CONDIFF: A CHALLENGING DATASET FOR NEURAL SOLVERS OF PARTIAL DIFFERENTIAL EQUATIONS

Anonymous authors

000

001

002 003 004

010 011

012

013

014

015

016

017

018

019

021

025

026

027 028

029

Paper under double-blind review

ABSTRACT

We present ConDiff, a novel dataset for scientific machine learning. ConDiff focuses on the parametric diffusion equation with space dependent coefficients, a fundamental problem in many applications of partial differential equations (PDEs). The main novelty of the proposed dataset is that we consider discontinuous coefficients with high contrast. These coefficient functions are sampled from a selected set of distributions. This class of problems is not only of great academic interest, but is also the basis for describing various environmental and industrial problems. In this way, ConDiff shortens the gap with real-world problems while remaining fully synthetic and easy to use. ConDiff consists of a diverse set of diffusion equations with coefficients covering a wide range of contrast levels and heterogeneity with a measurable complexity metric for clearer comparison between different coefficient functions. We baseline ConDiff on standard deep learning models in the field of scientific machine learning. By providing a large number of problem instances, each with its own coefficient function and right-hand side, we hope to encourage the development of novel physics-based deep learning approaches, such as neural operators, ultimately driving progress towards more accurate and efficient solutions of complex PDE problems.

1 INTRODUCTION

In recent years, machine learning techniques have emerged as a promising approach to solving PDEs, offering a new perspective in scientific computing. Machine learning algorithms, especially those based on neural networks, have demonstrated success in approximating complex functions and physical phenomena. Neural networks can provide more efficient and scalable methods compared to traditional numerical methods, which can be computationally expensive and limited by the dimensionality of the problem to be solved. Approaches using physical losses (Karniadakis et al., 2021), operator learning (Li et al., 2020), symmetries incorporation (Wang et al., 2020), data-driven discretization (Bar-Sinai et al., 2019) lead to more physically meaningful solutions and gave neural networks better recognition than just black-boxes.

Classical methods for solving PDEs have been extensively developed and refined over the years, 040 providing a basis for understanding and analyzing various physical phenomena. These methods 041 involve discretization the PDEs using techniques as the finite difference method (LeVeque, 2007), 042 finite element method (Bathe, 2006), finite volume method (Eymard et al., 2000) or spectral meth-043 ods (Trefethen, 2000), followed by numerical solution of the resulting algebraic equations. While 044 these methods have been successful in solving a wide range of PDEs, they often face the curse of dimensionality when parametric PDEs need to be solved in connection with optimization, optimal control, parameter identification, uncertainty quantification. The reduction of complexity for such 046 classes of problems can be addressed with surrogate models using machine learning. 047

The main approaches in scientific machine learning are (i) using governing equations as loss functions with physics-informed neural networks (Karniadakis et al., 2021; Cai et al., 2021; Eivazi et al., 2024; Raissi et al., 2019); (ii) learning mappings between infinite-dimensional function spaces with neural operators (Li et al., 2020; Fanaskov & Oseledets, 2023; Lu et al., 2021a; Li et al., 2024; Tran et al., 2021); (iii) hybrid approaches where machine learning techniques are incorporated into classical simulations (Brunton & Kutz, 2022; Schnell & Thuerey, 2024; Hsieh et al., 2019; Ingraham et al., 2018).

054 These surrogate models have shown significant potential in solving parametric PDEs, but a critical 055 aspect of their development remains the availability of comprehensive datasets for validation. The 056 accuracy and reliability of these machine learning-based approaches are highly dependent on the 057 quality and diversity of the data used to train and test them. Without such datasets, the performance 058 and generalization ability of these models cannot be adequately assessed, and their applicability to real-world problems may be limited. As new techniques and methods emerge in the future, the need for robust and extensive datasets will only increase. It is therefore essential to develop approaches to 060 the curation of high quality datasets that can support the development and validation of innovative 061 approaches to solving complex problems in different scientific and engineering domains. 062

Typically, scientific machine learning datasets have a large number of parametric PDEs (Takamoto et al., 2022; Luo et al., 2023; Hao et al., 2023) that have a single example per PDE. With ConDiff (short for Contrast Diffusion) we focused on the idea of providing a large number of different realizations for a single problem - the diffusion equation. Currently, ConDiff consists of a diverse set of diffusion equations with 24 realizations, which can be distinguished by complexity, and results in a total of 28800 samples. We also propose an approach to generating complex coefficients for parametric PDEs that can address real-world problems with a measurable metric of the complexity of the dataset.

071The ConDiff dataset is available on the Hugging Face Hub: https://huggingface.co/072datasets/condiff/ConDiff. The code with ConDiff generation, usage, validation and re-073quirements is available at: https://github.com/condiff-dataset/ConDiff.

074 075

2 CONDIFF

076 077

078 **Motivation** Creating a comprehensive benchmark for classes of parametric PDEs is a particular 079 challenge for the scientific machine learning community. The main challenges in creating a compre-080 hensive dataset are: (i) computational complexity; (ii) storage complexity for the desired dimensions 081 of the discretized PDE and parameter space; (iii) properties of the coefficients and solution functions; (iv) relation to real-world problems. The first and second reasons illustrate a technical bottleneck in 083 the creation of the dataset and are mostly dependent on the hardware and efficiency of the numerical method used. Properties such as coefficient smoothness, discontinuity, spatial variation of the 084 coefficients, variance of the parametric space significantly affect the complexity of the dataset and 085 should be carefully chosen. The solution to parametric PDEs (i.e. the ground truth for the dataset) depends on a number of numerical aspects such as choice of mesh, discretization, numerical algo-087 rithm, boundary and initial conditions. Therefore, it is very important to consider every little detail 088 regarding different numerical schemes, PDEs, boundary and initial conditions. 089

Existing benchmarks and datasets cover different aspects of scientific machine learning for different 090 classes of PDEs and can be divided into several groups. PDEBench (Takamoto et al., 2022), PIN-091 Nacle (Hao et al., 2023), CFDBench (Luo et al., 2023) have a large number of PDEs with different 092 boundary and initial conditions and different dimensionality and resolution. The best covered area is weather forecasting: SuperBench (Ren et al., 2023), ClimSim (Yu et al., 2024), DynaBench (Dulny 094 et al., 2023), OceanBench (Johnson et al., 2024), ChaosBench (Nathaniel et al., 2024). There are 095 also domain specific datasets with applications to Lagrangian mechanics LagrangeBench (Toshev 096 et al., 2024) and phase change phenomena BubbleML (Hassan et al., 2023). Recently, the Flow-097 Bench (Tali et al., 2024) dataset with complex geometries was introduced. Worth noting frameworks 098 for differential simulations and general environments for PDEs in scientific machine learning: PDE Control Gym (Bhan et al., 2024), PDEArena (Gupta & Brandstetter, 2022), DiffTaichi (Hu et al., 099 2019), DeepXDE (Lu et al., 2021b) and Φ_{Flow} (Holl et al., 2020). 100

