

000 001 002 003 004 005 DIFFWIND: PHYSICS-INFORMED DIFFERENTIABLE 006 MODELING OF WIND-DRIVEN OBJECT DYNAMICS 007 008 009

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ABSTRACT

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032 Modeling wind-driven object dynamics from video observations is highly challeng-
033 ing due to the invisibility and spatio-temporal variability of wind, as well as the
034 complex deformations of objects. We present DiffWind, a physics-informed dif-
035 ferentiable framework that unifies wind-object interaction modeling, video-based
036 reconstruction, and forward simulation. Specifically, we represent wind as a grid-
037 based physical field and objects as particle systems derived from 3D Gaussian
038 Splatting, with their interaction modeled by the Material Point Method (MPM). To
039 recover wind-driven object dynamics, we introduce a reconstruction framework that
040 jointly optimizes the spatio-temporal wind force field and object motion through
041 differentiable rendering and simulation. To ensure physical validity, we incorporate
042 the Lattice Boltzmann Method (LBM) as a physics-informed constraint, enforcing
043 compliance with fluid dynamics laws. Beyond reconstruction, our method natu-
044 rally supports forward simulation under novel wind conditions and enable new
045 applications such as wind retargeting. We further introduce *WD-Objects*, a dataset
046 of synthetic and real-world wind-driven scenes. Extensive experiments demon-
047 strate that our method significantly outperforms prior dynamic scene modeling
048 approaches in both reconstruction accuracy and simulation fidelity, opening a new
049 avenue for video-based wind-object interaction modeling.
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1 INTRODUCTION

052 The motion of objects under wind, such as leaves swaying, flags fluttering, or fabrics billowing,
053 is a visually distinctive yet physically complex phenomenon, arising from the interplay between
054 external fluid forces and internal material properties. From visual observations alone, humans
055 can intuit the presence of invisible wind and anticipate how objects would respond under similar
056 conditions. Enabling computational systems to replicate this ability, which involves recovering
057 both the underlying wind field and the resulting object dynamics from visual input, would have
058 broad impact on applications in augmented and virtual reality, visual effects, scientific analysis, and
059 simulation-based editing.
060

061 However, this task presents significant challenges: wind is invisible, dynamic, and spatially non-
062 uniform, while object deformations depend on unknown physical parameters such as mass, elasticity,
063 and geometry. Existing methods face significant limitations: Dynamic neural representations, such as
064 Deformable Nerf Du et al. (2021); Chu et al. (2022) and 3D Gaussian Splatting Yang et al. (2024b);
065 Wu et al. (2024), model object appearance and motion over time but capture only visible dynamics,
066 ignoring underlying physical causes like wind fields. Differentiable physics simulators Zhong et al.
067 (2024); Li et al. (2023b); Zhang et al. (2024) can optimize motion parameters, but are restricted
068 to simple, predefined motion patterns such as constant-force projectile motion and cannot handle
069 complex fluid-object interactions. Video-based wind inference methods Zhang et al. (2022); Runia
070 et al. (2020) estimate coarse wind speed or target specific scenarios such as cloth deformation, lacking
071 generality and physical consistency. These limitations motivate the following question: *Can we*
072 *jointly recover visible object dynamics and invisible wind fields from videos input, while ensuring*
073 *physical consistency and generalization to arbitrary wind conditions?*

074 To address this challenge, we propose *DiffWind*, a physics-informed differentiable framework for
075 wind-object interaction. In DiffWind, the invisible wind is modeled as a grid-based physical field,
076 while objects are modeled as deformable particle systems derived from 3D Gaussian Splatting, with
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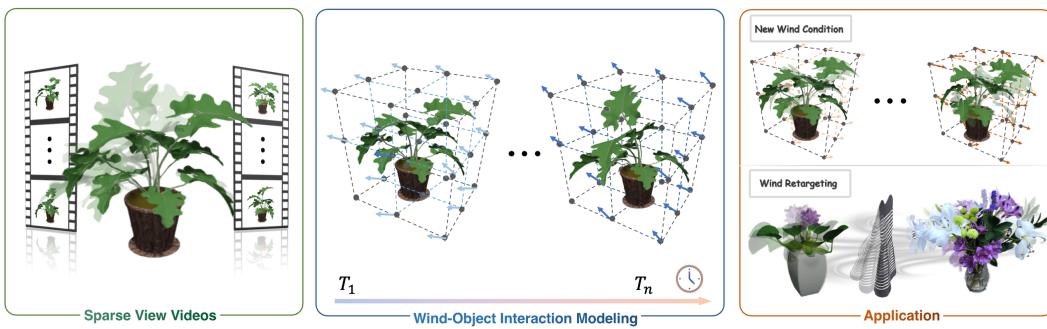


Figure 1: DiffWind is a physics-informed differentiable wind-object interaction framework that models wind and object dynamics separately. This design enables the reconstruction of wind and object motion from sparse-view videos, physically consistent simulation under new wind conditions, and retargeting to novel objects.

their interaction modeled by the Material Point Method (MPM). This design leverages physical intuition: fluids are naturally defined on Eulerian grids, where quantities such as velocity and pressure evolve over space and time, whereas solid or semi-solid objects undergo localized deformation, making Lagrangian particles a more effective representation. By coupling these domains, our formulation supports accurate reconstruction, physically plausible simulation, and novel applications such as wind retargeting. (Fig. 1)

A key strength of DiffWind is its differentiable formulation, which enables joint reconstruction of the spatio-temporal wind force field and dynamic object motion from sparse-view RGB videos. Reconstruction is achieved by minimizing the image-space reconstruction error via gradient-based optimization. However, photometric loss alone cannot guarantee physical validity, as wind forces follow complex fluid interactions that often entangle physical properties with scene-specific dynamics, limiting transferability. To address this, we introduce a *physics-informed optimization loss* that enforces compliance with fluid dynamics. Specifically, we use the Lattice Boltzmann Method (LBM) Li et al. (2023a) to provide directional guidance for the wind force field at each time step, constraining its reconstruction to yield physically plausible results.

Our contributions can be summarized as follows:

- We propose a novel differentiable modeling framework DiffWind for wind-object interaction, where the wind is represented as a grid-based physical field and the object is modeled with particle-based deformable geometry. This particle-grid coupling enables physically plausible, 3D-consistent simulations of wind-induced object motion.
- Based on the proposed interaction model, we develop a differentiable inverse reconstruction framework that simultaneously recovers the dynamic motion of visible objects and the invisible wind force field from sparse-view videos.
- We employ the Lattice Boltzmann Method (LBM) as a physics-informed constraint to ensure the wind force field adheres to fluid dynamics laws, and demonstrate applications in forward simulation under novel wind conditions and wind retargeting.
- We construct WD-Objects, a dataset covering both synthetic and real-world wind-driven object scenes. Extensive experiments show that DiffWind substantially outperforms existing methods in reconstruction accuracy, simulation fidelity, and generalization ability, advancing the frontier of video-based wind-object modeling.

2 RELATED WORK

4D Dynamic Scene Reconstruction. Dynamic scene reconstruction and free-viewpoint rendering have been widely explored by extending Neural Radiance Field (NeRF) Mildenhall et al. (2020) and 3D Gaussian Splatting (3DGS) Kerbl et al. (2023). NeRF-based approaches model dynamics either by adding temporal inputs Du et al. (2021); Chu et al. (2022) or by learning canonical spaces with deformation fields Pumarola et al. (2021); Park et al. (2021). 3DGS offers fast reconstruction and high-quality rendering, and has been extended to dynamic scenes via Gaussian tracking Luiten et al. (2024), canonical-space deformation Yang et al. (2024b); Wu et al. (2024), or 4D Gaussian

108 primitives Duan et al. (2024); Wang et al. (2025). But they mainly capture visible object dynamics
 109 and struggle to model complex, invisible effects such as wind fields, limiting their applicability to
 110 wind-driven motion modeling.
 111

112 **Windy Object Modeling.** Windy object modeling involves: (1) simulating wind-induced dynamics,
 113 and (2) reconstructing physical parameters from observed motion. Existing simulations use grid-
 114 based or particle-based fluid methods to animate specific objects (e.g., grass, trees, cloth) Lo et al.
 115 (2016); Pirk et al. (2014); Wilson et al. (2014), but are often category-specific. We instead employ the
 116 Material Point Method (MPM) Hu et al. (2018) to model complex deformations of various categories
 117 of objects under wind fields. Reconstruction methods estimate wind speed or material properties
 118 from monocular videos Runia et al. (2020); Cardona et al. (2019); Bouman et al. (2013); Yang et al.
 119 (2017), but cannot recover full 4D dynamics. Our method jointly reconstructs temporally varying
 120 wind fields and dynamic scenes via differentiable optimization.
 121

122 **Differentiable Physics-Based System Identification.** System identification is a methodology for
 123 building mathematical models of dynamic systems using measurements of the input and output
 124 signals of the system. Recent advances in differentiable physical simulation, enabled by frameworks
 125 like Taichi Hu et al. (2019; 2020) and Warp Macklin (2022), have facilitated such tasks. Prior works
 126 optimize material or geometry parameters for cloth Li et al. (2022), solids Jin et al. (2024), and
 127 fluids Li et al. (2024b) using differentiable solvers. Other studies combine differentiable simulation
 128 with NeRF Mildenhall et al. (2020) or 3DGS Kerbl et al. (2023) to jointly recover object geometry
 129 and physical properties from multi-view observations Li et al. (2023b); Zhong et al. (2024); Cai
 130 et al. (2024); Zhang et al. (2024); Cao et al. (2024); Liu et al. (2025), but are restricted to simple,
 131 predefined motion settings (e.g., gravity, dragging) and optimize initial velocity at the first frame.
 132 Moreover, these approaches focus solely on object motion and do not model the influence of the
 133 external environment such as wind fields on the dynamics. In contrast, our method reconstructs
 134 the wind force field through a differentiable optimization process, while enabling the modeling of
 135 complex and dynamic wind–object interactions.
 136

3 METHOD

137 **Problem Specification:** Given the sparse-view observed videos as input $\mathbf{V} = \{\{I_0^i, \dots, I_T^i\}|i = 1, \dots, N_v\}$ of the wind-object interaction, our goal is to reconstruct the dynamic process in 3D. Our
 138 key idea is to optimize the wind force field to match the observed sequences through differentiable
 139 simulation and differentiable rendering.
 140

