness and hue.
 220 The second case is appropriate for transformations of pixels which depend on the values of other
 221 (perhaps neighboring) pixels, which takes us to our second restrictive assumption.

222 The second useful restrictive assumption is that the group action can be represented as a convolution
 223 of a fixed kernel applied to each image. This can leverage the second case of diagonal actions in
 224 the case where a kernel matrix defines a convolution across each channel separately. But more
 225 importantly for the reduction of the number of parameters, an affine transformation represented as
 226 a convolution of a kernel with images a parameter count equal to the size of the kernel: for the
 227 Gaussian blur transformation, for example, the kernel is typically on the order of 7×7 or 15×15 , a
 228 sharp reduction in the number of parameters when compared with the entire affine space for images.
 229 Additionally, a restriction of this type is appropriate for image data, owing to transformations and
 230 architectures which make explicit use of convolutions: in particular, they make use of the local
 231 (spatial) image structure.

232 3.1.2 ADAPTING THE LEARNING OF THE MANIFOLD STRUCTURE

234 The high-dimensional problem apparent in image data also requires the adaptation of the second
 235 step in the symmetry discovery framework. Previous work on symmetry discovery using vector
 236 fields points to the need to explicitly identify a smooth function to test for symmetry (Shaw et al.,
 237 2024): therein, the authors introduce *level set estimation*, in addition to making use of probability
 238 distributions. However, the method of level set estimation does not appear to be developed, as
 239 yet, for high-dimensional data, as the experiments conducted using it are restricted to 10 or fewer
 240 dimensions, which is far below what is required for images. Probability density estimation is also
 241 difficult in high dimensions: while density and level set estimation work with this framework in
 242 principle, we offer an adaptation to this step which may be better suited in higher dimensions.

243 Let I be an image in the training set, and let $\Psi(t, \cdot)$ be a flow. For each I , we assume an ordered se-
 244 quence of transformed images $\{\Psi(t_j, I)\}_{j=1}^k$ (with $t_j < t_{j+1}$) and that there is an index $1 \leq i \leq k$
 245 with $\Phi(t_i, I) = I$: this is a somewhat restrictive setting, though not unlike trajectory-based symme-
 246 try discovery that appears elsewhere (Yang et al., 2023). The vector field infinitesimal generator for
 247 the flow at I is found by computing $\frac{d}{dt}(\Phi(t, I))|_{t=t_i}$, which can be approximated using numerical
 248 methods such as finite differences.

249 3.2 FRAMEWORK FOR CONTINUOUS SYMMETRY ENFORCEMENT

251 As with continuous symmetry discovery, symmetry enforcement using vector field regularization
 252 has been previously proposed (Shaw et al., 2024; 2025), the steps of which can be given as the
 253 following multi-step framework:

- 255 1. Identify the infinitesimal generators $\{X_j\}$ corresponding to the symmetry to be enforced
- 256 2. Regularize the training of a smooth model f using $X_j(f) = 0$ and an estimated gradient of
 257 f (with respect to the inputs, not parameters)

258 The only adaptation of the symmetry enforcement framework needed is that which naturally follows
 259 our adaptation of the discovery framework: the transformations and, subsequently, the vector field
 260 infinitesimal generators, are specified using diagonal group actions, convolutions, or a combination
 261 of the two.

263 3.3 THE PRACTICALITY OF A VECTOR FIELD APPROACH FOR IMAGES

265 In image classification tasks, it is common practice to augment using several types of transfor-
 266 mations. Thus, it is natural to consider symmetry regularization for multi-parameter groups, the
 267 infinitesimal generators of which were discussed in Section C. In practice, multi-parameter groups
 268 could be handled in multiple ways. Our approach is to “stack” the tensors defined by applying the
 269 components of the infinitesimal generators to the data, so that the regularization term has a larger
 shape than in the single-parameter case.

270 In general, flows of vector fields do not commute. As a simple example, planar rotations about the
 271 origin do not commute with translations in the horizontal direction. In fact, this can be seen by
 272 analyzing the Lie bracket of two vector fields: the flows of the vector fields commute if and only
 273 if the Lie bracket of the vector field vanishes identically. We see this in the simple example: the
 274 Lie bracket of ∂_x and $-y\partial_x + x\partial_y$ is ∂_y . As our experiments rely on the power law transform and
 275 Gaussian blur, it should also be said that these transformations do not commute.

276 The presence of non-commuting transformations presents an additional consideration for data aug-
 277 mentation: a collection of transformations could be composed in any possible order, if the transfor-
 278 mations do not commute. When considering the set of possible augmentations, this number grows
 279 rather quickly. Without considering order, m augmentation types (transformations) with n selected
 280 hyperparameter values each (discretizing m continuous transformations) would yield a hyperparam-
 281 eter grid of size n^m . However, when considering the order in which these augmentations take place,
 282 the number of possible ways to uniquely augment a sample increases much faster: as there are $m!$
 283 ways to order the transformations, we have $m! \cdot n^m$ unique augmentations. Even for small values
 284 of n , this increases combinatorially. However, this does not, in principle, affect the regularization
 285 approach, since we simply require that $X_i(F) = 0$ for each generator X_i .