While all of these datasets contribute significantly to the community, to the best of the authors' knowledge there is no dataset dedicated to the very important class of academic and real-world problems, the class of parametric PDEs with random coefficients. Typically, when a new model is proposed, authors test it with a set of equations with smooth coefficients (Brandstetter et al., 2022; Nguyen et al., 2023; Ripken et al., 2023; Bryutkin et al., 2024). Such coefficients do not allow important classes of industrial applications to be addressed. In section 3 we show that increasing the heterogeneity and contrast in the coefficient function leads to increasing challenges in building accurate surrogate models.

Problem definition Existing benchmarks (Takamoto et al., 2022; Hao et al., 2023; Luo et al., 2023) cover a set of PDEs, both steady-state and time-dependent, with different resolutions and time lengths. In our work, we approach the problem from the other side tacking a fixed parametric PDE and generating a comprehensive set of random coefficients for it. We consider a 2D steady-state diffusion equation:

 $-\nabla \cdot (k(x)\nabla u(x)) = f(x), \text{ in } \Omega$ $u(x)\Big|_{x \in \partial \Omega} = 0 \qquad (1)$

Note that the equation (1) models not only diffusion, but also steady-state Darcy flow in porous media, steady-state heat conduction, etc. To address certain real-world problems, we use the Gaussian Random Field (GRF) to generate the field $\phi(x)$ (Figure 1) with the following covariance models as functions of distance d:

• Cubic:

113 114

115

116

117 118

123

132

133 134

138 139 140

141

142 143 144

145

$$\operatorname{Cov}(d) = \begin{cases} \sigma^2 \left(1 - 7 \left(\frac{d}{l} \right)^2 + \frac{35}{4} \left(\frac{d}{l} \right)^3 - \frac{7}{2} \left(\frac{d}{l} \right)^5 + \frac{3}{4} \left(\frac{d}{l} \right)^7 \right), & d < l \\ 0, & d \ge l \end{cases}.$$
(2)

• Exponential:

$$\operatorname{Cov}(d) = \sigma^2 \exp\left(-\frac{d}{l}\right).$$
 (3)

• Gaussian:

$$\operatorname{Cov}(d) = \sigma^2 \exp\left(-\frac{d^2}{l^2}\right). \tag{4}$$

The correlation length in each dataset is l = 0.05 and the complexity of a resulting dataset is controlled by variance σ^2 . The forcing term f(x) is sampled from the standard normal distribution for each sampled PDE in each dataset. The resulting coefficient k(x) is obtained with:

$$k(x) = \exp\left(\phi(x)\right). \tag{5}$$

We propose to measure the complexity of the generated GRF with the global contrast in the field $\phi(x)$:

contrast = exp
$$\left(\max \left(\phi(x) \right) - \min \left(\phi(x) \right) \right)$$
. (6)

146 **Complexity grows with variance** By increasing the variance σ^2 one can obtain a higher contrast (6) and thus a higher complexity of the PDE. This is a well-known phenomenon in applied numerical analysis and can be easily observed empirically. We illustrate this behaviour with the condition number $\kappa(A)$ of the matrices A obtained with discretization of the equation (1).

150 In the Table 1 one can observe that increasing σ^2 leads to a higher condition number $\kappa(A)$ = 151 $|\lambda_{\rm max}|/|\lambda_{\rm min}|$ of the discretized differential operator (Capizzano, 2003). The condition number is 152 closely related to the performance of the numerical methods used to solve PDEs (Benzi et al., 2005; 153 Elman et al., 2014). A high condition number indicates that small changes in the input can lead to 154 large changes in the output, making the problem ill-conditioned. This is particularly important in 155 PDEs, where small perturbations can significantly affect the solution. Also, if iterative methods are 156 used to solve the discretized PDE, a larger condition number means a larger number of iterations for 157 unpreconditioned and most of preconditioned iterative methods (Saad, 2003).

158

Connection to real-world problems All of the above reasoning is done with regard to the frequent occurrence of such tasks in real world (Hashmi, 2014; Massimo, 2013; Carr & Turner, 2016; Oristaglio & Hohmann, 1984; Muravleva et al., 2021), including composite materials modeling, heat transfer, geophysical problems, fluid flow modeling. In Figure 2 one can see a cross section of the



Figure 1: Visualization of the GRF (top row), the coefficient k(x) generated from this GRF (middle row) and the corresponding solution of the equation (1) (bottom row) for a sampled PDEs with grid 128×128 and $\sigma^2 = 2.0$.



Figure 2: Cross section of the x-permeability field along the z axis over the SPE10 model 2 with z = 4.

x-permeability field along the z axis over the SPE10 model 2 benchmark (Christie & Blunt, 2001). 215 The term permeability is used to denote the coefficients of the above equation when considering flow in porous media. This field is very similar to the ConDiff samples in Figure 1.

This benchmark is well known in the field of reservoir modelling and fluid flow in porous media.
 SPE10 model 2 poses a significant challenge for the tasks of uncertainty quantification, upscaling and multiphase fluid flow modelling.

Table 1: Summary of the ConDiff with min, mean and max values of the contrast (6). ¹Condition number $\kappa(A)$ is calculated for a single sampled discretized (1).