141 As shown in Fig. 2, we model the wind as a grid field, where each node can contain various physical
 142 quantities. The object is represented as a set of particles, each carrying attributes such as appearance,
 143 material, and motion. We employ a differentiable Material Point Method (MPM) to model the
 144 interaction between the wind and the object. The *force* property of the wind grid field is applied
 145 to the background grid of MPM, which drives the motion of object particles and results in object
 146 deformation. More details about the interaction modeling are provided in Sec. 3.2. Then, given
 147 observed sequences of object motion under wind influence, we perform differentiable optimization
 148 of the *force* property in the wind grid field that enables the reconstruction of the dynamic process
 149 (Sec. 3.3). Furthermore, we propose a novel physics-informed optimization method that leverages
 150 LBM to enforce the force field’s compliance with the fundamental laws of fluid dynamics (Sec. 3.4).
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3.1 PRELIMINARY: PHYSICAL MODELS

152 **Fluid Mechanics.** Our wind field dynamics are governed by incompressible Navier-Stokes equa-
 153 tions (NSE). We build the fluid simulator by leveraging the Lattice Boltzmann Method (LBM). With
 154 t discretized into fixed intervals, \mathbf{x}_w discretized on a regular grid and \mathbf{u}_w discretized into fixed
 155 unit directions at grid points. Under this discretization, one can solve the NSE by computing the
 156 corresponding lattice Boltzmann equation (LBE) SHAN et al. (2006):
 157

$$\mathbf{f}_i(\mathbf{x}_w + \mathbf{c}_i, t + 1) - \mathbf{f}_i(\mathbf{x}_w, t) = \Omega_i + \mathbf{F}_{\mathbf{o}i}, \quad (1)$$

158 where \mathbf{c}_i is the discrete lattice velocity in the i -th direction, Ω_i is the Bhatnagar-Gross-Krook(BGK)
 159 collision model term SHAN et al. (2006) and $\mathbf{F}_{\mathbf{o}i}$ is the external force. The macroscopic physical
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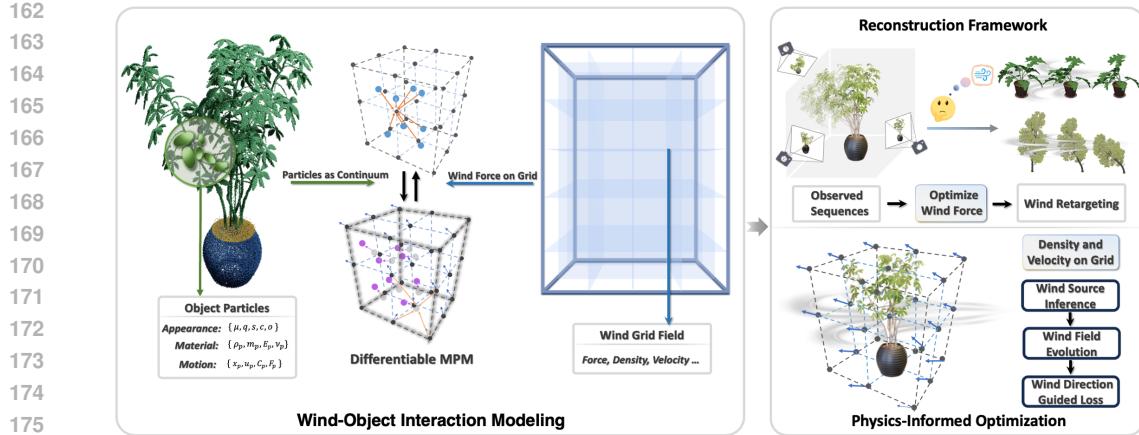


Figure 2: Overview of DiffWind. We propose a novel wind-object interaction modeling approach, where the wind is represented as a grid field and the object is modeled as a set of particles. Based on this modeling approach, we introduce a reconstruction framework for wind-object interaction by optimizing the wind force field. In addition, we employ the Lattice Boltzmann Method (LBM) to generate wind force field direction guidance to enforce compliance with fluid dynamics laws.

quantities of the wind field, such as density ρ_w , linear momentum $\rho_w \mathbf{u}_w$, and the stress tensor \mathbf{S}_w can be calculated as:

$$\rho_w = \sum_{i=0}^{q-1} \mathbf{f}_i, \quad \rho_w \mathbf{u}_w = \sum_{i=0}^{q-1} \mathbf{c}_i \mathbf{f}_i + \frac{1}{2} \mathbf{F}_{\mathbf{o}i}, \quad \rho \mathbf{S}_{w,\alpha\beta} = \sum_{i=0}^{q-1} \left(\mathbf{c}_i^2 - \frac{1}{3} \delta_{\alpha\beta} \right) \mathbf{f}_i, \quad (2)$$

where $\alpha, \beta \in \{x, y, z\}$ refer to the coordinates of the stress tensor \mathbf{S}_w and δ is Kronecker delta. In particular, we use HOME-LBM Li et al. (2023a) as our simulator, which solves the LBE through High-Order Moment Encoding with Hermite polynomial expansion SHAN et al. (2006) to achieve more efficient and accurate fluid simulation.

Continuum Mechanics. Our object deformations are governed by continuum mechanics which describes motions by a time-dependent deformation map $\mathbf{x}_o = \phi(\mathbf{X}_o, t)$ that relates rest material position and deformed material position and the deformation gradient $\mathbf{F}(\mathbf{X}_o, t) = \nabla_{\mathbf{X}_o} \phi(\mathbf{X}_o, t)$. The evolution of ϕ is governed by:

$$\frac{\partial \rho_o}{\partial t} + \mathbf{u}_o \cdot \nabla \rho_o + \rho_o \nabla \cdot \mathbf{u}_o = 0, \quad \rho_o \left(\frac{\partial \mathbf{u}_o}{\partial t} + \mathbf{u}_o \cdot \nabla \mathbf{u}_o \right) = \nabla \cdot \boldsymbol{\sigma} + \rho_o \mathbf{F}_w, \quad (3)$$

where ρ_o is the mass density, \mathbf{u}_o is the local material velocity of the object, $\boldsymbol{\sigma} = \frac{1}{J} \mathbf{P}(\mathbf{F}) \mathbf{F}^T$ is the Cauchy stress. These two equations respectively describe the conservation of mass and momentum. We use the differentiable Material Point Method Hu et al. (2020) to evolve Eq. (3).

3.2 WIND-OBJECT INTERACTION MODELING

We represent the object with 3D Gaussians, a particle-based approach suited for high-quality reconstruction and rendering. The wind, being invisible and difficult to represent with discrete particles, is modeled as a continuous grid-based field, where each grid node stores physical quantities such as density, velocity, and force, and its evolution is simulated using the Lattice Boltzmann Method (LBM). To model the interaction between the two representations, we adopt the Material Point Method (MPM), which exchanges information between particle-based objects and the grid-based wind: objects are treated as moving Lagrangian particles, while wind acts as a force field on an Eulerian grid.

Object Modeling. To integrate 3D Gaussians into physical simulations, the representation should minimize rendering artifacts under deformation, accurately capture surfaces for boundary conditions, preserve volume integrity, and support region-wise physical properties. Vanilla 3DGS, however, may produce needle-like artifacts and floaters under large deformations, lacks internal structure, and creates holes in weakly textured regions. We mitigate these issues by adding regularization terms during optimization and densifying invisible points in geometry-deficient areas.

216 In particular, we define the following loss to optimize the 3D Gaussians:
 217

$$\mathcal{L}_{\text{static}} = \mathcal{L}_{\text{color}} + \lambda_a \mathcal{L}_{\text{aniso}} + \lambda_o \mathcal{L}_o, \quad (4)$$

219 where $\mathcal{L}_{\text{color}}$ is the color loss between render and training image in vanilla 3DGS; $\mathcal{L}_{\text{aniso}}$ is designed
 220 to constrain the anisotropy of the Gaussian kernels:
 221

$$\mathcal{L}_{\text{aniso}} = \frac{1}{|P|} \sum_{i \in P} \max \left(\frac{\max (\{s_i^x, s_i^y, s_i^z\})}{\text{median} (\{s_i^x, s_i^y, s_i^z\})}, \epsilon \right) - \epsilon, \quad (5)$$

224 where $\epsilon = 2.0$ is a hyperparameter that controls the level of anisotropy; and \mathcal{L}_o is designed to
 225 encourage the 3D Gaussian points to stay close to the object surface by encouraging each Gaussian's
 226 opacity to be either near zero or near one:
 227

$$\mathcal{L}_o = \exp (-\beta_o \cdot (o_i - 0.5)^2), \quad (6)$$

229 where $\beta_o = 20$ is a hyperparameter. Additionally, we prune Gaussian points whose opacity falls
 230 below a threshold $\tau_o = 0.05$.

231 Based on the optimized Gaussian kernels, we employ the octree-based voxel filling algorithm from
 232 Kaolin Fuji Tsang et al. (2024) to densify additional invisible points for improved geometric support.
 233 Since 3DGS directly models the entire scene, extracting objects for simulation becomes necessary.
 234 Moreover, physical properties related to object deformation cannot be directly obtained from static
 235 image observations. Therefore, we train a 3D-consistent feature field by contrastive learning to
 236 extract 2D segmentation priors for 3D regions segmentation in the scene, and employ a multimodal
 237 large language model (MLLM) to construct a physical agent that infers their corresponding physical
 238 properties, including material name, density, Poisson's ratio ν_p , and Young's modulus E . For more
 239 details about 3D regions segmentation and the physical agent, please refer to the Appendix.

240 **Interaction Modeling.** We employ Material Point Method (MPM) to model the interaction between
 241 the object and wind, by binding each Gaussian kernel and densified point to a corresponding particle
 242 in the MPM simulation. During the P2G and G2P process in MPM, wind is applied as a force field
 243 to the MPM grid to simulate the deformation of objects. Similar to PhysGaussian Xie et al. (2024),
 244 we update the covariance matrix of 3DGS as $\Sigma_p(t) = \mathbf{F}_p(t) \Sigma_p \mathbf{F}_p(t)^T$, $\mathbf{F}_p(t)$ is the deformation
 245 gradient of each particle p at time step t .

246 3.3 RECONSTRUCTION FRAMEWORK FOR WIND-OBJECT INTERACTION

248 We formulate the simulation process of wind-object interaction for a single step as:
 249

$$(x^{t+1}, v^{t+1}, F^{t+1}, C^{t+1}) = \mathcal{S}(x^t, v^t, F^t, C^t, \theta, \mathbf{F}_w, \Delta t), \quad (7)$$

251 where (x^t, v^t, F^t, C^t) denote the position, velocity, deformation gradient, and affine momentum of
 252 all object particles at time t , respectively. Note that the particles include both Gaussian kernels and
 253 densified invisible points. θ denotes the collection of the physical properties of all particles, which
 254 are obtained during the object modeling step, and \mathbf{F}_w denotes the wind force field.