286 There is also another practical consideration that leads us to consider a vector field approach. We
 287 will see in the experiments below that although symmetry regularization can obtain a similar effect
 288 as augmentation, augmentation has a tendency to produce models which are more symmetric, and
 289 not always at the expense of accuracy. The computational overhead of computing the model gradient
 290 for symmetry regularization, despite the fact that this gradient can be used for multiple symmetry
 291 terms in multi-parameter enforcement, is also a practical concern, leaving one to wonder whether
 292 symmetry regularization has practical value. One reason symmetry regularization is of practical
 293 value is due to the fact that the symmetry group, be it single- or multi-parameter, could be unknown
 294 before performing symmetry enforcement in a particular classification task. In the case where vector
 295 fields are used to characterize the symmetry, direct augmentation is generally difficult, owing to the
 296 need to compute the flow of a vector field (or many), which requires solving a system of first-order
 297 ordinary differential equations. Symmetry enforcement through vector field regularization, on the
 298 other hand, eliminates the need to identify the symmetry group explicitly in order to enforce the
 299 symmetry. Thus, a critical practical use of vector field regularization is that it enables the use of
 300 discovered continuous symmetries described by vector fields, which method of discovery is both
 301 efficient and expressive (Shaw et al., 2024).

301 The need to identify a function to test for symmetry may seem to restrict the ability to perform
 302 symmetry discovery in the context of image data. However, this need is also what allows symmetry
 303 discovery of models themselves to take place. A successfully-trained model may be probed for its
 304 symmetry properties, which symmetries may subsequently be used to regularize the training of other
 305 models. This approach is, of course, not unique to the case of images data: however, the restrictive
 306 assumptions on the vector fields themselves (diagonal actions and convolutions) allow for symmetry
 307 discovery of models to be completed without undue computational resources (we experiment with
 308 this in 4.3).

309

310 4 EXPERIMENTAL RESULTS

311

312 Our first group of experiments deals with symmetry enforcement. Our first two experiments show-
 313 case symmetry enforcement for two types of transformations, namely the power law transform and
 314 Gaussian blur. Our third experiment tackles symmetry enforcement in the case of a subset of Im-
 315 agenet using a pretrained ResNet50 model (He et al., 2015; Deng et al., 2009). After symmetry
 316 enforcement experiments, we perform a symmetry discovery experiment, simulating the case where
 317 symmetry is discoverable in data but where the symmetry group is not explicitly recovered.

318

319

320 4.1 SYMMETRY ENFORCEMENT RESULTS

321

322 In this paper, we concern ourselves primarily with the power law transformation and Gaussian blur.
 323 The application of the power law transformation is sometimes called “gamma correction,” as frac-
 324 tional parameter values lead to an image corrected for darkness.

324 4.1.1 POWER LAW SYMMETRY ON MNIST
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326 In this section, we consider single-parameter symmetry enforcement experiments which make use of
327 the power law transform. The infinitesimal generator for this transformation is derived in Appendix
328 A.1. We train a simple convolutional neural network using the MNIST dataset in four distinct
329 approaches. Each approach can be described in terms of Equation 20. Each approach uses the cross
330 entropy loss for \mathcal{L} , the Adam optimizer (Kingma & Ba, 2014) with a learning rate of 0.001, and
331 trains for twelve epochs. The Baseline approach uses $\lambda(t) = 0$. The Regularization approach uses
332 $\lambda(t) = 0$ for the first three epochs, followed by $\lambda(t) = 0.5$ for the remaining epochs, where $\tilde{\mathcal{L}}$ is
333 the mean square error loss scaled by 28^2 . The Augmentation approach is similar to the Baseline
334 approach, but applies the power law transform according to a random (uniform) sample of γ -values
335 from $\{0.25, 1.0, 1.75\}$. Table 1 summarizes our results for 10 independent trials (median \pm IQR/2)
336 for select values of γ . Values through 1.9 were not substantially different from values at 1.0 and are
337 thus not shown.