Covariance	Variance	Min contrast	Mean contrast	Max contrast	$\kappa^1(A)$
		Grid	64×64		
	0.1	$7.0 \cdot 10^{0}$	$1.0\cdot 10^1$	$1.5\cdot 10^1$	$3.6 \cdot 10^{3}$
Cubia	0.4	$5.0\cdot 10^1$	$9.6\cdot 10^1$	$2.5\cdot 10^2$	$7.3 \cdot 10^3$
Cubic	1.0	$6.0\cdot 10^2$	$8.3\cdot 10^2$	$1.0\cdot 10^3$	$2.0 \cdot 10^{4}$
	2.0	$8.0\cdot 10^4$	$8.9\cdot 10^4$	$1.0\cdot 10^5$	$1.8 \cdot 10^{5}$
	0.1	$6.0\cdot 10^0$	$9.0\cdot 10^0$	$1.5\cdot 10^1$	$4.3 \cdot 10^3$
Evn	0.4	$5.0\cdot 10^1$	$8.5 \cdot 10^1$	$2.3 \cdot 10^2$	$5.2 \cdot 10^{3}$
	1.0	$6.0\cdot 10^2$	$7.9\cdot 10^2$	$1.0\cdot 10^3$	$1.7 \cdot 10^{4}$
	2.0	$8.0 \cdot 10^4$	$8.9\cdot 10^4$	$1.0\cdot 10^5$	$1.9 \cdot 10^{5}$
Gauss	0.1	$5.0 \cdot 10^0$	$8.0\cdot 10^0$	$1.4 \cdot 10^1$	$4.1 \cdot 10^{3}$
	0.4	$5.0\cdot 10^1$	$7.5\cdot 10^1$	$2.3\cdot 10^2$	$8.1 \cdot 10^{3}$
	1.0	$6.0\cdot 10^2$	$7.7\cdot 10^2$	$1.0\cdot 10^3$	$2.4 \cdot 10^{4}$
	2.0	$8.0\cdot 10^4$	$8.9\cdot 10^4$	$1.0\cdot 10^5$	$8.8 \cdot 10^{5}$
		Grid 1	$.28 \times 128$		
	0.1	$8.0\cdot 10^0$	$1.1 \cdot 10^1$	$1.5\cdot 10^1$	$1.6 \cdot 10^{4}$
Cubic	0.4	$5.5\cdot 10^1$	$1.3\cdot 10^2$	$2.5\cdot 10^2$	$3.8\cdot10^4$
Cubic	1.0	$6.0\cdot 10^2$	$8.8\cdot 10^2$	$1.0\cdot 10^3$	$1.0 \cdot 10^{5}$
	2.0	$8.0\cdot 10^4$	$8.9\cdot 10^4$	$1.0\cdot 10^5$	$1.2 \cdot 10^{6}$
	0.1	$6.0\cdot 10^0$	$1.0\cdot 10^1$	$1.5 \cdot 10^1$	$1.7 \cdot 10^{4}$
Exp	0.4	$5.1\cdot 10^1$	$1.1\cdot 10^2$	$2.5\cdot 10^2$	$3.3\cdot 10^4$
	1.0	$6.0\cdot 10^2$	$8.3\cdot 10^2$	$1.0\cdot 10^3$	$9.7\cdot 10^4$
	2.0	$8.0\cdot 10^4$	$8.9\cdot 10^4$	$1.0\cdot 10^5$	$6.3\cdot10^5$
Gauss	0.1	$5.0 \cdot 10^0$	$8.0 \cdot 10^0$	$1.4\cdot 10^1$	$1.8 \cdot 10^{4}$
	0.4	$5.0\cdot10^1$	$7.8\cdot 10^1$	$2.5\cdot 10^2$	$7.2 \cdot 10^{4}$
	1.0	$6.0\cdot 10^2$	$7.7\cdot 10^2$	$1.0\cdot 10^3$	$1.6 \cdot 10^{5}$
	2.0	$8.0\cdot 10^4$	$8.9\cdot 10^4$	$1.0\cdot 10^5$	$1.5 \cdot 10^{6}$

Dataset description To generate the fields $\phi(x)$ we use the highly efficient parafields library¹ with C++ backend. We use covariance models from {cubic, exponential, Gaussian} with 4 variance values from $\{0.1, 0.4, 1.0, 2.0\}$. We use the forcing term $f(x) \sim \mathcal{N}(0, 1)$. The standard normal force function is chosen to be more complex than a constant forcing term, but not too complex to distract from the complex coefficients, which is the focus of ConDiff. A Dirichlet boundary condition is set for each coefficient realization since boundary conditions do not contribute significantly to the resulting complexity (Capizzano, 2003). The ground truth solution is obtained using cell-centered second-order finite volume method. The coefficients are in the center of cells, the values are in the nodes.

For each parameter set, we generate 1000 training and 200 test realizations of the diffusion equation (1) on 64×64 and 128×128 grids. We provide the train-test split in the ConDiff for fair comparison in future research papers. Note that datasets with the same field parameters but different grid sizes are generated independently and do not represent the same field. The fixed geometry of ConDiff allows PDEs with different fields $\phi(x)$ to be compared without fear that different ge-

¹https://github.com/parafields/parafields

ometries will interfere with a fair comparison across different coefficient functions. To control the complexity of the generated PDEs realizations, we set contrast bounds during generation as follows:

272 273 274

275

276 277

278 279

280

281

282

283

284

298

299 300 301

302

303

- $\sigma^2 = 0.1$, contrast $\in [5, 15]$,
- $\sigma^2 = 0.4$, contrast $\in [50, 250]$,
- $\sigma^2 = 1.0$, contrast $\in [6 \cdot 10^2, 10^3]$,
- $\sigma^2 = 2.0$, contrast $\in [8 \cdot 10^4, 10^5]$.

In total, ConDiff consists of 24 PDEs with different GRFs and grid sizes. Table 1 summarizes the properties of ConDiff. Figure 3 illustrates the contrast distributions. Coming back to the permeability cross section of SPE10 model 2 (Figure 2), it has contrast $= 2.5 \cdot 10^6$ according to (6). We want to emphasize that although the most complex coefficient of ConDiff is smaller by an order of magnitude compared to the cross section of SPE10 model 2, our experiments show that this coefficient is too complex for the chosen models to predict well.



Figure 3: GRF contrast distribution for PDEs from Table 1.

3 EXPERIMENTS

304 **Models** We do not attempt to benchmark every scientific machine learning surrogate model on the 305 ConDiff. Since the ConDiff consists of triplets (k(x), f(x), u(x)), its primary use is to validate dif-306 ferent architectures of neural operators. Therefore, we have selected the following list of models to validate on the ConDiff: Spectral Neural Operator (SNO) (Fanaskov & Oseledets, 2023), Factorized 307 Fourier Neural Operator (F-FNO) (Tran et al., 2021), Dilated ResNet (DilResNet) (Yu et al., 2017) 308 and U-Net (Ronneberger et al., 2015). Neural operators FNO and SNO are both types of neural 309 networks designed to learn mappings between function spaces, in particular to solve PDEs. Neural 310 operators are designed to be universal approximators of continuous operators acting between Ba-311 nach spaces and to be discretization invariant, meaning that they can handle different discretizations 312 of the underlying function spaces without requiring changes to the model. DilResNet and U-Net are 313 classical neural network models originating from the field of computer vision (CV). Both models 314 have shown their applicability beyond CV and have been used extensively for modeling physical 315 phenomena (Stachenfeld et al., 2021; Ma et al., 2021). More details about the models used can be 316 found in the Appendix A.1.