255 We can then render the object at each time and each viewpoint by 3DGS:
 256

$$\hat{I}_t = \mathcal{R}_{\text{render}}(x_{gs}^t, \sigma, F_{gs}^t, \Sigma, c, W), \quad (8)$$

259 where x_{gs}^t and F_{gs}^t denote the subset of x^t and F^t corresponding to the Gaussian kernels.

260 Owing to the differentiable nature of Eq. (7) and Eq. (8), we can optimize the invisible wind force
 261 field from the visible motion of the object by minimizing the photometric loss between the rendered
 262 image and input image as:
 263

$$\mathcal{L}_{\text{render}} = \sum_{i=1}^{N_v} (1 - \lambda) \mathcal{L}_1^i + \lambda \mathcal{L}_{D-SSIM}^i, \quad (9)$$

266 where N_v denotes the number of viewpoints, \mathcal{L}_1^i and \mathcal{L}_{D-SSIM}^i represent the L_1 and structural
 267 similarity index measure losses at viewpoint i , respectively. We set $\lambda = 0.1$ in our experiments.
 268

269 Although θ also receives gradients during backpropagation, jointly optimizing both the wind force
 270 field \mathbf{F}_w and the material parameters leads to an ill-posed problem, because different combinations

270 of material stiffness and external forces can produce nearly identical motions. For instance, a small
 271 Young’s modulus paired with a small force can generate the same motion as a large Young’s modulus
 272 with a large force. Therefore, we use MLLM-based reasoning to obtain reasonable material priors
 273 as initialization (see Sec. 3.2), and treat only \mathbf{F}_w as the optimization target, initializing it from zero
 274 for each time-step optimization. We use the results from object modeling as the initial state and
 275 optimize the wind force field between two consecutive states. The optimization proceeds sequentially,
 276 with each time step starting only after the previous one has been completed. The optimized object
 277 state from the previous step serves as the initial condition for the next which mitigates gradient
 278 instability that can occur when performing multi-step simulation followed by a single backward pass.
 279 By optimizing the wind force field at each time step, we simultaneously achieve the reconstruction of
 280 wind-object interaction dynamics.

281 **Wind Retargeting.** By accurately reconstruction of wind-object interaction, the proposed frame-
 282 work enables a novel task: wind retargeting. This involves applying the estimated wind force fields to
 283 other objects, as illustrated in Fig. 1, which is infeasible for the existing object dynamics modeling
 284 methods like Deformable-GS Yang et al. (2024b), which only focus on dynamic scene reconstruction.

285 3.4 PHYSICS-INFORMED OPTIMIZATION OF WIND FORCE FIELD

286 Although we can optimize the wind force field in each timestep using visual observations, it does not
 287 guarantee adherence to physical laws as wind forces are fluid and exhibit complex interactions with
 288 the target objects. Therefore, we further exploit a physics-informed optimization loss to guide the
 289 force field optimization process. Specifically, we use the Lattice Boltzmann Method (LBM) to solve
 290 Eq. (1), thereby generating the guiding direction of the force field.

291 **Wind Source Inference.** We start from the posed RGB video and estimate temporally-consistent
 292 monocular depth maps using Video Depth Anything Chen et al. (2025). The RGB-depth pairs are
 293 provided to a multimodal large language model (MLLM), which infers the wind source direction.
 294 Then, the inferred direction serves as an inlet velocity boundary condition for the LBM simulator.

295 **Wind Field Evolution.** We use HOME-LBM Li et al. (2023a) to simulate wind field evolution,
 296 where the positions of particles serve as the boundary condition for LBM, thereby influencing the
 297 wind field dynamics. In particular, we first distinguish solid and fluid nodes in the LBM lattice based
 298 on the positions of particles at the current state by converting object particles into occupied voxel
 299 grids:

$$300 \text{Solid}(\mathbf{x}_w) = \begin{cases} 1, & \text{if } \exists \mathbf{x}_p \text{ s.t. } \forall d \in \{x, y, z\}, 0 \leq \mathbf{x}_{p,d} - \mathbf{x}_w < \Delta x \\ 0, & \text{otherwise} \end{cases} \quad (10)$$

301 where \mathbf{x}_w is the position of wind grid node, \mathbf{x}_p is the position of particle, and Δx is the grid spacing.
 302 This mapping enables the LBM simulator to capture the effect of the object’s instantaneous position
 303 on the wind field. At each simulation step, macroscopic wind field quantities are computed from the
 304 distribution functions using Eq. (2). Additional details of the algorithm are provided in the Appendix.

305 **Physics-Informed Loss.** The wind velocity field from the LBM simulator $\mathbf{D}_{\text{guide}}$ is used to guide
 306 the direction of the reconstructed wind force field $\mathbf{D}_{\text{recon}}$:

$$307 \mathcal{L}_{\text{phys}} = \|\mathbf{D}_{\text{recon}} - (\mathbf{D}_{\text{recon}} \cdot \mathbf{D}_{\text{guide}}) \mathbf{D}_{\text{guide}}\|^2, \quad (11)$$

308 where $\mathcal{L}_{\text{phys}}$ is the physics-informed optimization loss that measures the directional difference between
 309 the reconstructed wind force field and the one obtained from the LBM simulator. Here, the wind
 310 force is governed by the aerodynamic drag equation $\mathbf{F}_w = \frac{1}{2} \rho_w C_D |\mathbf{v}|^2 \cdot \frac{\mathbf{v}}{|\mathbf{v}|}$, where C_D is the drag
 311 coefficient and \mathbf{v} is the wind velocity, this formulation implies that an accurate estimation of the wind
 312 velocity field direction can effectively guide the reconstruction of the wind force field. Through both
 313 physics-informed constraints and visual supervision, we ensure the accuracy and continuity of the
 314 reconstructed wind force field.

315 4 EXPERIMENTS

316 In this section, we first introduce the implementation details of our method (Sec. 4.1), followed by
 317 evaluations of the capability in reconstructing wind-driven object dynamics (Sec. 4.2). We then

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Figure 3: Our reconstruction dataset from one camera view, with PSNR values of our dynamic reconstruction shown above. Please use **Adobe Reader/KDE Okular** to see **animations**.

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Figure 4: Qualitative results of novel view synthesis for a selected synthetic wind-object interaction scene. We compare our method with Deformable-GS Yang et al. (2024b), Efficient-GS Katsumata et al. (2024), 4D-GS Wu et al. (2024) and GaussianPrediction Zhao et al. (2024). Please use **Adobe Reader/KDE Okular** to see **animations**.

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conduct ablation studies (Sec. 4.3) to validate key design choices. Finally, we demonstrate the application of our method to forward simulation of wind-driven object dynamics (Sec. 4.4).

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4.1 IMPLEMENTATION DETAILS

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For real-world captured data, we first use Detector-Free SFM He et al. (2024) to compute the pose of each viewpoint. Subsequently, we use EntitySeg Lu et al. (2023) to generate instance-level 2D segmentation masks and Grounded-SAM Ren et al. (2024) to generate part-level 2D segmentation masks. Afterward, we train 3D Gaussians for 60k iterations. For 3D regions segmentation, we train 20k iterations to build 3D consistent feature fields, and then use GPT5 Hurst et al. (2024) for physical property reasoning. In our experiments, we divide the simulation space into a 128^3 grid, and for real-world scenes, only the segmented foreground objects are used as input to the simulator. All components of our simulator are implemented using Taichi Hu et al. (2019; 2020), and all our experiments are evaluated on a single NVIDIA GeForce RTX 4090 GPU.

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4.2 EVALUATION ON RECONSTRUCTION

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In this section, we first introduce our datasets, then compare our method with SOTA dynamic scene reconstruction approaches for novel view synthesis, and finally show our wind retargeting results.

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Datasets. As there are no publicly available datasets for modeling wind-driven object dynamics, we introduce both synthetic and real-world datasets, named *WD-Objects*, in this paper for evaluation. For the synthetic dataset, we first reconstruct seven object-level 3D Gaussians from various sources, including NeRF Mildenhall et al. (2020), PhysDreamer Zhang et al. (2024), ShapeNet Chang et al. (2015); Ma et al. (2024), and LoopGaussian Li et al. (2024a). We then simulate dynamic scenes of the objects moving with the wind, further details are provided in Sec. 4.4. Next, we select eight orthogonal views to render reference videos, where four viewpoints are used for training and the remaining four for evaluating the quality of novel view synthesis. The evaluation is conducted over 50 simulation time steps. We show the object states from one camera view in Fig. 3. For a more complete visualization of the motion process, please refer to our video supplementary. For the real-world dataset, we record 360-degree surround videos of real-world static scenes by GoPro cameras, each scene includes an object and a background. The objects include a pothos plant, a

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 379 Table 1: Quantitative comparison on the synthetic dataset for the novel view synthesis. Best results
 380 are highlighted as **first**, **second**, **third**.

Method	Dress			Flag			Pants			Sweater		
	PSNR(↑)	SSIM(↑)	LPIPS(↓)	PSNR(↑)	SSIM(↑)	LPIPS(↓)	PSNR(↑)	SSIM(↑)	LPIPS(↓)	PSNR(↑)	SSIM(↑)	LPIPS(↓)
Deformable-GS	27.32	.9715	.0253	40.86	.9939	.0057	42.36	.9924	.0149	33.10	.9557	.0594
Efficient-GS	22.79	.9470	.0481	26.72	.9483	.0499	34.73	.9782	.0385	26.06	.9241	.0920
4D-GS	27.64	.9729	.0239	36.96	.9868	.0147	41.33	.9911	.0184	30.14	.9556	.0702
GaussianPrediction	29.00	.9773	.0201	39.89	.9930	.0067	42.11	.9917	.0168	28.64	.9465	.0734
Ours	41.76	.9986	.0007	41.86	.9974	.0009	58.38	.9998	.0002	41.72	.9951	.0021
Method	Ficus			ShapeNetPlant			Alocasia			Average		
	PSNR(↑)	SSIM(↑)	LPIPS(↓)	PSNR(↑)	SSIM(↑)	LPIPS(↓)	PSNR(↑)	SSIM(↑)	LPIPS(↓)	PSNR(↑)	SSIM(↑)	LPIPS(↓)
Deformable-GS	41.72	.9967	.0030	38.58	.9946	.0047	32.79	.9681	.0364	36.68	.9818	.0213
Efficient-GS	22.67	.8951	.0906	24.63	.9467	.0532	28.41	.9463	.0533	26.57	.9408	.0608
4D-GS	28.80	.9508	.0372	32.91	.9865	.0115	28.96	.9491	.0386	32.39	.9704	.0306
GaussianPrediction	38.81	.9920	.0083	37.46	.9935	.0061	34.65	.9738	.0292	35.79	.9811	.0229
Ours	48.82	.9995	.0003	52.54	.9998	.0001	44.94	.9980	.0010	47.15	.9983	.0008

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 392 Table 2: Quantitative comparison on the real-world dataset for the novel view synthesis. Best results
 393 are highlighted as **first**, **second**, **third**.