338 Table 1: Power Law Transformation Results for the MNIST dataset (Median \pm IQR/2)
339

γ	Baseline	Regularization	Augmentation	Reg + Aug
0.1	0.4435 ± 0.0757	0.7289 ± 0.0528	0.9643 ± 0.0318	0.9066 ± 0.0438
0.2	0.8670 ± 0.0332	0.9497 ± 0.0365	0.9757 ± 0.0315	0.9275 ± 0.0446
0.25	0.9309 ± 0.0359	0.9604 ± 0.0357	0.9765 ± 0.0317	0.9295 ± 0.0446
1.0	0.9684 ± 0.0413	0.9690 ± 0.0355	0.9773 ± 0.0323	0.9317 ± 0.0449

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346 Table 1 shows that Regularization can encourage a model to be more symmetric than a baseline
347 approach while still maintaining fairly high accuracy on the original test set. However, the Aug-
348 mentation approach appears to outperform Regularization for $\gamma = 0.1$. Curiously, a combined
349 Regularization/Augmentation approach seemed to underperform Augmentation and Regularization,
350 generally.

351 4.1.2 GAUSSIAN BLUR SYMMETRY ON MNIST
352

353 The experimental setup is similar to that of Section 4.1.1. The primary difference is that the in-
354 finitesimal generator is that of the (estimated) generator for Gaussian blur, derived in Appendix A.2.
355 Additionally, the symmetry loss $\tilde{\mathcal{L}}$ is scaled by $10 \cdot 28^2$ rather than 28^2 , and the models are trained
356 for 10, not 12, epochs.

357 Table 2: Gaussian Blur Transformation Results for the MNIST dataset (Median \pm IQR/2)
358

σ	Baseline	Regularization	Augmentation	Reg + Aug
0.1	0.8924 ± 0.0024	0.9710 ± 0.0439	0.9216 ± 0.0409	0.9797 ± 0.0022
1.0	0.8867 ± 0.0027	0.9594 ± 0.0421	0.9225 ± 0.0405	0.9714 ± 0.0034
2.0	0.8451 ± 0.0138	0.8731 ± 0.0438	0.9167 ± 0.0391	0.8969 ± 0.0022
3.0	0.7747 ± 0.0201	0.7524 ± 0.0618	0.9103 ± 0.0371	0.8108 ± 0.0156
6.0	0.7076 ± 0.0354	0.6288 ± 0.0389	0.8998 ± 0.0346	0.7234 ± 0.0485

359 The results of Table 2 demonstrate that the regularization parameter $\lambda(t)$ can be chosen such that
360 the model performance on the original test is much more favorable than for transformed copies.

361 4.1.3 RESNET50 ON IMAGENETTE
362

363 To examine the ability of symmetry regularization to apply to large models, we look to the ResNet50
364 model (He et al., 2015) on a 10-class subset of ImageNet (Deng et al., 2009) known as Im-
365 ageNette (Paszke et al., 2019). We do not adjust the model architecture, but rather map the predicted
366 index to an index between 0 and 9 to obtain an adjusted prediction.

367 We train the model for two epochs without any pretrained weights. With regard to Equation (20),
368 we select the Cross Entropy loss function for \mathcal{L} , $224 \cdot 224 \cdot 3$ multiplied by the Mean Square Error

378 for $\tilde{\mathcal{L}}$, and $\lambda(t) = 0.5$. We optimize using SGD (Sutskever et al., 2013) with momentum 0.9 and
 379 learning rate 0.01. We repeat this process five times, evaluating both on the original validation set as
 380 well as a blurred copy of the validation set with $\sigma = 3.0$ and a kernel size of 15. (For regularization,
 381 the infinitesimal generator is derived in a similar manner as in Appendix A.2, but with a kernel size
 382 of 15 instead of seven.) Subsequently, we repeat this process for a model using no regularization.
 383 No additional strategies to improve model training are employed.

384 The accuracy score of the unregularized model in terms of the mean \pm standard deviation is
 385 $39.02\% \pm 4.74\%$ for the original validation set and $27.7\% \pm 3.33\%$ evaluating at $\sigma = 3.0$. For
 386 the regularized model, the scores are 47.87 ± 1.29 and 33.49 ± 2.47 , respectively. This result not
 387 only shows a significant improvement obtained using symmetry regularization, but also that sym-
 388 metry regularization can be applied for larger problems.

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390 4.2 SYMMETRY DISCOVERY FROM DATA

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392 We now demonstrate the discovery of the infinitesimal (vector field) generator for Gaussian blur
 393 symmetry. As discussed in Section 3.3, the symmetry group, be it single- or multi-parameter, could
 394 be unknown before performing symmetry enforcement in a particular classification task, creating
 395 a need to discover the symmetry. In this experiment, we simulate an unknown symmetry group
 396 present in the MNIST dataset by first augmenting the dataset according to $\sigma \in \{0.1i\}_{i=0}^{10}$. Next,
 397 for each original training image and its augmentations, we compute the gradient along σ for each
 398 $\sigma \in \{0.1i\}_{i=0}^{10}$.² Generally, the infinitesimal generator is found by evaluating the derivative at the
 399 group identity—in this case, zero. However, as is explained in Appendix A.2, we approximate the
 400 infinitesimal generator by evaluating at $\sigma = 0.3$.