317

Experiment environment For training neural networks we use frameworks from the JAX (Bradbury et al., 2018) ecosystem: Equinox (Kidger & Garcia, 2021) and Optax (DeepMind et al., 2020). The loss function used is the relative L_2 loss:

- 321
- 322

$$L_2 = \frac{1}{N} \sum_{i=1}^{N} \frac{\|\hat{y}_i - y_i\|_2}{\|y_i\|_2} \,. \tag{7}$$

Training samples for the models are the values of the coefficient function k(x) and the forcing term f(x) in the grid cells. Targets are the values of the solution function u(x) in the grid cells. We also use (7) as a primary performance metric, assessing the quality of the models' predictions, and report averaged values over the test set with standard deviation.

For all the problems we train for 400 epochs for grid = 64 and for 500 epochs for grid = 128. We use the AdamW optimizer with an initial learning rate equals to 10^{-3} and a weight decay equals to 10^{-2} . We use a learning rate schedule that halves the learning rate every 50 epochs. Each PDE realization has a dataset size of 1000 training samples and 200 test samples. We use a single GPU Nvidia Tesla V100 16Gb for training on grid = 64 and a single GPU Nvidia A40 48Gb for training on grid = 128.

Table 2: Results for Poisson equation.

Grid	SNO	F-FNO	DilResNet	U-Net
64	0.056 ± 0.018	0.027 ± 0.008	0.018 ± 0.005	0.020 ± 0.007
128	0.073 ± 0.021	0.047 ± 0.013	0.063 ± 0.016	0.267 ± 0.049

Table 3: Performance comparison of the models on the PDEs with the 64×64 grid from ConDiff.

Covariance	Variance	SNO	F-FNO	DilResNet	U-Net
	0.1	0.09 ± 0.02	0.07 ± 0.02	0.07 ± 0.02	0.08 ± 0.02
Cubia	0.4	0.15 ± 0.04	0.14 ± 0.03	0.14 ± 0.03	0.17 ± 0.04
Cubic	1.0	0.23 ± 0.06	0.22 ± 0.06	0.22 ± 0.06	0.24 ± 0.06
	2.0	0.35 ± 0.10	0.34 ± 0.09	0.35 ± 0.10	0.42 ± 0.10
	0.1	0.12 ± 0.03	0.11 ± 0.03	0.11 ± 0.03	0.12 ± 0.04
Evn	0.4	0.21 ± 0.06	0.21 ± 0.06	0.20 ± 0.06	0.26 ± 0.07
Ехр	1.0	0.33 ± 0.09	0.34 ± 0.09	0.36 ± 0.09	0.35 ± 0.09
	2.0	0.59 ± 0.14	0.58 ± 0.14	0.60 ± 0.13	0.64 ± 0.13
Gauss	0.1	0.12 ± 0.04	0.11 ± 0.04	0.11 ± 0.03	0.12 ± 0.03
	0.4	0.23 ± 0.06	0.22 ± 0.06	0.21 ± 0.06	0.25 ± 0.06
	1.0	0.38 ± 0.08	0.37 ± 0.09	0.38 ± 0.09	0.39 ± 0.09
	2.0	0.66 ± 0.14	0.65 ± 0.14	0.66 ± 0.13	0.72 ± 0.24

Table 4: Performance comparison of SNO and F-FNO on the PDEs with the 128×128 grid from ConDiff.

Covariance	Variance	SNO	F-FNO
	0.1	0.09 ± 0.03	0.08 ± 0.02
Cubic	0.4	0.15 ± 0.04	0.14 ± 0.04
Cubic	1.0	0.23 ± 0.06	0.22 ± 0.06
	2.0	0.36 ± 0.11	0.36 ± 0.10
	0.1	0.13 ± 0.03	0.12 ± 0.03
Eve	0.4	0.21 ± 0.07	0.21 ± 0.06
Ехр	1.0	0.33 ± 0.09	0.33 ± 0.08
	2.0	0.58 ± 0.15	0.57 ± 0.13
	0.1	0.13 ± 0.04	0.12 ± 0.03
Cause	0.4	0.23 ± 0.06	0.23 ± 0.06
Jauss	1.0	0.37 ± 0.10	0.37 ± 0.10
	2.0	0.68 ± 0.13	0.66 ± 0.13

Validation on ConDiff We start the experiments with the Poisson equation and consider it as a special case of (1) with k(x) = 1 and contrast = 1. All models achieve an accuracy of the order of 10^{-2} (Table 2). Increasing the grid size leads to moderate increases in error, except for the U-Net for which the error increases by an order of magnitude.

The diffusion equation for grid 64 (Table 3) with covariances (2), (3) and (4) are more challenging for the models. While the performance on the diffusion equation with cubic covariance with $\sigma^2 = 0.1$ is comparable to the performance on the Poisson equation, the error on the diffusion equation with exponential and Gaussian covariances is already an order of magnitude higher. Increasing σ^2 leads to worse performance of each model on each PDE. The most complex PDE is the one generated with the Gaussian covariance model in GRF, which is also consistent with the condition number estimation in Table 1. Interestingly, the performance of FNO and SNO models on PDEs with grid 128 is not much different from PDEs on grid 64 (Table 4).

Table 5: Generalization of the models to unseen PDEs with different GRF covariance model with 64×64 grid and $\sigma^2 = 0.1$.

		SNO			F-FNO	
Train\Test	Cubic	Exp	Gauss	Cubic	Exp	Gauss
Cubic	0.09 ± 0.02	0.12 ± 0.04	0.12 ± 0.03	0.07 ± 0.02	0.11 ± 0.03	0.11 ± 0.03
Exp	0.09 ± 0.03	0.12 ± 0.03	0.12 ± 0.04	0.08 ± 0.03	0.11 ± 0.03	0.11 ± 0.04
Gauss	0.09 ± 0.03	0.12 ± 0.03	0.12 ± 0.03	0.08 ± 0.02	0.11 ± 0.03	0.11 ± 0.04
		DilResNet			U-Net	
	Cubic	Exp	Gauss	Cubic	Exp	Gauss
Cubic	0.07 ± 0.02	0.11 ± 0.04	0.11 ± 0.03	0.08 ± 0.02	0.12 ± 0.03	0.12 ± 0.03
Exp	0.07 ± 0.02	0.11 ± 0.03	0.11 ± 0.03	0.08 ± 0.03	0.11 ± 0.03	0.11 ± 0.04
Gauss	0.17 ± 0.06	0.25 ± 0.09	0.11 ± 0.04	0.08 ± 0.02	0.12 ± 0.04	0.12 ± 0.03

Table 6: Generalization of the models to unseen PDEs with different GRF covariance model with 64×64 grid and $\sigma^2 = 0.4$.