Method	POTHOS			HAT			POMPON			TULIP		
	PSNR(↑)	SSIM(↑)	LPIPS(↓)									
Deformable-GS	24.95	.9572	.0362	30.55	.9575	.0297	26.06	.9604	.0303	20.49	.9536	.0439
Efficient-GS	22.25	.9368	.0573	29.53	.9562	.0442	24.02	.9477	.0447	19.34	.9432	.0707
4D-GS	24.37	.9516	.0427	30.32	.9556	.0347	26.23	.9614	.0298	20.41	.9527	.0462
GaussianPrediction	24.69	.9532	.0401	30.36	.9564	.0307	25.57	.9568	.0348	19.96	.9511	.0463
Ours	26.18	.9692	.0256	31.37	.9585	.0281	26.14	.9660	.0267	23.53	.9719	.0308

400 beanie hat, a pompon flower, and a tulip. Additionally, we capture synchronized videos of these
 401 plants being affected by a hairdryer.

402 **Baselines.** We compare our approach against the current state-of-the-art dynamic scene reconstruction
 403 methods: Deformable-GS Yang et al. (2024b), Efficient-GS Katsumata et al. (2024), 4D-GS Wu
 404 et al. (2024) and GaussianPrediction Zhao et al. (2024). For a fair comparison, we run their optimiza-
 405 tion process using our reconstructed static 3D Gaussians as initialization, and apply the same
 406 regularization terms as described in Sec. 3.2.

408 **Comparison on Novel View Synthesis.** The quantitative results are shown in Table 1 and Table 2.
 409 We present the PSNR/SSIM/LPIPS (VGG) values for the novel view rendering results to validate
 410 the accuracy of our wind force field reconstruction results. We can see that the proposed method
 411 significantly outperforms existing methods. A set of qualitative results is shown in Fig. 4, with full
 412 results provided in the Appendix, demonstrating that our method produces more realistic novel views
 413 than existing approaches, owing to its superior modeling of wind-driven object dynamics.

4.3 ABLATION STUDIES

416 **Effectiveness of Physics-Informed Loss.** We evaluate
 417 the effectiveness of physics-informed loss during dynamic
 418 reconstruction. The quantitative results in Table 3, show
 419 that our full model, by embedding physical directional
 420 constraints during training, further significantly enhances
 421 rendering quality, resulting in reconstructions that are not
 422 only more precise but also more physically plausible.

423 **Robustness against Physical Property Reasoning.** We
 424 also inspect the robustness of our method to variations in
 425 material parameters. To this end, we conduct experiments
 426 on the synthetic dataset to analyze how changes in Young’s
 427 modulus affect reconstruction performance. We predefine
 428 three categories of Young’s modulus: *soft*, *medium*, and
 429 *hard*, with values of $E = 4 \times 10^4$, 4×10^5 , and 4×10^6 ,
 430 respectively, and also include values inferred using MLLM.
 431 We perform dynamic reconstruction for each case without
 the physics-informed loss, in order to better assess the influence of physical properties on the

Table 3: Ablation of effectiveness of physics-informed loss on the synthetic dataset.

Config.	PSNR	SSIM	LPIPS $\times 10^{-1}$
w/o $\mathcal{L}_{\text{phys}}$	51.31	.9995	.0030
Full Model	53.24	.9997	.0021

Table 4: Ablation of robustness against physical property reasoning on the synthetic dataset.

Config.	PSNR	SSIM	LPIPS $\times 10^{-1}$
Soft	49.78	.9993	.0040
Medium	49.83	.9992	.0044
Hard	48.64	.9988	.0068
MLLM	51.31	.9995	.0030

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Ours

DynamuCrafter

CogVideoX

Figure 5: Wind synthesis results at the same viewpoint across different timesteps. DynamuCrafter Xing et al. (2023) fails to maintain temporal coherence, CogVideoX Yang et al. (2024a) produces unrealistic jittering. In contrast, DiffWind generates realistic time-evolving wind-object interactions. Please use **Adobe Reader/KDE Okular** to see **animations**.

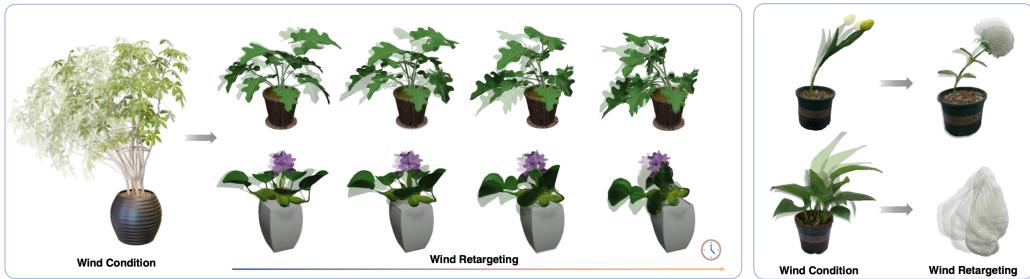


Figure 6: Wind retargeting on synthetic data (left) and real-world data (right).

reconstruction results. Table 4 shows that using different predefined physical parameters still achieves comparable reconstruction performance, while the values inferred by MLLM yield slightly better results. Hence, we adopt MLLM for more convenient and accurate physical property estimation.

Necessity of Invisible Points Densification. We also evaluate the impact of invisible points densification, which provides geometric support for the simulation process. The quantitative results in Table 5 indicate that training with densified points improves the reconstruction results, justifying its necessity for accurate geometric modeling.

Influence of Grid Resolution. To evaluate how discretization resolution affects reconstruction quality and computational efficiency, we compare grid sizes of 64^3 , 128^3 , and 256^3 . As shown in Table 6 the quantitative results show that reconstruction quality remains stable across different grid resolutions, while 128^3 provides the best trade-off between fidelity and runtime. Under our default 128^3 setting, each optimization iteration, including the LBM update, the MPM simulation, the differentiable rendering, and the gradient backpropagation, takes approximately 1.17 seconds on a single NVIDIA RTX 4090 GPU, compared to 1.02 seconds for 64^3 and 2.02 seconds for 256^3 .

4.4 FORWARD SIMULATION OF WIND-DRIVEN OBJECT DYNAMICS

Simulation under Specified Wind Conditions. Benefiting from our wind-object interaction modeling approach, we enable forward simulation of wind-driven object dynamics, where object deformations are simulated by MPM and wind field evolution is simulated by LBM. To evaluate the visual quality of our wind-object simulations, we compare our method against SOTA and publicly available video generation models, including SVD Blattmann et al. (2023), CogVideoX Yang et al. (2024a), VideoCrafter Chen et al. (2024), and DynamuCrafter Xing et al. (2023). The comparison is performed on four test scenes. The multi-view images of the static test scenes, Vase and Garden, are sourced from NeRF Mildenhall et al. (2020) and Feature-Splatting Qiu et al. (2024), respectively. Additionally, we captured two static scenes ourselves, Cloth and Sock.

Table 5: Ablation of necessity of invisible points densification on the real-world dataset.

Config.	PSNR	SSIM	LPIPS $\times 10^{-1}$
w/o Densification	27.31	.9614	.0292
Full Model	27.90	.9646	.0268

Table 6: Ablation of influence of grid resolution on the synthetic dataset.

Config.	PSNR	SSIM	LPIPS $\times 10^{-1}$
64^3	49.08	.9991	.0046
128^3	51.31	.9995	.0030
256^3	51.90	.9997	.0024

486 We conduct a user study with 32 participants, who rate
 487 20 randomly ordered videos generated by different methods.
 488 Both motion realism and the physical realism of
 489 wind–object interactions are evaluated on a 5-point Lik-
 490 ert scale, from 1 (strongly disagree) to 5 (strongly agree).
 491 Results in Table 7 demonstrate the superior visual quality
 492 of our simulations, with comparative examples shown in
 493 Fig. 5. Additional results are provided in the Appendix
 494 and supplementary video.

495 **Wind Retargeting.** By using the reconstructed wind force field as input for forward simulation, our
 496 method enables a new application, i.e., wind retargeting. We present several representative results of
 497 wind retargeting in Fig. 1, Fig. 2 and Fig. 6, with additional examples provided in our supplementary
 498 video. It can be observed that our method effectively decouples the invisible wind force field from
 499 the visible object dynamics, while enabling accurate transfer to previously unseen scenes.

500 5 CONCLUSION

501 In this paper, we present a novel framework for modeling wind-object interactions. The proposed
 502 framework can reconstruct wind-driven object dynamics, simulate in new wind conditions and
 503 perform wind retargeting. To build this framework, we represent wind as a grid-based physical field
 504 and model objects using particle-based deformable geometry. By leveraging differentiable physical
 505 simulation and rendering, our system achieves backward reconstruction of wind-object interactions.
 506 In addition, we employ the LBM as a physics-informed constraint to enforce compliance with fluid
 507 dynamics laws. We also present a new dataset for comprehensive evaluation. The experiments
 508 demonstrate the superiority of the proposed framework in modeling realistic wind-object dynamics.

509 **Limitation and Future Work.** The proposed framework currently focuses on modeling object-level
 510 dynamics under wind conditions, without accounting for interactions between objects. **Robustly**
 511 **handling multi-object collisions is challenging due to discontinuous contact dynamics and complex**
 512 **force interactions, which complicate gradient-based optimization.** In addition, MPM is primarily
 513 used for simulating continuum objects, **however, our framework is simulator-agnostic and can be**
 514 **extended to non-continuum objects by replacing MPM with an appropriate differentiable simulator**
 515 **without changing the overall formulation,** other types of simulators such as the spring–mass method
 516 or the Finite Element Method (FEM) could be explored to model a wider range of object motion
 517 behaviors. We leave these extensions as future work.