401

402 We have thus estimated a gradient image for each training image, which gradient image, up to
 403 scaling, should correspond with the output of the Gaussian blur infinitesimal generator applied to
 404 the training image. As we represent the Gaussian blur infinitesimal generator in Appendix A.2 by
 405 convolving an image with a 7x7 kernel, we can estimate this kernel by training a single-layer con-
 406 volutional neural network (with a kernel of size 7x7), using the training images as input and the
 407 gradient images as targets. We train this small model using the L1Loss function and the Adam
 408 optimizer (Kingma & Ba, 2014), with a learning rate of 0.01, for 50 epochs. We obtain the follow-
 409 ing (approximate) 7x7 kernel, which has a cosine similarity of 0.9998 with respect to the estimate
 410 ground truth kernel:

411

$$\begin{bmatrix} 4.35 \cdot 10^{-3} & 2.35 \cdot 10^{-3} & 1.01 \cdot 10^{-3} & 4.90 \cdot 10^{-4} & -8.28 \cdot 10^{-5} & 6.33 \cdot 10^{-5} & 4.38 \cdot 10^{-4} \\ 2.95 \cdot 10^{-3} & 1.48 \cdot 10^{-3} & 1.90 \cdot 10^{-3} & 5.68 \cdot 10^{-3} & 6.30 \cdot 10^{-4} & 5.45 \cdot 10^{-5} & 4.42 \cdot 10^{-4} \\ 2.07 \cdot 10^{-3} & 1.59 \cdot 10^{-3} & 2.08 & 47.3 & 2.08 & -1.04 \cdot 10^{-5} & 3.70 \cdot 10^{-4} \\ 1.30 \cdot 10^{-3} & 5.01 \cdot 10^{-3} & 47.3 & -198 & 47.3 & 4.17 \cdot 10^{-3} & 5.47 \cdot 10^{-4} \\ 1.06 \cdot 10^{-3} & 1.26 \cdot 10^{-3} & 2.08 & 47.3 & 2.08 & 8.22 \cdot 10^{-4} & 6.92 \cdot 10^{-4} \\ 1.51 \cdot 10^{-3} & 1.33 \cdot 10^{-3} & 1.71 \cdot 10^{-3} & 5.24 \cdot 10^{-3} & 1.08 \cdot 10^{-3} & 3.94 \cdot 10^{-4} & 4.82 \cdot 10^{-4} \\ 2.14 \cdot 10^{-3} & 1.60 \cdot 10^{-3} & 1.27 \cdot 10^{-3} & 1.25 \cdot 10^{-3} & 8.67 \cdot 10^{-4} & 3.20 \cdot 10^{-4} & 5.86 \cdot 10^{-5} \end{bmatrix}$$

412

413 4.3 MODEL SYMMETRY DISCOVERY

414

415 In this experiment, we consider the problem of symmetry discovery of a pretrained model. We begin
 416 by training a small CNN F on the MNIST dataset, and we note that the accuracy on the test set is
 417 approximately 0.9716. We now seek to approximate a generator X such that $X(F) = 0$ for each
 418 training image. We assume that the coefficient functions of X are expressible in terms of a 3×3
 419 convolutional kernel: that is, there is a 3×3 matrix K such that $X(F)$ evaluated at an image I ,
 420 denoted $X(F)|_I$, is given as

421

$$X(F)|_I = K * \nabla F(I).$$

422

423 Thus, with F fixed, we seek to learn K such that

424

$$K * \nabla F(I) = 0, \quad \|K\| = 1. \quad (1)$$

425

426 ²Previous work has computed vector field generators from level sets (Shaw et al., 2024). However, due to
 427 the high dimensionality of image data, a directional derivative along the flow path may yield a more efficient
 428 estimation algorithm.

432 We impose the constraint $\|K\| = 1$ so that the learned weights of K do not become identically 0:
 433 we reiterate that our discovered K is equivalent to any uniform scaling of K .
 434

435 Our symmetry discovery task is thus one of training a CNN with a single convolutional layer defined
 436 by a 3×3 kernel (using a stride of one, and where each input image is padded with zeros so that the
 437 shape of the output image is that of the input image). The constraint $\|K\| = 1$ is applied manually
 438 by dividing the kernel, during each epoch, by its own magnitude. We use the L1 loss function (scaled
 439 by $10 \cdot 28^2$), optimizing using the Adam(Kingma & Ba, 2014) algorithm with a learning rate of 0.01,
 440 training for 100 epochs. The result is the following kernel, which we take to define the discovered
 441 symmetry:
 442

$$\begin{bmatrix} -0.5993 & 0.9755 & -0.5215 \\ 0.9917 & -1.7328 & 1.0194 \\ -0.4998 & 0.9674 & -0.6031 \end{bmatrix}.$$