Train\ Test	Cubic						
main \ rest	Cubic	Exp	Gauss	Cubic	Exp	Gauss	
Cubic	0.15 ± 0.04	0.22 ± 0.06	0.22 ± 0.07	0.14 ± 0.03	0.21 ± 0.06	0.21 ± 0.07	
Exp	0.18 ± 0.05	0.21 ± 0.06	0.22 ± 0.06	0.15 ± 0.04	0.21 ± 0.06	0.22 ± 0.07	
Gauss	0.17 ± 0.05	0.22 ± 0.06	0.23 ± 0.07	0.15 ± 0.04	0.21 ± 0.07	0.22 ± 0.06	
		DilResNet			U-Net		
	Cubic	Exp	Gauss	Cubic	Exp	Gauss	
Cubic	0.14 ± 0.04	0.23 ± 0.07	0.23 ± 0.07	0.17 ± 0.06	0.24 ± 0.07	0.24 ± 0.07	
Exp	0.14 ± 0.04	0.20 ± 0.06	0.22 ± 0.06	0.23 ± 0.08	0.26 ± 0.07	0.27 ± 0.08	
Gauss	0.30 ± 0.10	0.24 ± 0.07	0.21 ± 0.06	0.21 ± 0.06	0.27 ± 0.08	0.26 ± 0.07	

Transfer between parametric spaces Ideally, the surrogate model should handle transfers be-tween different underlying parametric spaces of PDEs without loss of quality. In Tables 5, 6, 7, 8 show that in most experiments the error increases when training on cubic GRF and inferencing on exponential and Gaussian GRF. Conversely, the error decreases when training on Gaussian GRF and inferencing on cubic GRF.

4	3	2
4	3	3

Table 7: Generalization of the models to unseen PDEs with different GRF covariance model with 64×64 grid and $\sigma^2 = 1.0$.

		SNO			F-FNO	
Train\Test	Cubic	Exp	Gauss	Cubic	Exp	Gauss
Cubic	0.23 ± 0.06	0.35 ± 0.09	0.39 ± 0.09	0.22 ± 0.06	0.34 ± 0.09	0.37 ± 0.09
Exp	0.25 ± 0.06	0.33 ± 0.09	0.38 ± 0.09	0.24 ± 0.06	0.34 ± 0.09	0.38 ± 0.09
Gauss	0.24 ± 0.07	0.35 ± 0.09	0.38 ± 0.08	0.24 ± 0.06	0.35 ± 0.09	0.37 ± 0.09
		DilResNet			U-Net	
	Cubic	Exp	Gauss	Cubic	Exp	Gauss
Cubic	0.22 ± 0.06	0.35 ± 0.09	0.38 ± 0.09	0.24 ± 0.06	0.36 ± 0.09	0.38 ± 0.08
Exp	0.25 ± 0.07	0.36 ± 0.09	0.38 ± 0.10	0.25 ± 0.07	0.35 ± 0.09	0.38 ± 0.10
Gauss	0.57 ± 0.22	0.59 ± 0.22	0.38 ± 0.09	0.27 ± 0.07	0.36 ± 0.11	0.39 ± 0.09

Table 8: Generalization of the models to unseen PDEs with different GRF covariance model with 64×64 grid and $\sigma^2 = 2.0$.

		SNO			F-FNO	
Train\Test	Cubic	Exp	Gauss	Cubic	Exp	Gauss
Cubic	0.35 ± 0.10	0.60 ± 0.14	0.70 ± 0.26	0.34 ± 0.09	0.61 ± 0.14	0.67 ± 0.19
Exp	0.39 ± 0.11	0.59 ± 0.14	0.69 ± 0.24	0.39 ± 0.11	0.58 ± 0.14	0.66 ± 0.15
Gauss	0.40 ± 0.11	0.60 ± 0.13	0.66 ± 0.14	0.37 ± 0.11	0.60 ± 0.13	0.65 ± 0.14
		DilResNet			U-Net	
	Cubic	Exp	Gauss	Cubic	Exp	Gauss
Cubic	0.35 ± 0.10	0.61 ± 0.14	0.66 ± 0.15	0.42 ± 0.10	0.65 ± 0.14	0.68 ± 0.14
Exp	0.41 ± 0.10	0.60 ± 0.13	0.66 ± 0.17	0.53 ± 0.18	0.64 ± 0.13	0.72 ± 0.16
Gauss	0.72 ± 0.50	0.68 ± 0.20	0.66 ± 0.13	0.66 ± 0.40	0.69 ± 0.16	0.72 ± 0.24

4 DISCUSSION

We propose a novel dataset for the field of neural solving of parametric PDEs. The unique feature
of the dataset is discontinuous coefficients with high contrast for parametric PDEs from different
distributions. By designing the coefficients in this way, we achieve a high complexity of the generated PDEs, which also illustrates real-world problems. The proposed complexity function allows to
distinguish between the generated PDEs. We also provide code to generate new data based on the
approach used in this paper. Furthermore, we validate a number of surrogate models on the ConDiff
to illustrate its usefulness in the field of scientific machine learning.

The practical use of ConDiff is straightforward: it should be used for novel deep learning models
and approaches for modeling solution of parametric PDEs from their coefficients. Ultimately, novel
deep learning models should exhibit machine-precision prediction quality and not degrade with
increasing contrast.

It should be noted that the problems considered in this paper belong to the class of stochastic PDEs.
The equation (1) has to be solved for a very large number of sampled coefficients when Monte
Carlo or other methods are used to solve the stochastic PDEs. The surrogate models can help to
significantly reduce the computational burden, so embedding the surrogate models tested on ConDiff
into a Monte Carlo or similar stochastic PDEs solver is a reasonable next step.