518 ETHICS STATEMENT

519 Our study focuses on modeling wind-driven object dynamics using physics-informed differentiable
 520 simulations. All experimental evaluations are conducted using synthetic, publicly available, or
 521 self-captured datasets, curated to avoid sensitive or private content. We assert that this research has
 522 been carried out in accordance with the code of ethics.

523 REPRODUCIBILITY STATEMENT

524 In order to support reproducibility and verification, we present implementation and evaluation details
 525 in our paper, and will release the associated source code upon acceptance.

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Table 7: Human evaluation on visual quality.

Method	Visual Quality				
	Cloth	Garden	Sock	Vase	Avg
SVD	2.59	2.53	2.66	2.19	2.49
CogVideoX	2.31	1.72	2.69	2.63	2.34
VideoCrafter	2.56	2.44	2.38	1.94	2.33
DynamiCrafter	1.97	3.33	2.69	3.44	2.84
Ours	4.47	4.16	4.59	4.13	4.34

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693

APPENDIX

We highly encourage watching the **supplementary video** where our results are best illustrated.

A BACKGROUND

A.1 3D GAUSSIAN REPRESENTATION

3D Gaussian Splatting (3DGS) Kerbl et al. (2023) employs a substantial number of explicit 3D Gaussians to represent a static 3D scene. Each 3D Gaussian G is defined by a full covariance matrix Σ and a center location μ :

$$G(x) = e^{-\frac{1}{2}(x-\mu)^T \Sigma^{-1}(x-\mu)}. \quad (12)$$

For differentiable rendering optimization, 3DGS decomposes Σ into scaling matrix S and rotation matrix R : $\Sigma = RSS^T R^T$, where S and R are stored by a 3D vector s and a quaternion q respectively. To project these 3D Gaussians to 2D image, given a viewing transformation W , we obtain the 2D covariance matrix Σ' and 2D center location μ' as:

$$\Sigma' = JW\Sigma W^T J^T, \quad \mu' = JW\mu, \quad (13)$$

where J is the Jacobian of the affine approximation of the projective transformation. Then we can use the neural point-based α -blending to render the color C of each pixel with N ordered 3D Gaussians:

$$C = \sum_{i \in N} T_i c_i \alpha_i, \quad (14)$$

where T_i, α_i are calculated as:

$$T_i = \prod_{j=1}^{i-1} (1 - \alpha_j), \quad \alpha_i = o_i e^{-\frac{1}{2}(x-\mu')^T \Sigma'^{-1}(x-\mu')}. \quad (15)$$

Here, o_i is the opacity of the 3D Gaussian. Therefore, the 3D scene can be represented by the parameter set P of 3D Gaussians, where $P = \{G_i : \mu_i, q_i, s_i, c_i, o_i\}$.

A.2 FLUID MECHANICS

Navier-Stokes Equations. Navier-Stokes equations are well-known fundamental equations that describe fluid mechanics:

$$\begin{aligned} \frac{\partial \mathbf{u}_w}{\partial t} + \mathbf{u}_w \cdot \nabla \mathbf{u}_w &= -\frac{1}{\rho_w} \nabla p + \nu \nabla \cdot \nabla \mathbf{u}_w + \mathbf{F}_o, \\ \nabla \cdot \mathbf{u}_w &= 0, \end{aligned} \quad (16)$$

where \mathbf{u}_w is the wind velocity field, ρ_w is the density, p is the pressure, ν is the kinematic viscosity and \mathbf{F}_o is the external force, which represents the influence of objects interacting with the wind field in our work. The first equation describes the momentum conservation, while the second equation enforces the incompressibility constraint.

Continuous Boltzmann Equation. The Boltzmann equation was proven to recover Navier-Stokes equation macroscopically SHAN et al. (2006). It describes fluid dynamics through the time evolution of a mesoscopic distribution function $\mathbf{f}(\mathbf{u}_w, \mathbf{x}_w, t)$:

$$\frac{\partial \mathbf{f}}{\partial t} + \mathbf{u}_w \cdot \nabla \mathbf{f} = \Omega(\mathbf{f}) + \mathbf{F}_o \cdot \nabla_{\mathbf{u}_w} \mathbf{f}, \quad (17)$$

where \mathbf{F}_o represents external forces, \mathbf{f} represents the probability of a particle being at position \mathbf{x} at time t with velocity \mathbf{u} and Ω is the collision term to relax distribution function towards the equilibrium state, typically using the Bhatnagar-Gross-Krook(BGK) model SHAN et al. (2006). By integrating the microscopic properties, the macroscopic physical quantities of the wind field, such as density ρ_w , linear momentum $\rho_w \mathbf{u}_w$, and the stress tensor \mathbf{S}_w can be calculated:

$$\begin{aligned} \rho_w &= \int \mathbf{f} d\mathbf{u}_w, \quad \rho_w \mathbf{u}_w = \int \mathbf{u}_w \mathbf{f} d\mathbf{u}_w, \\ \rho_w \mathbf{S}_{w,\alpha\beta} &= \int \left(\mathbf{u}_w^2 - \frac{1}{3} \delta_{\alpha\beta} \right) \mathbf{f} d\mathbf{u}_w, \end{aligned} \quad (18)$$

756 where $\alpha, \beta \in \{x, y, z\}$ refer to the coordinates of the stress tensor \mathbf{S}_w and δ is Kronecker delta.
 757

758 We note that all numerical solvers for fluid dynamics, including the Lattice Boltzmann Method (LBM)
 759 adopted in our framework, face inherent limitations under highly turbulent flows or when the local
 760 Reynolds number exceeds the stability range of the scheme. In such regimes, solvers may fail to fully
 761 resolve small-scale vortical structures or highly chaotic flow patterns. In our reconstruction pipeline,
 762 LBM is used as a physics-informed constraint to regularize wind estimation toward physically
 763 plausible behaviors. Consequently, our method can reliably reconstruct moderately complex and
 764 unsteady wind fields, but consistent with the intrinsic limitations of the underlying solver, it does not
 765 capture the full spectrum of fine-grained turbulent structures.
 766

766 A.3 CONTINUUM MECHANICS

768 The object deformations are governed by continuum mechanics. In this section, we introduce the
 769 relevant background.
 770

771 **Constitutive Models.** Constitutive models describe the material responses to different mechanical
 772 loading conditions, which provide the stress-strain relations to formulate the governing equations.
 773 We use the Neo-Hookean constitutive model to capture the non-linearity of the object’s response to
 774 the wind force field:
 775

$$\begin{aligned}\psi(\mathbf{F}) &= \frac{\mu}{2} \sum_i [(\mathbf{F}^T \mathbf{F})_{ii} - 1] - \mu \log(J) + \frac{\lambda}{2} \log^2(J), \\ \mathbf{P}(\mathbf{F}) &= \frac{\partial \psi}{\partial \mathbf{F}} = \mu(\mathbf{F} - \mathbf{F}^T) + \lambda \log(J) \mathbf{F}^{-T},\end{aligned}\tag{19}$$

776 where ψ is the hyperelastic energy density function, \mathbf{P} is the first Piola-Kirchhoff stress, \mathbf{F} is the
 777 deformation gradient that encodes local transformations including stretch, rotation, and shear
 778 Bonet & Wood (1997) and $J = \det(\mathbf{F})$. The Lame parameters μ and λ are related to Young’s modulus E and
 779 Poisson’s ratio ν_p :
 780

$$\mu = \frac{E}{2(1 + \nu_p)}, \quad \lambda = \frac{E\nu_p}{(1 + \nu_p)(1 - 2\nu_p)}.\tag{20}$$

781 **Material Point Method.** Material Point Method (MPM)Stomakhin et al. (2013); Hu et al. (2018)
 782 solves the governing equations by transferring physical properties between Lagrangian particles and
 783 Eulerian grid. In particular, we use MLS-MPMHu et al. (2018) as our simulator. Each particle’s
 784 state is described by a five-tuple $(\mathbf{x}_p, \mathbf{m}_p, \mathbf{u}_p, \mathbf{C}_p, \mathbf{F}_p)$: \mathbf{x}_p is the particle position, \mathbf{m}_p is the particle
 785 mass, \mathbf{u}_p is the particle velocity, \mathbf{C}_p is the affine momentumJiang et al. (2015) on particle, \mathbf{F}_p is the
 786 elastic deformation gradient tracked on the particle. MPM operates in a particle-to-grid (P2G) and
 787 grid-to-particle (G2P) transfer loop to simulate the object dynamics. In the P2G stage, we transfer
 788 mass and momentum from particles to the grid as:
 789

$$m_i^t = \sum_p w_{ip}^t m_p, \quad m_i^t \mathbf{u}_i^t = \sum_p w_{ip}^t m_p (\mathbf{u}_p^t + \mathbf{C}_p^t (\mathbf{x}_i - \mathbf{x}_p^t)).\tag{21}$$

790 In the G2P stage, we transfer velocities back to particles and update particle states as:
 791

$$\mathbf{u}_p^{t+1} = \sum_i \mathbf{u}_i^{t+1} w_{ip}^t, \quad \mathbf{x}_p^{t+1} = \mathbf{x}_p^t + \Delta t \mathbf{u}_p^{t+1},\tag{22}$$

792 where i represents the property on the Eulerian grid, n represents the property at time t_n and w_{ip}^n is
 793 the B-spline kernel weight at \mathbf{x}_p^n . The deformation gradient is updated by MLS approximation as:
 794

$$\mathbf{F}_p^{t+1} = (\mathbf{I} + \Delta t \mathbf{C}_p^t) \mathbf{F}_p^t.\tag{23}$$

801 For more details about the simulation algorithm of MLS-MPMHu et al. (2018), we refer the reader to
 802 its original publication.
 803

804 B MORE DETAILS ON 3D REGIONS SEGMENTATION

805 A number of existing methods support 3D Gaussians segmentation. Specifically, we employ a
 806 contrastive learning strategy to lift 2D segmentation results into 3D Gaussians, following Om-
 807 niSeg3DYing et al. (2024). In particular, we use EntitySegLu et al. (2023) to generate instance-level
 808