446 5 CONCLUSION

447
 448 In this paper, we have introduced continuous symmetry discovery and enforcement for image data
 449 using the Lie derivative along vector fields. We have shown that symmetry regularization can yield
 450 similar outcomes to augmentation. We have also shown that symmetry can be discovered from image
 451 data directly, and while the efficient method of discovering symmetry using vector fields does not
 452 yield the symmetry group explicitly, the discovered vector field infinitesimal generator(s) can still be
 453 used to enforce symmetry in downstream tasks. Therefore, symmetry regularization allows one to
 454 leverage symmetry discovery using vector fields, which is computationally efficient when compared
 455 with existing methods (Shaw et al., 2024; Hu et al., 2025). While symmetry enforcement using
 456 augmentation requires direct knowledge or discovery of the symmetry group itself, the symmetry
 457 group for a Lie algebra of vector fields is, by assumption, the Lie group generated by the flows of a
 458 Lie algebra basis for the vector fields.
 459

460 Future work includes optimizing the runtime of symmetry regularization. Due to the need to com-
 461 pute the gradient with respect to the input images, symmetry regularization, in our experiments, takes
 462 significantly more time than training without. It may be possible to reduce the computation time with
 463 a more efficient implementation (such as computing gradients across the output dimensions in par-
 464 allel), and it may also be possible to use cached gradients from the layers of the parameter gradients
 465 obtained using backpropagation.
 466

467 Despite the remaining future computational improvements to be made on symmetry enforcement,
 468 we reiterate the discussion in Section C with regard to computational challenges of augmentation.
 469 For m transformation types in which n hyperparameter values are selected for each, we showed that
 470 there are $m! \cdot n^m$ unique augmentations if the transformations do not commute (which is the general
 471 case). Thus, an efficient augmentation strategy of selecting random augmentations for training runs
 472 the risk of falling hopelessly short of the number of needed augmentations, while a more robust
 473 method (explicitly expanding the size of the dataset with augmentations) faces the daunting task of
 474 training on substantially larger datasets. Meanwhile, though the input gradient is, as implemented,
 475 computationally expensive, each vector field makes use of this gradient during enforcement: it need
 476 only be computed once for all of the vector fields. Therefore, an additional advantage of symmetry
 477 regularization is that it scales better as the dimension of the symmetry Lie group is increased.
 478

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648 A DERIVATION OF INFINITESIMAL GENERATORS 649

650 Here, we provide a detailed derivation of the infinitesimal generators used in our experiments. Sec-
651 tion A.1 provides a derivation for Power Law symmetry, and Section A.2 does so for Gaussian Blur.
652

653 A.1 THE POWER LAW TRANSFORM 654

655 The power law transformation, sometimes known as “gamma correction”, is a transformation in
656 which each channel value at each pixel—which value we denote as I —is adjusted according to the
657 following equation:

$$658 \quad I \rightarrow c \cdot I^\gamma, \quad (2)$$

659 where c is a (non-zero) hyperparameter Paszke et al. (2019), and where $I \in (0, 1]$. For this group
660 action, the group identity element is $\gamma = 1$: differentiating and setting $\gamma = 1$ yields an infinitesimal
661 generator at each pixel of

$$662 \quad X = cI \ln(I) \partial_I. \quad (3)$$

663 Since $cX(f) = 0$ if and only if $X(f) = 0$, an infinitesimal generator for the Power Law transfor-
664 mation acting diagonally on the RGB channels can be taken to be

$$665 \quad X = R \ln(R) \partial_R + G \ln(G) \partial_G + B \ln(B) \partial_B.$$

666 Now let (R_i, G_i, B_i) denote the RGB values at the pixel i , and suppose that each image in a dataset
667 has m pixels. The infinitesimal generator for the diagonal action of the power law transformation in
668 (R, G, B) coordinates is given as

$$669 \quad \tilde{X} = \sum_{i=1}^m R_i \ln\left(\frac{R_i}{255}\right) \partial_{R_i} + G_i \ln\left(\frac{G_i}{255}\right) \partial_{G_i} + B_i \ln\left(\frac{B_i}{255}\right) \partial_{B_i}. \quad (4)$$

672 A.2 GAUSSIAN BLUR 673

674 The Gaussian blur transformation is a transformation in which each channel value at each pixel (i, j) ,
675 denoted by I_{ij} , is replaced by a weighted linear combination of its former value and the values of
676 its neighbors. The weights of this linear combination are determined by a Gaussian function. More
677 concretely, consider the following function:

$$678 \quad f(\sigma; x, y) = \frac{1}{2\pi\sigma^2} \exp\left\{-\frac{x^2 + y^2}{2\sigma^2}\right\}. \quad (5)$$

679 Let k be the pre-determined size of the kernel, and let q be the integer quotient of k and 2. Then

$$680 \quad I_{ij} \rightarrow \sum_{s=i-q}^{i+q} \sum_{r=j-q}^{j+q} w_{sr} I_{sr}, \quad (6)$$

681 where

$$682 \quad w_{ij} = \frac{f(\sigma; i, j)}{\sum_{s=i-q}^{i+q} \sum_{r=j-q}^{j+q} f(\sigma; s, r)}.$$

683 This can be represented as a convolution of the image with pixel values I_{ij} with the kernel with
684 weights w_{ij} .