486 5 LIMITATIONS

488 Limitations of the proposed dataset are: 489 490 1. For practical numerical analysis, ConDiff is generated with small and moderate variances. 491 The case of large variances has to be studied separately. 492 2. A linear elliptic parametric PDE is the basis of ConDiff, so other high contrast datasets are 493 needed to test surrogate models for hyperbolic PDEs, nonlinear problems, etc. 494 3. ConDiff is generated on a regular rectangular grid. Other meshes and geometries may 495 be required as an evolution of ConDiff. This may require more complex computational 496 methods to obtain the ground truth solution. 497 44. The forcing term f(x) is sampled from the standard normal distributions. While in this pa-498 per we focus on the complexity arising from discontinuous coefficients with high contrast, 499 the right-hand side of a PDE can also significantly affect the complexity of the solving 500 PDE. The case of complex forcing terms has to be studied separately. 501 REFERENCES 504 Yohai Bar-Sinai, Stephan Hoyer, Jason Hickey, and Michael P Brenner. Learning data-driven dis-505 cretizations for partial differential equations. Proceedings of the National Academy of Sciences, 506 116(31):15344-15349, 2019. 507 Klaus-Jürgen Bathe. Finite element procedures. Klaus-Jurgen Bathe, 2006. 509 Michele Benzi, Gene H Golub, and Jörg Liesen. Numerical solution of saddle point problems. Acta 510 numerica, 14:1-137, 2005. 511 512 Luke Bhan, Yuexin Bian, Miroslav Krstic, and Yuanyuan Shi. Pde control gym: A benchmark for 513 data-driven boundary control of partial differential equations. arXiv preprint arXiv:2405.11401, 514 2024. 515 James Bradbury, Roy Frostig, Peter Hawkins, Matthew James Johnson, Chris Leary, Dougal 516 Maclaurin, George Necula, Adam Paszke, Jake VanderPlas, Skye Wanderman-Milne, and Qiao 517 Zhang. JAX: composable transformations of Python+NumPy programs, 2018. URL http: 518 //github.com/google/jax. 519 Johannes Brandstetter, Daniel Worrall, and Max Welling. Message passing neural pde solvers. arXiv 521 preprint arXiv:2202.03376, 2022. 522 523 Steven L Brunton and J Nathan Kutz. Data-driven science and engineering: Machine learning, dynamical systems, and control. Cambridge University Press, 2022. 524 Andrey Bryutkin, Jiahao Huang, Zhongying Deng, Guang Yang, Carola-Bibiane Schönlieb, and An-526 gelica Aviles-Rivero. Hamlet: Graph transformer neural operator for partial differential equations. 527 *arXiv preprint arXiv:2402.03541*, 2024. 528 529 Shengze Cai, Zhiping Mao, Zhicheng Wang, Minglang Yin, and George Em Karniadakis. Physics-530 informed neural networks (pinns) for fluid mechanics: A review. Acta Mechanica Sinica, 37(12): 531 1727–1738, 2021. 532 S Serra Capizzano. Generalized locally toeplitz sequences: spectral analysis and applications to discretized partial differential equations. Linear Algebra and its Applications, 366:371-402, 2003. 534 535 EJ Carr and IW Turner. A semi-analytical solution for multilayer diffusion in a composite medium 536 consisting of a large number of layers. Applied Mathematical Modelling, 40(15-16):7034-7050, 2016. 538 Michael Andrew Christie and Martin J Blunt. Tenth spe comparative solution project: A comparison of upscaling techniques. SPE Reservoir Evaluation & Engineering, 4(04):308–317, 2001.

556

558

559

560

561

562

565

574

575

576

579

585

- 540 DeepMind, Igor Babuschkin, Kate Baumli, Alison Bell, Surya Bhupatiraju, Jake Bruce, Peter 541 Buchlovsky, David Budden, Trevor Cai, Aidan Clark, Ivo Danihelka, Antoine Dedieu, Clau-542 dio Fantacci, Jonathan Godwin, Chris Jones, Ross Hemsley, Tom Hennigan, Matteo Hessel, 543 Shaobo Hou, Steven Kapturowski, Thomas Keck, Iurii Kemaev, Michael King, Markus Kunesch, 544 Lena Martens, Hamza Merzic, Vladimir Mikulik, Tamara Norman, George Papamakarios, John Quan, Roman Ring, Francisco Ruiz, Alvaro Sanchez, Laurent Sartran, Rosalia Schneider, Eren 545 Sezener, Stephen Spencer, Srivatsan Srinivasan, Miloš Stanojević, Wojciech Stokowiec, Luyu 546 Wang, Guangyao Zhou, and Fabio Viola. The DeepMind JAX Ecosystem, 2020. URL http: 547 //github.com/google-deepmind. 548
- Andrzej Dulny, Andreas Hotho, and Anna Krause. Dynabench: A benchmark dataset for learning dynamical systems from low-resolution data. In *Joint European Conference on Machine Learning and Knowledge Discovery in Databases*, pp. 438–455. Springer, 2023.
- Hamidreza Eivazi, Yuning Wang, and Ricardo Vinuesa. Physics-informed deep-learning applica tions to experimental fluid mechanics. *Measurement science and technology*, 35(7):075303, 2024.
 - Howard C Elman, David J Silvester, and Andrew J Wathen. *Finite elements and fast iterative solvers: with applications in incompressible fluid dynamics.* Oxford university press, 2014.
 - Robert Eymard, Thierry Gallouët, and Raphaèle Herbin. Finite volume methods. *Handbook of numerical analysis*, 7:713–1018, 2000.
 - VS Fanaskov and Ivan V Oseledets. Spectral neural operators. In *Doklady Mathematics*, volume 108, pp. S226–S232. Springer, 2023.
- Jayesh K Gupta and Johannes Brandstetter. Towards multi-spatiotemporal-scale generalized pde
 modeling. *arXiv preprint arXiv:2209.15616*, 2022.
- Zhongkai Hao, Jiachen Yao, Chang Su, Hang Su, Ziao Wang, Fanzhi Lu, Zeyu Xia, Yichi Zhang,
 Songming Liu, Lu Lu, et al. Pinnacle: A comprehensive benchmark of physics-informed neural
 networks for solving pdes. *arXiv preprint arXiv:2306.08827*, 2023.
- 569 M Saleem J Hashmi. *Comprehensive materials processing*. Newnes, 2014.
- Sheikh Md Shakeel Hassan, Arthur Feeney, Akash Dhruv, Jihoon Kim, Youngjoon Suh, Jaiyoung
 Ryu, Yoonjin Won, and Aparna Chandramowlishwaran. Bubbleml: A multi-physics dataset and
 benchmarks for machine learning. *arXiv preprint arXiv:2307.14623*, 2023.
 - Philipp Holl, Vladlen Koltun, and Nils Thuerey. Learning to control pdes with differentiable physics. arXiv preprint arXiv:2001.07457, 2020.
- Jun-Ting Hsieh, Shengjia Zhao, Stephan Eismann, Lucia Mirabella, and Stefano Ermon. Learning
 neural pde solvers with convergence guarantees. *arXiv preprint arXiv:1906.01200*, 2019.
- Yuanming Hu, Luke Anderson, Tzu-Mao Li, Qi Sun, Nathan Carr, Jonathan Ragan-Kelley, and
 Frédo Durand. Difftaichi: Differentiable programming for physical simulation. *arXiv preprint arXiv:1910.00935*, 2019.
- John Ingraham, Adam Riesselman, Chris Sander, and Debora Marks. Learning protein structure
 with a differentiable simulator. In *International conference on learning representations*, 2018.
- J Emmanuel Johnson, Quentin Febvre, Anastasiia Gorbunova, Sam Metref, Maxime Ballarotta,
 Julien Le Sommer, et al. Oceanbench: The sea surface height edition. Advances in Neural
 Information Processing Systems, 36, 2024.
- George Em Karniadakis, Ioannis G Kevrekidis, Lu Lu, Paris Perdikaris, Sifan Wang, and Liu Yang.
 Physics-informed machine learning. *Nature Reviews Physics*, 3(6):422–440, 2021.