2D segmentation masks and Grounded-SAM Ren et al. (2024) to generate part-level 2D segmentation masks, where we design prompts for GPT-4o to identify the names of different parts in the scene as input to Grounded-SAM. For each 2D region, we blend its mask with a canonical view reference image, and concatenate it with the reference image as a prompt image, then we guide the physical agent to identify the region in the marked area and explain all possible material types with physical properties, including Poisson’s ratio ν_p and Young’s modulus E . Next, we associate each Gaussian kernel with an optimizable feature vector $h_g \in \mathbb{R}^{16}$, for the features with region-wise labels $\mathbf{L} = \{l_j^i\}$ in 2D level, we maximize the similarity for features with the same region and low similarity between different regions. We design the following loss for \mathbf{L} as:

$$\mathcal{L}_{CC}(\mathbf{L}) = -\frac{1}{N_r} \sum_{i=1}^{N_r} \sum_{j=1}^{n_i} \log \frac{\exp(l_j^i \cdot \bar{l}^i / \phi_i)}{\sum_{k=1}^{N_r} \exp(l_j^i \cdot \bar{l}^k / \phi_k)}, \quad (24)$$

where N_r is the number of regions involved in \mathbf{L} , $n_i = |\{l^i\}|$, l_j^i is the render feature with point index j and region id i , \bar{l}^i is the mean value for l_j^i , and $\phi_i = \frac{\sum_{j=1}^{n_i} \|l_j^i - \bar{l}^i\|_2}{n_i \log(n_i + \alpha)}$ is the region temperature, where $\alpha = 10$ is a smoothing parameter. After training the feature field, we apply HDBSCAN clustering algorithm McInnes et al. (2017) to automatically classify different regions without requiring a predefined number of clusters, we then compute the convex hull of each region to categorize the densified invisible points and Gaussian points that may not have been classified by HDBSCAN and we employ KNN to assign any remaining unclassified points based on their proximity to existing clusters. Finally, we assign physical properties to each 3D region based on the correspondence between 2D and 3D segmentations.

C MORE DETAILS ON DENSIFY INVISIBLE POINTS

Although PhysGaussian proposes using a 3D opacity field from 3D Gaussians for internal filling, in practice, this method often fails to achieve satisfactory results. Moreover, it requires the 3D Gaussian point cloud to be densely distributed and is not suitable for hole completion. Therefore, we adopt the octree-based voxel filling algorithm provided by the Kaolin library to densify invisible points for improved geometric support. The algorithm mainly consists of the following three steps:

(1) Voxelization of 3D Gaussians via Octree Construction. The set of 3D Gaussians is first converted into a voxel representation using a hierarchical algorithm built upon Kaolin’s Structured Point Cloud (SPC), which functions as an octree. The axis-aligned bounding box (AABB) of all Gaussians is enclosed within a cubical root node, which is recursively subdivided in an 8-way manner. For each sub-node, a list of overlapping Gaussian IDs is maintained. Only the nodes that contain relevant Gaussian density are further subdivided and checked for overlaps. This process continues until a specified octree resolution level is reached. The frontier nodes of the octree represent a voxelized shell of the 3D Gaussians. Optionally, nodes whose accumulated opacity values fall below a predefined opacity_threshold are culled to remove low-density regions. As a result, this step yields a voxel shell approximation of the Gaussian surface, without including inner volume content.

(2) Volume Filling via Multi-View Depth Fusion. To convert the voxelized shell into a volumetric solid, the algorithm performs space carving based on rendered depth maps. Specifically, the SPC is ray-traced from an icosahedral set of camera viewpoints to generate depth maps. These depth maps are then fused into a second sparse SPC, where each node of the octree maintains an occupancy state: *empty*, *occupied*, or *unseen*. The occupancy status of each voxel in a regular 3D grid is determined by querying this SPC. Finally, the union of all *occupied* and *unseen* voxels yields a volumetric representation that better captures the object’s solid geometry.

(3) Dense Point Sampling. After volume filling, the 3D Gaussians are now represented as dense voxels that include the object’s interior volume. A point is sampled at the center of each occupied voxel to construct a dense point cloud. Optionally, a small random perturbation is added to each point. This perturbation is constrained to be small enough to ensure the point remains within the bounds of its voxel. By the end of this step, each voxel contains at most one sampled point, resulting in a uniformly distributed and volume-complete point cloud.

864 **D MORE DETAILS ON HOME-LBM ALGORITHM**
865866 The post-collision distribution function in HOME-LBM is given by:
867

$$\begin{aligned}
868 \quad f_i = \rho w_i & \left[1 + \frac{\mathbf{c}_i \cdot \mathbf{u}}{c_s^2} + \frac{\mathbf{H}^{[2]}(\mathbf{c}_i) : \mathbf{S}}{2c_s^4} \right. \\
869 & + \frac{1}{2c_s^6} \left(H_{xx}^{[3]}(\mathbf{c}_i)(S_{xx}u_y + 2S_{xy}u_x - 2u_xu_xu_y) \right. \\
870 & + H_{yy}^{[3]}(\mathbf{c}_i)(S_{yy}u_x + 2S_{xy}u_y - 2u_xu_yu_y) \\
871 & + H_{xz}^{[3]}(\mathbf{c}_i)(S_{xx}u_z + 2S_{xz}u_x - 2u_xu_xu_z) \\
872 & + H_{zz}^{[3]}(\mathbf{c}_i)(S_{zz}u_x + 2S_{xz}u_z - 2u_xu_zu_z) \\
873 & + H_{yz}^{[3]}(\mathbf{c}_i)(S_{zz}u_y + 2S_{yz}u_z - 2u_yu_zu_z) \\
874 & + H_{yz}^{[3]}(\mathbf{c}_i)(S_{yy}u_z + 2S_{yz}u_y - 2u_yu_yu_z) \\
875 & \left. \left. + H_{xy}^{[3]}(\mathbf{c}_i)(S_{xx}u_y + S_{yz}u_x + S_{xy}u_z - 2u_xu_yu_z) \right) \right], \tag{25}
\end{aligned}$$

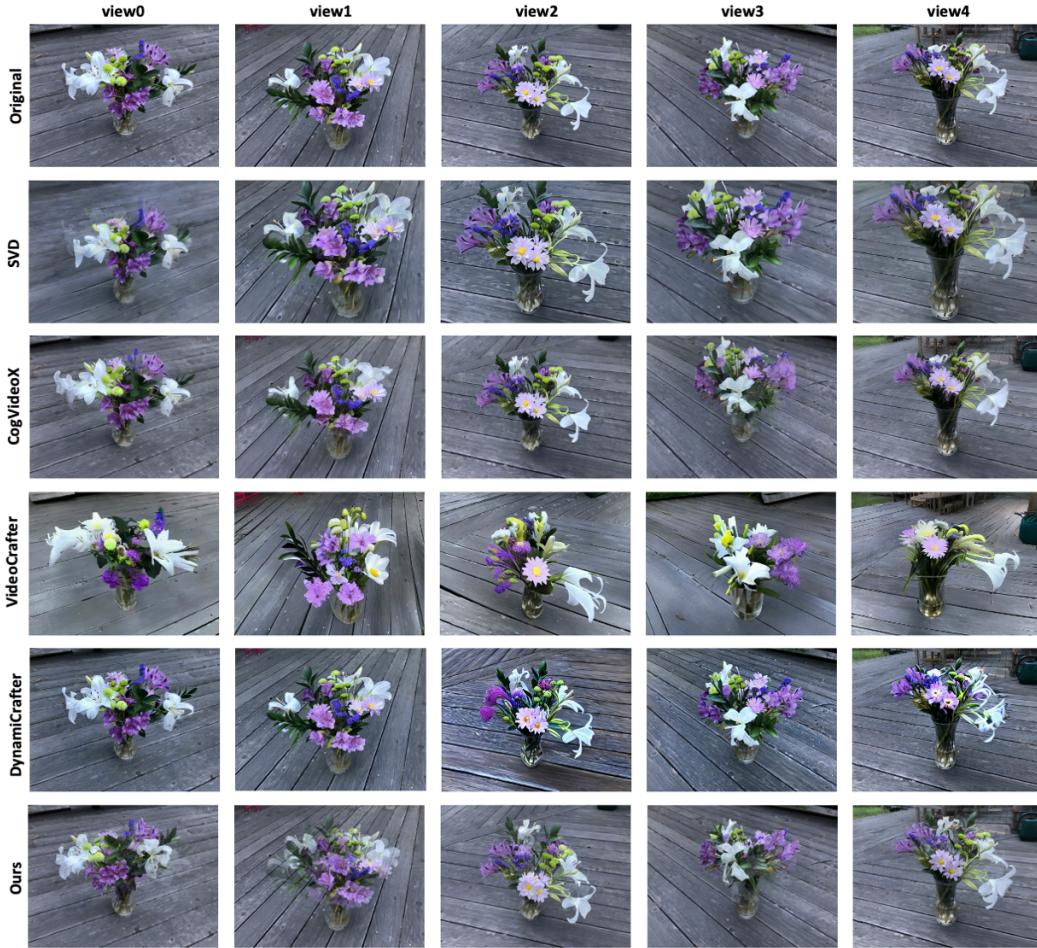
882 where \mathbf{H} is the Hermite polynomials, \mathbf{S} is the stress tensor and \mathbf{u} is the velocity. This post-collision
883 distribution function enables the computation of the microscopic distribution function in the LBM,
884 which can then be integrated into the standard LBM simulation pipeline.885 To address the low accuracy problem caused by the BGK collision model in the classical LBM
886 algorithm, HOME-LBM introduces a high-order collision model based on Hermite expansion after
887 the streaming step the update process is given by:

$$\begin{aligned}
888 \quad \rho(\mathbf{x}, t+1) &= \rho^*, \\
889 \quad u_\alpha(\mathbf{x}, t+1) &= u_\alpha^* + \frac{1}{2\rho^*} F_\alpha, \\
890 \quad S_{xy}(\mathbf{x}, t+1) &= \left(1 - \frac{1}{\tau} \right) S_{xy}^* + \frac{1}{\tau} u_x^* u_y^* + \frac{2\tau - 1}{2\tau\rho^*} (F_x u_y^* + F_y u_x^*), \\
891 \quad S_{xx}(\mathbf{x}, t+1) &= \frac{\tau - 1}{3\tau} (2S_{xx}^* - S_{yy}^* - S_{zz}^*) + \frac{1}{3} (u_x^{*2} + u_y^{*2} + u_z^{*2}) \\
892 & + \frac{1}{\rho^*} F_x u_x^* + \frac{\tau - 1}{3\tau\rho^*} (2F_x u_x^* - F_y u_y^* - F_z u_z^*), \tag{26} \\
893 \quad S_{yy}(\mathbf{x}, t+1) &= \frac{\tau - 1}{3\tau} (2S_{yy}^* - S_{xx}^* - S_{zz}^*) + \frac{1}{3} (u_x^{*2} + u_y^{*2} + u_z^{*2}) \\
894 & + \frac{1}{\rho^*} F_y u_y^* + \frac{\tau - 1}{3\tau\rho^*} (2F_y u_y^* - F_x u_x^* - F_z u_z^*), \\
895 \quad S_{zz}(\mathbf{x}, t+1) &= \frac{\tau - 1}{3\tau} (2S_{zz}^* - S_{xx}^* - S_{yy}^*) + \frac{1}{3} (u_x^{*2} + u_y^{*2} + u_z^{*2}) \\
896 & + \frac{1}{\rho^*} F_z u_z^* + \frac{\tau - 1}{3\tau\rho^*} (2F_z u_z^* - F_x u_x^* - F_y u_y^*).
\end{aligned}$$

905 Here, the superscript * denotes the values obtained after the streaming step and F denotes the force
906 term derived after the streaming step. For a more detailed mathematical derivation, please refer to the
907 original HOME-LBM paper Li et al. (2023a).908 **E MORE EXPERIMENTAL RESULTS**
909910 **E.1 MORE SIMULATION RESULTS**
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918 Table 8: Detailed time statistics for forward simulation. #OBG denotes the number of object particles,
 919 and #BG indicates the number of background Gaussian particles. "LBM Time", "MPM Time", and
 920 "Render Time" represent the average computation time per timestep for LBM simulation, MPM
 921 simulation, and 3D Gaussian Splatting rendering, respectively.