685 At first, one would wish to compute a vector field for this transformation parameterized by σ , where
686 $\sigma = 0$ corresponds to the identity transformation. However, the derivative of w_{sr} at $\sigma = 0$ is 0
687 identically: therefore, we approximate the derivative at 0 by evaluating the derivative at a nearby
688 point, namely $\sigma = 0.3$. At this point, $w_{00} \approx 0.9847$, while $w'_{00} \neq 0$. For concreteness, the matrix
689 of values this yields is given below:

$$690 \quad \begin{bmatrix} 2.46 \cdot 10^{-41} & 2.05 \cdot 10^{-29} & 2.73 \cdot 10^{-22} & 6.35 \cdot 10^{-20} & 2.73 \cdot 10^{-22} & 2.05 \cdot 10^{-29} & 2.46 \cdot 10^{-41} \\ 2.05 \cdot 10^{-29} & 1.45 \cdot 10^{-17} & 1.57 \cdot 10^{-10} & 3.25 \cdot 10^{-8} & 1.57 \cdot 10^{-10} & 1.45 \cdot 10^{-17} & 2.05 \cdot 10^{-29} \\ 2.73 \cdot 10^{-22} & 1.57 \cdot 10^{-10} & 1.08 \cdot 10^{-3} & 0.14 & 1.08 \cdot 10^{-3} & 1.57 \cdot 10^{-10} & 2.73 \cdot 10^{-22} \\ 6.35 \cdot 10^{-20} & 3.25 \cdot 10^{-8} & 0.14 & -0.56 & 0.14 & 3.25 \cdot 10^{-8} & 6.35 \cdot 10^{-20} \\ 2.73 \cdot 10^{-22} & 1.57 \cdot 10^{-10} & 1.08 \cdot 10^{-3} & 0.14 & 1.08 \cdot 10^{-3} & 1.57 \cdot 10^{-10} & 2.73 \cdot 10^{-22} \\ 2.05 \cdot 10^{-29} & 1.45 \cdot 10^{-17} & 1.57 \cdot 10^{-10} & 3.25 \cdot 10^{-8} & 1.57 \cdot 10^{-10} & 1.45 \cdot 10^{-17} & 2.05 \cdot 10^{-29} \\ 2.46 \cdot 10^{-41} & 2.05 \cdot 10^{-29} & 2.73 \cdot 10^{-22} & 6.35 \cdot 10^{-20} & 2.73 \cdot 10^{-22} & 2.05 \cdot 10^{-29} & 2.46 \cdot 10^{-41} \end{bmatrix}$$

702 B VECTOR FIELDS AND FLOWS

704 We now provide some background on vector fields and their associated flows (1-parameter trans-
 705 formations). We refer the reader to literature on the subject for additional information (Lee, 2012).
 706 Suppose that X is a smooth (tangent) vector field on \mathbb{R}^n :

$$708 X = \alpha^i \partial_{x^i} := \sum_{i=1}^n \alpha^i \partial_{x^i}, \quad (7)$$

710 where $\alpha^i : \mathbb{R}^n \rightarrow \mathbb{R}$ for $i \in [1, n]$, and where $\{x^i\}_{i=1}^n$ are coordinates on \mathbb{R}^n . X assigns a tangent
 711 vector at each point and can also be viewed as a function on the set of smooth, real-valued functions.
 712 E.g. if $f : \mathbb{R}^n \rightarrow \mathbb{R}$ is smooth,

$$714 X(f) = \sum_{i=1}^n \alpha^i \frac{\partial f}{\partial x^i}. \quad (8)$$

716 For example, for $n = 2$, if $f(x, y) = xy$ and $X = y\partial_x$, then $X(f) = y^2$. X is also a *derivation* on
 717 the set of smooth functions on \mathbb{R}^n : that is, for smooth functions $f_1, f_2 : \mathbb{R}^n \rightarrow \mathbb{R}$ and $a_1, a_2 \in \mathbb{R}$,
 718

$$719 X(a_1 f_1 + a_2 f_2) = a_1 X(f_1) + a_2 X(f_2), \quad X(f_1 f_2) = X(f_1) f_2 + f_1 X(f_2). \quad (9)$$