591
592 Patrick Kidger and Cristian Garcia. Equinox: neural networks in JAX via callable PyTrees and filtered transformations. *Differentiable Programming workshop at Neural Information Processing Systems 2021*, 2021.

625

- Randall J LeVeque. Finite difference methods for ordinary and partial differential equations: steadystate and time-dependent problems. SIAM, 2007.
- Zongyi Li, Nikola Kovachki, Kamyar Azizzadenesheli, Burigede Liu, Kaushik Bhattacharya, Andrew Stuart, and Anima Anandkumar. Fourier neural operator for parametric partial differential equations. *arXiv preprint arXiv:2010.08895*, 2020.
- Zongyi Li, Nikola Kovachki, Chris Choy, Boyi Li, Jean Kossaifi, Shourya Otta, Mohammad Amin
 Nabian, Maximilian Stadler, Christian Hundt, Kamyar Azizzadenesheli, et al. Geometry informed neural operator for large-scale 3d pdes. *Advances in Neural Information Processing Systems*, 36, 2024.
- Lu Lu, Pengzhan Jin, Guofei Pang, Zhongqiang Zhang, and George Em Karniadakis. Learning nonlinear operators via deeponet based on the universal approximation theorem of operators. *Nature machine intelligence*, 3(3):218–229, 2021a.
- Lu Lu, Xuhui Meng, Zhiping Mao, and George Em Karniadakis. DeepXDE: A deep learn ing library for solving differential equations. *SIAM Review*, 63(1):208–228, 2021b. doi:
 10.1137/19M1274067.
- Yining Luo, Yingfa Chen, and Zhen Zhang. Cfdbench: A comprehensive benchmark for machine learning methods in fluid dynamics. *arXiv preprint arXiv:2310.05963*, 2023.
- Hao Ma, Yuxuan Zhang, Nils Thuerey, Xiangyu Hu, and Oskar J Haidn. Physics-driven learning
 of the steady navier-stokes equations using deep convolutional neural networks. *arXiv preprint arXiv:2106.09301*, 2021.
- Luigi Massimo. *Physics of high-temperature reactors*. Elsevier, 2013.
- Ekaterina A Muravleva, Dmitry Yu Derbyshev, Sergei A Boronin, and Andrei A Osiptsov. Multigrid
 pressure solver for 2d displacement problems in drilling, cementing, fracturing and eor. *Journal of Petroleum Science and Engineering*, 196:107918, 2021.
- Juan Nathaniel, Yongquan Qu, Tung Nguyen, Sungduk Yu, Julius Busecke, Aditya Grover, and
 Pierre Gentine. Chaosbench: A multi-channel, physics-based benchmark for subseasonal-to seasonal climate prediction. *arXiv preprint arXiv:2402.00712*, 2024.
- Duc Minh Nguyen, Minh Chau Vu, Tuan Anh Nguyen, Tri Huynh, Nguyen Tri Nguyen, and Truong Son Hy. Neural multigrid memory for computational fluid dynamics. *arXiv preprint arXiv:2306.12545*, 2023.
- Michael L Oristaglio and Gerald W Hohmann. Diffusion of electromagnetic fields into a two dimensional earth: A finite-difference approach. *Geophysics*, 49(7):870–894, 1984.
- Maziar Raissi, Paris Perdikaris, and George E Karniadakis. Physics-informed neural networks: A deep learning framework for solving forward and inverse problems involving nonlinear partial differential equations. *Journal of Computational physics*, 378:686–707, 2019.
- Pu Ren, N Benjamin Erichson, Shashank Subramanian, Omer San, Zarija Lukic, and Michael W
 Mahoney. Superbench: A super-resolution benchmark dataset for scientific machine learning.
 arXiv preprint arXiv:2306.14070, 2023.
- Winfried Ripken, Lisa Coiffard, Felix Pieper, and Sebastian Dziadzio. Multiscale neural operators for solving time-independent pdes. *arXiv preprint arXiv:2311.05964*, 2023.
- Olaf Ronneberger, Philipp Fischer, and Thomas Brox. U-net: Convolutional networks for biomed ical image segmentation. In *Medical image computing and computer-assisted intervention– MICCAI 2015: 18th international conference, Munich, Germany, October 5-9, 2015, proceed- ings, part III 18*, pp. 234–241. Springer, 2015.
- 645 Yousef Saad. *Iterative methods for sparse linear systems*. SIAM, 2003.
- 647 Patrick Schnell and Nils Thuerey. Stabilizing backpropagation through time to learn complex physics. *arXiv preprint arXiv:2405.02041*, 2024.