Scene.	#OBG	#BG	LBM Time($\times 10^{-2}$ s)	MPM Time(s)	Render Time($\times 10^{-2}$ s)
Cloth	125223	527159	0.702	0.919	0.598
Garden	52505	4484605	0.710	0.830	3.994
Sock	131794	630068	0.709	0.874	0.669
Vase	145868	1944064	0.720	0.947	1.890



957 Figure 7: Wind synthesis results at the same timestep from multiple views. SVD Blattmann et al.
 958 (2023) exhibits uncontrolled camera movements, VideoCrafter Chen et al. (2024) and Dynami-
 959 crafter Xing et al. (2023) fail to maintain 3D consistency in object geometry and generate only subtle
 960 motion, while CogVideoX Yang et al. (2024a) fails to produce geometrically consistent object motion
 961 across views. In contrast, DiffWind maintains 3D consistency and produces realistic wind-object
 962 interaction.

963 Fig. 7 and Fig. 9 compare our results against various diffusion models, demonstrating that our method
 964 produces more physically plausible and 3D-consistent motion. The time performance of our forward
 965 simulation is reported in Table 8. Please refer to our video supplementary material for more results.

968 E.2 MORE RECONSTRUCTION RESULTS

970 **Baselines.** We compare our approach against the current state-of-the-art dynamic scene reconstruc-
 971 tion methods: Deformable-GS Yang et al. (2024b), Efficient-GS Katsumata et al. (2024), 4D-GS Wu
 et al. (2024) and GaussianPrediction Zhao et al. (2024).



Figure 8: DiffWind is capable of simulating realistic wind-object interactions in static scenes reconstructed from real-world captured multi-view images.

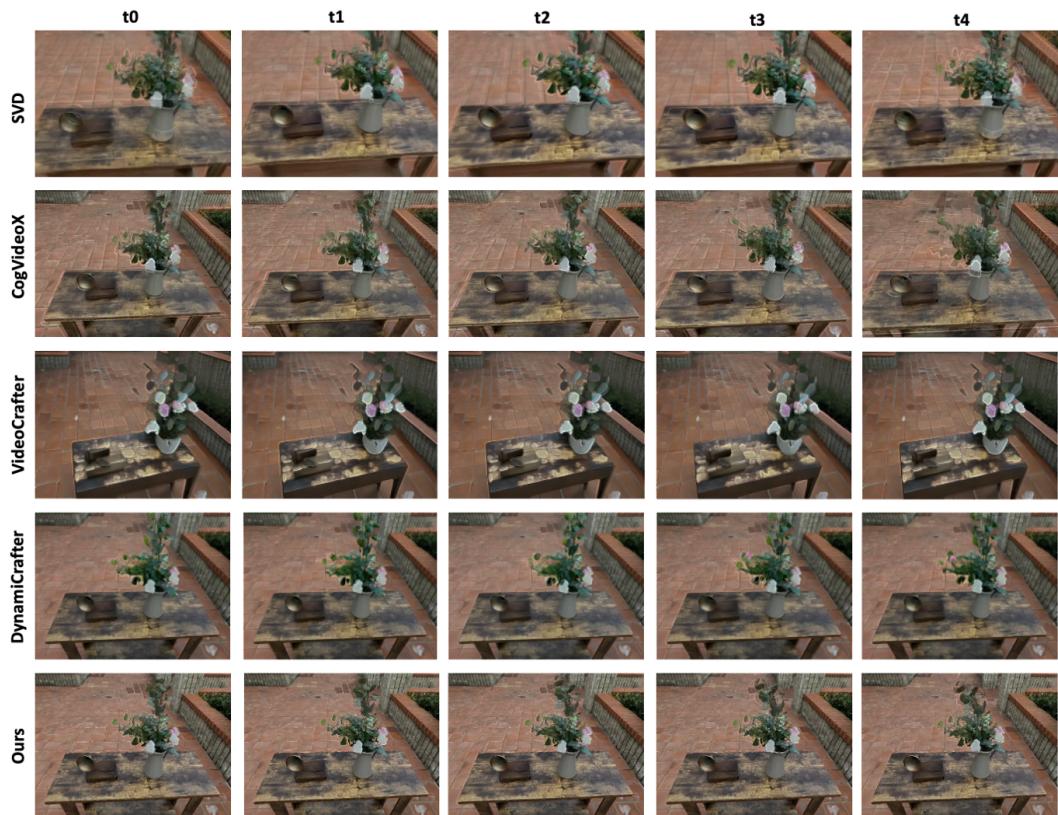


Figure 9: Wind synthesis results at the same viewpoint across different timesteps. SVD Blattmann et al. (2023) suffers from noticeable camera shake, CogVideoX Yang et al. (2024a) produces unrealistic jittering, VideoCrafter Chen et al. (2024) distorts the original object structure and DynamiCrafter Xing et al. (2023) fails to maintain temporal coherence, resulting in implausible motion artifacts. In contrast, DiffWind generates realistic time-evolving wind-object interactions.

Deformable-GS reconstructs dynamic scenes by introducing an additional MLP-based deformation field, while Efficient-GS reconstructs dynamic scenes by learning trajectory functions that govern the motion of Gaussian kernels over time. Similar to Deformable-GS, 4D-GS models the deformation field via HexPlane representations, and GaussianPrediction associates each Gaussian point with an additional motion feature as an extra input to the deformation field. For a fair comparison, we run their optimization process using our reconstructed static 3D Gaussians as initialization, and apply the same regularization terms as described in Sec. 3.2.

Real-World Data. We also capture real-world datasets for evaluation. Specifically, we record 360-degree surround videos of real-world static scenes by GoPro cameras. Each scene includes an object and a background. The objects include a pothos plant, a beanie hat, a pompon flower, and a tulip. Additionally, we capture synchronized videos of these plants being affected by a hairdryer. We report the quantitative comparison results in Table. 2 and show the qualitative results in Fig. 11.

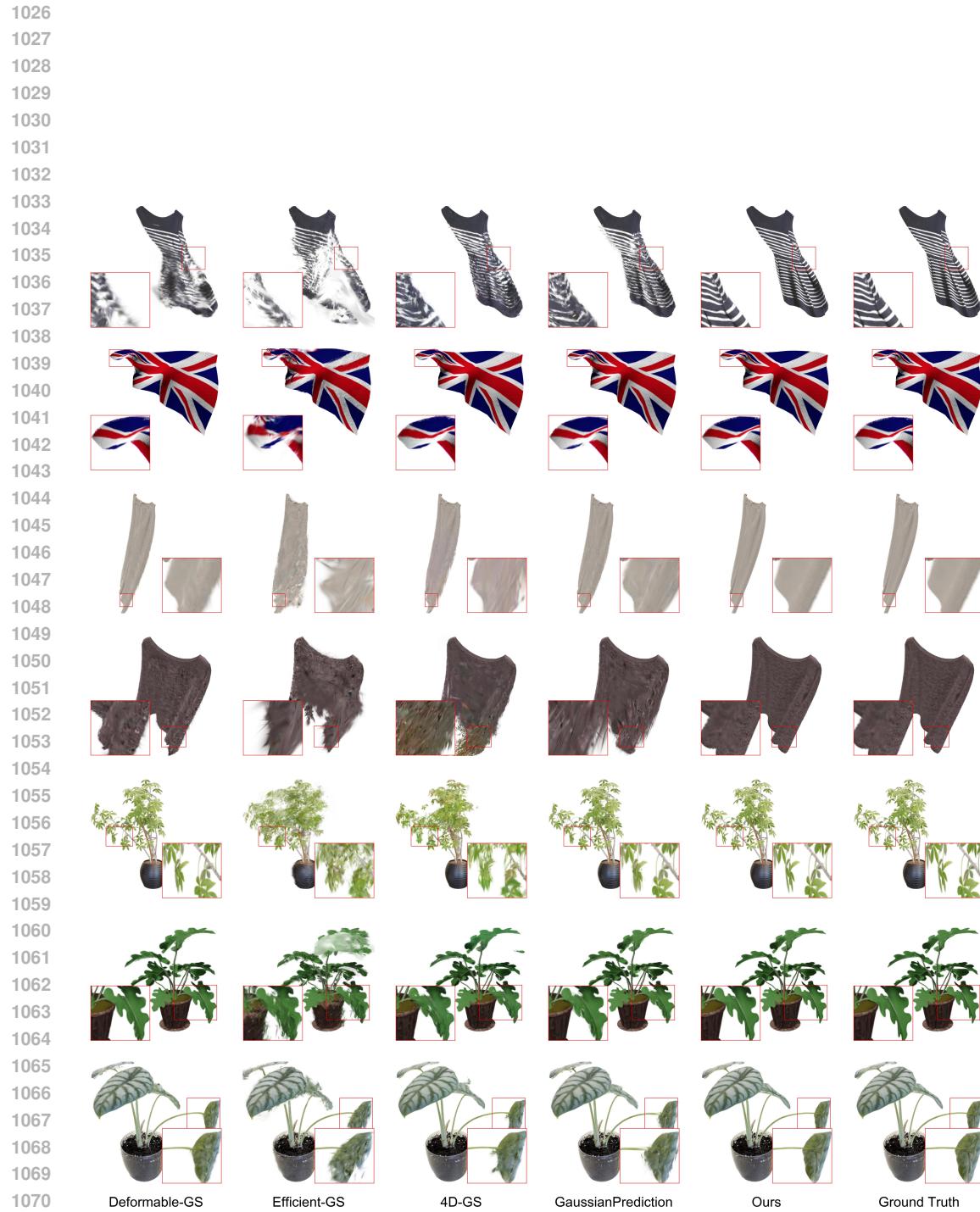


Figure 10: Qualitative results of novel view synthesis on synthetic wind-object interaction scenes. We compare our methods with Deformable-GS Yang et al. (2024b), Efficient-GS Katsumata et al. (2024), 4D-GS Wu et al. (2024) and GaussianPrediction Zhao et al. (2024).