720 These properties are satisfied by derivatives. A flow on \mathbb{R}^n is a smooth function $\Psi : \mathbb{R} \times \mathbb{R}^n \rightarrow \mathbb{R}^n$
 721 which satisfies

$$722 \Psi(0, p) = p, \quad \Psi(s, \Psi(t, p)) = \Psi(s + t, p) \quad (10)$$

724 for all $s, t \in \mathbb{R}$ and for all $p \in \mathbb{R}^n$. A flow is a 1-parameter group of transformations. An example
 725 of a flow $\Psi : \mathbb{R} \times \mathbb{R}^2 \rightarrow \mathbb{R}^2$ is

$$726 \Psi(t, (x, y)) = (x \cos(t) - y \sin(t), x \sin(t) + y \cos(t)), \quad (11)$$

727 with t being the continuous parameter known as the flow parameter. This flow rotates a point (x, y)
 728 about the origin by t radians.

730 For a given flow Ψ , one may define a (unique) vector field X as given in Equation 8, where each
 731 function α^i is defined as

$$732 \alpha^i = \left. \left(\frac{\partial \Psi}{\partial t} \right) \right|_{t=0}. \quad (12)$$

734 Such a vector field is called the infinitesimal generator of the flow Ψ . For example, the infinitesimal
 735 generator of the flow given in Equation 11 is $-y\partial_x + x\partial_y$.

736 Conversely, given a vector field X as in Equation 8, one may define a corresponding flow as follows.
 737 Consider the following system of differential equations:

$$739 \frac{dx^i}{dt} = \alpha^i, \quad x^i(0) = x_0^i. \quad (13)$$

741 Suppose that a solution $\mathbf{x}(t)$ to Equation 13 exists for all $t \in \mathbb{R}$ and for all initial conditions $\mathbf{x}_0 \in \mathbb{R}^n$.
 742 Then the function $\Psi : \mathbb{R} \times \mathbb{R}^n \rightarrow \mathbb{R}^n$ given by

$$743 \Psi(t, \mathbf{x}_0) = \mathbf{x}(t) \quad (14)$$

745 is a flow. The infinitesimal generator corresponding to Ψ is X . For example, to calculate the flow of
 746 $-y\partial_x + x\partial_y$, we solve

$$747 \dot{x} = -y, \quad \dot{y} = x, \quad x(0) = x_0, \quad y(0) = y_0 \quad (15)$$

749 and obtain the flow $\Psi(t, (x_0, y_0))$ defined by Equation 11. It is generally easier to obtain the in-
 750 finitesimal generator of a flow than to obtain the flow of an infinitesimal generator.

751 A smooth function $f : \mathbb{R}^n \rightarrow \mathbb{R}$ is said to be X -invariant if $X(f) = 0$ identically for a smooth
 752 vector field X . The function f is Ψ -invariant if, for all $t \in \mathbb{R}$, $f = f(\Psi(t, \cdot))$ for a flow Ψ . If X is
 753 the infinitesimal generator of Ψ , f is Ψ -invariant if and only if f is X -invariant.

756 C INFINITESIMAL GENERATORS OF MULTI-PARAMETER GROUPS
757758 Let $G \in \mathbb{R}^s$ be a group, and suppose G acts on \mathbb{R}^n : that is, for $g_1, g_2 \in G$ and for $x \in \mathbb{R}^n$, there is
759 a function $\Psi : G \times \mathbb{R}^n \rightarrow \mathbb{R}^n$ such that (assuming the group operation is vector addition)
760

761
$$\Psi(\mathbf{0}, x) = x, \quad \Psi(g_2, \Psi(g_1, x)) = \Psi(g_1 + g_2, x). \quad (16)$$

762 The use of the symbol Ψ to denote a multi-parameter group action is not accident, as a flow is a
763 1-parameter group action. Let $\{v_i\}_{i=1}^s$ be a basis for the tangent space of G at $\mathbf{0}$, the group identity
764 element. Lastly, let σ be a curve in G for which $\sigma(t_0) = \mathbf{0}$ and $\dot{\sigma}(t_0) = v_i$ for $t_0 \in \mathbb{R}$. The
765 infinitesimal generator X_i corresponding to v_i is given by

766
$$X_i = \left(\frac{d}{dt} \Psi(\sigma(t), x) \right) \Big|_{t=t_0}. \quad (17)$$

769 For example, consider the group $G = \mathbb{R}^3$ acting on \mathbb{R}^2 via
770

771
$$\Psi((a, b, \theta), (x, y)) = (x \cos(\theta) - y \sin(\theta) + a, x \sin(\theta) + y \cos(\theta) + b).$$

772 Given the following three curves,
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774
$$\sigma_a(t) = (t, 0, 0), \quad \sigma_b(t) = (0, t, 0), \quad \sigma_\theta(t) = (0, 0, t),$$

775 we find that

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$$X_a = \frac{d}{dt} (t, 0) \Big|_{t=0} = \partial_x, \quad X_b = \frac{d}{dt} (0, t) \Big|_{t=0} = \partial_y,$$