648 649 650	Kimberly Stachenfeld, Drummond B Fielding, Dmitrii Kochkov, Miles Cranmer, Tobias Pfaff, Jonathan Godwin, Can Cui, Shirley Ho, Peter Battaglia, and Alvaro Sanchez-Gonzalez. Learned coarse models for efficient turbulence simulation. <i>arXiv preprint arXiv:2112.15275</i> , 2021.
651 652 653 654	Makoto Takamoto, Timothy Praditia, Raphael Leiteritz, Daniel MacKinlay, Francesco Alesiani, Dirk Pflüger, and Mathias Niepert. Pdebench: An extensive benchmark for scientific machine learning. <i>Advances in Neural Information Processing Systems</i> , 35:1596–1611, 2022.
655 656 657 658	Ronak Tali, Ali Rabeh, Cheng-Hau Yang, Mehdi Shadkhah, Samundra Karki, Abhisek Upadhyaya, Suriya Dhakshinamoorthy, Marjan Saadati, Soumik Sarkar, Adarsh Krishnamurthy, et al. Flowbench: A large scale benchmark for flow simulation over complex geometries. <i>arXiv preprint arXiv:2409.18032</i> , 2024.
659 660 661 662	Artur Toshev, Gianluca Galletti, Fabian Fritz, Stefan Adami, and Nikolaus Adams. Lagrangebench: A lagrangian fluid mechanics benchmarking suite. <i>Advances in Neural Information Processing</i> <i>Systems</i> , 36, 2024.
663 664	Alasdair Tran, Alexander Mathews, Lexing Xie, and Cheng Soon Ong. Factorized fourier neural operators. <i>arXiv preprint arXiv:2111.13802</i> , 2021.
665 666	Lloyd N Trefethen. Spectral methods in MATLAB. SIAM, 2000.
667 668	Rui Wang, Robin Walters, and Rose Yu. Incorporating symmetry into deep dynamics models for improved generalization. <i>arXiv preprint arXiv:2002.03061</i> , 2020.
669 670 671	Fisher Yu, Vladlen Koltun, and Thomas Funkhouser. Dilated residual networks. In <i>Proceedings of the IEEE conference on computer vision and pattern recognition</i> , pp. 472–480, 2017.
672 673 674 675	Sungduk Yu, Walter Hannah, Liran Peng, Jerry Lin, Mohamed Aziz Bhouri, Ritwik Gupta, Björn Lütjens, Justus C Will, Gunnar Behrens, Julius Busecke, et al. Climsim: A large multi-scale dataset for hybrid physics-ml climate emulation. <i>Advances in Neural Information Processing Systems</i> , 36, 2024.
676 677	
678	
679	
680	
681	
682	
683	
684	
685	
686	
687	
688	
689	
690	
691	
692	
693	
694	

702 A APPENDIX

706

708

709 710

711 712

713

714 715

716

717 718 719

723 724 725

729

730

731 732

733

734

704 A.1 ARCHITECTURES

In this section, we discuss the architectures used in more detail and provide information on the training procedures and hyperparameters used. The list of used models is:

1. F-FNO – Factorized Fourier Neural Operator (F-FNO) from (Tran et al., 2021).

- 2. fSNO Spectral Neural Operator (SNO). The construction mirrors FNO, but instead of FFT, a transformation based on Gauss quadratures is used (Fanaskov & Oseledets, 2023).
- 3. DilResNet Dilated Residual Network from (Yu et al., 2017), (Stachenfeld et al., 2021).
- 4. U-Net classical computer vision architecture introduced in (Ronneberger et al., 2015).

F-FNO Unlike the original (Li et al., 2020), the authors of (Tran et al., 2021) proposed to changing the operator layer to:

$$z^{\ell+1} = z^{\ell} + \sigma \Big[W_2^{(\ell)} \sigma \Big(W_1^{(\ell)} \mathcal{K}^{(\ell)}(z^{(\ell)}) + b_1^{(\ell)} \Big) + b_2^{(\ell)} \Big],$$

where σ is an activation function, W_1 and W_2 are weight matrices in the physical space, b_1 and b_2 are bias vectors and

$$\mathcal{K}^{(\ell)}(z^{(\ell)}) = \sum_{d \in D} \Big[\mathrm{IFFT} \big(R_d^{(\ell)} \cdot \mathrm{FFT}_d(z^\ell) \big) \Big],$$

where R_d is a Fourier domain weight matrix, FFT and IFFT are Fast Fourier and inverse Fast Fourier transforms.

F-FNO has an encoder-processor-decoder architecture. We used the following parameters: 4 Fourier layers in the processor, 12 modes and GeLU as the activation function. We used 48 features in the processor.

SNO We utilized spectral neural operators (SNO) (Fanaskov & Oseledets, 2023) with linear integral kernels:

735 736 737

 $u \leftarrow \int dx A_{ij} p_j(x) \left(p_i, u \right) ,$

⁷³⁸ where $p_i(x)$ are orthogonal or trigonometric polynomials.

These linear integral kernels are an extension of the integral kernels used in the FNO (Li et al., 2020). More specifically, starting from the input function u^n , we produce the output function u^{n+1} , which is later transformed by nonlinear activation. The transformation depends on the set of polynomials p_j that form a suitable basis for the problem at hand (e.g. trigonometric polynomials, Chebyshev polynomials, etc.). These polynomials are chosen beforehand and do not change during training. The transformation is naturally divided into three parts: analysis, processing, synthesis.

At the analysis stage, we find a discrete representation of the input function by projecting it onto a set of polynomials. To do this, we compute scalar products:

 $\alpha_j = (p_j, u^n) = \int dx p_j(x) u^n(x) w(x) ,$

where w(x) is a non-negative weight function given by the polynomial used.

At the processing stage, we process the obtained coefficients with a linear layer:

754 755 $\alpha'_i = \sum_j A_{ij} \alpha_j .$ Finally, at the synthesis stage, we recover the continuous function as the sum of the processed coefficients:

 $u^{n+1} = \sum_{j} p_j \alpha_j^{`}.$

We use SNO in Fourier basis (see (Fanaskov & Oseledets, 2023)) with encoder-processor-decoder architecture. The number of SNO layers is 4 and the number of $p_j(x)$ is 20. We use GeLU as activation function.

DilResNet The conventional dilated residual network was first proposed in (Stachenfeld et al., 2021). In this study, the DilResNet architecture is configured with four blocks, each consisting of a sequence of convolutions with steps of [1, 2, 4, 8, 4, 2, 1] and a kernel size of 3. Skip connections are also applied after each block and the GeLU activation function is used.

U-Net We adopt the traditional U-Net architecture proposed in (Ronneberger et al., 2015). This
U-Net configuration is characterised by a series of levels, where each level has approximately half
the resolution of the previous one, and the number of features is doubled. At each level, we apply
a sequence of three convolutions, followed by max pooling, and then a transposed convolution for
upsampling. After upsampling, three more convolutions are applied at each level. The U-Net used
in this study consists of four layers and incorporates the GeLU activation function.