777
$$X_\theta = \frac{d}{dt} (x \cos(\theta) - y \sin(\theta), x \sin(\theta) + y \cos(\theta)) \Big|_{t=0} = -y \partial_x + x \partial_y.$$

780 For each of these vector fields, a corresponding flow can be computed, which flows we call Ψ_a , Ψ_b ,
781 and Ψ_θ , respectively. In terms of the original parameters, these flows are given as
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783
$$\Psi_a(a, (x, y)) = (x + a, y), \quad \Psi_b(b, (x, y)) = (x, y + b),$$

784
$$\Psi_\theta = (x \cos(\theta) - y \sin(\theta), x \sin(\theta) + y \cos(\theta)).$$

785 While each of these flows are, individually, 1-parameter group actions, it is clear that the infinitesi-
786 mal generators X_a , X_b , and X_θ are the infinitesimal generators for the multi-parameter group action
787 given in Equation (16). Thus, discovering vector field infinitesimal generators which annihilate a
788 fixed (smooth) function does not solely apply to 1-parameter groups.
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810 D USING VECTOR FIELDS FOR DISCOVERY AND ENFORCEMENT
811812 D.1 QUANTIFYING THE EXTENT TO WHICH A SMOOTH MODEL IS INVARIANT
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814 The notion of similarity of vector fields has been previously discussed (Shaw et al., 2025), though we
815 include this discussion herein since the proposed methodology is not likely known to the machine
816 learning community at large. Given a metric tensor g_{ij} —usually assumed to be the standard Euclidean
817 metric tensor—we define the angle between two vector fields X and \hat{X} by

$$818 \cos(\theta(X, \hat{X})) = \frac{1}{\int_{\Omega} d\mathcal{M}} \mathbb{E} \left[\frac{|\langle X, \hat{X} \rangle_g|}{\|X\|_g \cdot \|\hat{X}\|_g} \right], \quad (18)$$

822 where $\langle X, \hat{X} \rangle_g = \sum_{i,j} f_i \hat{f}_j g_{ij}$, $\|X\|_g = \sqrt{\langle X, X \rangle_g}$, and where

$$824 \mathbb{E}[u(\mathbf{x})] = \int_{\Omega} u(\mathbf{x}) d\mathcal{M},$$

826 with the region Ω being defined by the range of a given dataset. Ordinarily, this is the full range of
827 the dataset. This formula is a generalization of the formula given in (Shaw et al., 2024) in the case
828 where the manifold and/or metric is not assumed to be Euclidean.

829 The notion of a “cosine similarity” between vector fields induces a method by which the extent to
830 which a particular smooth function is invariant can be quantified in relation to other functions. For
831 a fixed vector field X , a function f is X -invariant if and only if X is orthogonal to the gradient of
832 f , which vector field we denote X_f . Thus, the extent to which f is X -invariant can be quantified
833 in terms of the cosine of the angle between X and X_f , given in Equation (18): the closer to 0 this
834 value is, the more X -invariant f is.

835 D.2 CONTINUOUS SYMMETRY ENFORCEMENT USING VECTOR FIELD REGULARIZATION
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837 In this section, we summarize the previously-given method (Shaw et al., 2025; Otto et al., 2024;
838 Finzi et al., 2020) of enforcing continuous symmetries using regularization induced by vector fields.
839 We seek to extend these approaches to image data.

840 Suppose that, as in the case of supervised learning, one seeks to learn a function F mapping data
841 instances x_i to targets y_i , where $x_i \in \mathbb{R}^d$ and $y_i \in \mathbb{R}^m$. Suppose also that the function F is estimated
842 by means of the minimization of $\mathcal{L}(F(\mathbf{x}), \mathbf{y})$ for a smooth loss function \mathcal{L} . The model function F is
843 invariant with respect to the infinitesimal generators $\{X_k\}_{k=1}^s$ precisely when, for each component
844 f_j of F ,

$$845 X_k(f_j) = 0 \quad (19)$$

846 for $1 \leq k \leq s$ and $1 \leq j \leq m$. This can be used as a regularization term giving a loss function of:

$$848 (1 - \lambda(t))\mathcal{L}(F(\mathbf{x}), \mathbf{y}) + \lambda(t)\tilde{\mathcal{L}}(\mathbf{X}(F)(\mathbf{x}), \mathcal{O}), \quad (20)$$

849 where $\mathbf{X}(F)(\mathbf{x}) = (X_k(f_j))(x_i)$, \mathcal{O} is an array of zeros and is of the same shape as $\mathbf{X}(F)(\mathbf{x})$, $\tilde{\mathcal{L}}$
850 is a smooth loss function, and $\lambda(t) \in [0, 1]$ is a (possibly time/epoch-dependent) symmetry regular-
851 ization parameter.

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