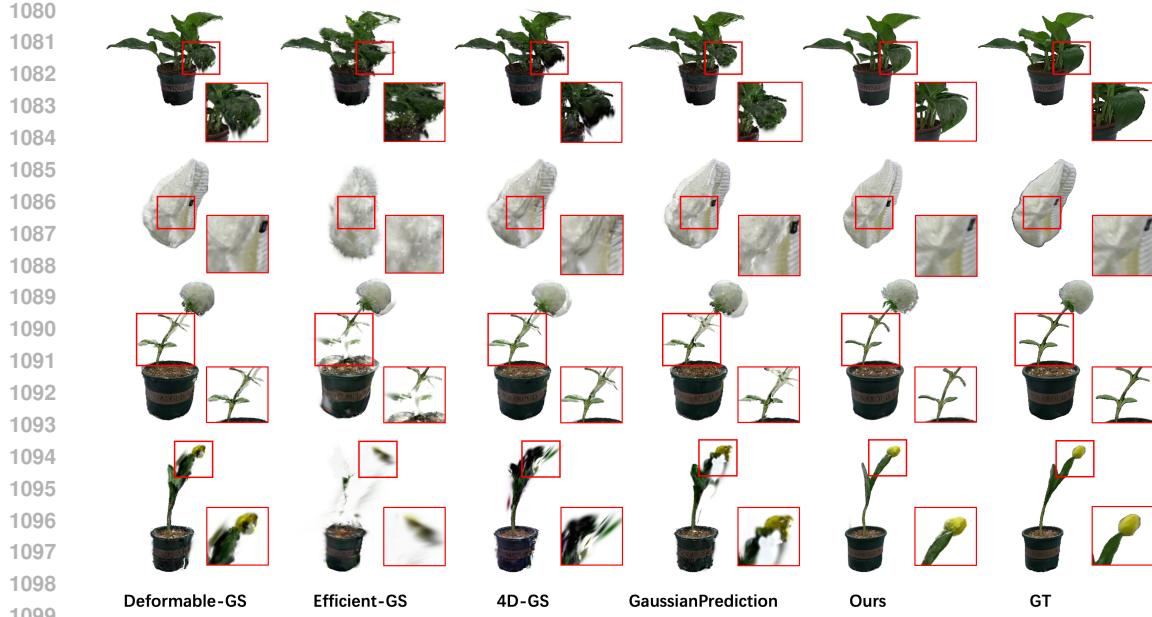


Figure 11: Qualitative results of novel view synthesis on real-world wind-object interaction scenes. We compare our method with Deformable-GS Yang et al. (2024b), Efficient-GS Katsumata et al. (2024), 4D-GS Wu et al. (2024) and GaussianPrediction Zhao et al. (2024).

Table 9: Evaluations of reconstructed wind force fields on the synthetic dataset.

Metrics	Flag	Pants	Sweater	Dress	Ficus	ShapeNet	Alocasia	Average
CosSim	0.9489	0.9585	0.9474	0.8390	0.8695	0.9883	0.9229	0.9249
NMSE	0.0341	0.0276	0.0350	0.1074	0.0870	0.0078	0.0514	0.0500

Ablation on Invisible Point Densification. We inspect the effectiveness of densifying invisible points and conduct the experiments on the real-world dataset. In Fig. 12, the left side shows the original locations of the 3D Gaussians, while the right side shows the complete point set after densification.

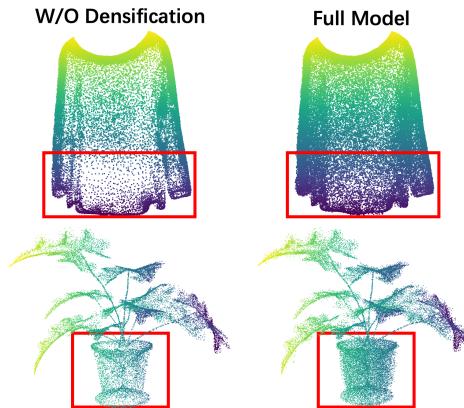


Figure 12: Visualization of object particle positions before and after invisible points densification.

Please generate a caption for the given image that explicitly mentions the different part categories present in the image. The caption should only include the names of the parts (e.g., ‘flower pot’, ‘flowers’) without any detailed descriptions. Ensure the caption is concise and clearly indicates the presence of distinct parts.

Example output:
 • “A potted plant with a flower pot and flowers.”
 • “A car with wheels, windows, and a roof.”
 • “A chair with a seat, backrest, and legs.”

Next, Extract and list the parts from the caption generated in this caption. For each part, provide a simple label or phrase (e.g., ‘flower pot’, ‘flowers’), and separate each part by a period (‘.’). The output should be a single line with all the parts, separated by periods and no extra spaces or new lines.

Example output:
 Extracted parts: “flower pot. flowers.”

Figure 13: Prompt for 3D Segmentation.

Evaluations of Reconstructed Wind Force Fields. We assess the fidelity of the reconstructed wind force field \mathbf{F}_w under controlled synthetic settings, as shown in Table 9. Let \mathbf{F}_w^{gt} and $\mathbf{F}_w^{\text{rec}}$ denote the ground-truth and reconstructed wind force fields. We normalize the force vector at each grid location to unit length, and denote the normalized vectors as $\hat{\mathbf{F}}_i^{\text{gt}}$ and $\hat{\mathbf{F}}_i^{\text{rec}}$. For directional consistency, we report the average cosine similarity over all time steps $\text{CosSim} = \frac{1}{N} \sum_{i=1}^N \hat{\mathbf{F}}_i^{\text{gt}} \top \hat{\mathbf{F}}_i^{\text{rec}}$,

1134 with range $[-1, 1]$, where higher values indicate stronger alignment between reconstructed and
 1135 ground-truth directions. Additionally, we report the normalized mean squared error over all time steps
 1136 $\text{NMSE} = \frac{1}{N} \sum_{i=1}^N \|\hat{\mathbf{F}}_i^{\text{gt}} - \hat{\mathbf{F}}_i^{\text{rec}}\|^2$, with range $[0, 4]$, where lower values indicate better reconstruction.
 1137 These results show that the recovered wind force fields are directionally coherent and maintain
 1138 consistent relative magnitudes with respect to the ground truth on the synthetic dataset.
 1139

1140 **Evaluations of Other Baselines.** Note that there
 1141 are existing methods that also employ differentiable physical simulation for physics parameter esti-
 1142 mation from video observation, including Phys-
 1143 Dreamer Zhang et al. (2024) and PhysFlow Liu et al.
 1144 (2025). However, both of them only optimize the
 1145 initial motion velocities at first frame and Young’s
 1146 modulus for simple and predefined motion patterns,
 1147 which cannot fully characterize the complex dynam-
 1148 ics provided in the input video, especially for windy
 1149 object reconstruction, as shown in Fig. 14. There-
 1150 fore, the results obtained by these methods are
 1151 generally worse than those of the dynamic scene recon-
 1152 struction methods we listed above for novel view
 1153 synthesis. For example, for the Ficus example in our
 1154 dataset, the PSNR/SSIM/LPIPS metrics for Phys-
 1155 Dreamer and PhysFlow are 18.22/0.8546/0.0948 and
 1156 18.11/0.8529/0.0965, respectively, which are consistently worse than other methods. Hence, we focus
 1157 on the comparison of our method to those SOTA dynamic scene reconstruction methods for the novel
 1158 view synthesis task.
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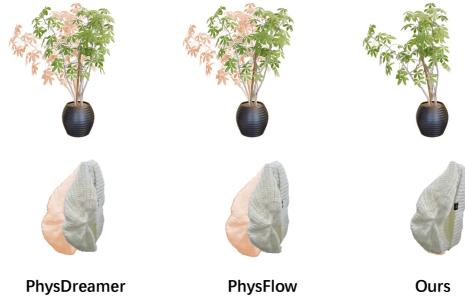
1160 E.3 COMPUTATIONAL COST AND SCALABILITY

1162 All experiments in our paper can be executed on a single NVIDIA RTX 4090 GPU (24GB). Both
 1163 the LBM and MPM simulators are implemented using the Taichi GPU framework, enabling high-
 1164 performance GPU execution of the entire physics pipeline. For forward simulation, DiffWind requires
 1165 less than 1 second per frame for the complete LBM wind update, MPM object simulation, and
 1166 differentiable rendering, demonstrating higher efficiency than existing video generation baselines, for
 1167 example, VideoCrafter requires ~ 81.68 s to generate 16 frames, DynamiCrafter ~ 151.47 s for 16
 1168 frames, and CogVideoX ~ 292.37 s for 50 frames. For reconstruction tasks, each optimization iteration
 1169 includes the LBM update, MPM simulation, differentiable rendering, and gradient computation, and
 1170 runs in ~ 1.17 seconds. The overall reconstruction process takes roughly 2 hours per scene. This
 1171 makes the optimization computationally practical while maintaining high reconstruction fidelity and
 1172 physical consistency. For comparison, 4DGS typically optimizes for ~ 1 hour per scene. While
 1173 DiffWind is slower than 4DGS-style baselines, this is expected because each iteration involves
 1174 full differentiable physics simulation, tackling a substantially more complex problem. Improving
 1175 the efficiency of differentiable physical simulators is a broader challenge for the community and a
 1176 direction for future research, not a limitation unique to our method.
 1177

1178 We also evaluated scalability with respect to grid resolution. GPU memory usage during training
 1179 increases moderately with grid size: 64^3 uses ~ 12.3 GB, 128^3 uses ~ 13.0 GB, 256^3 uses ~ 14.8 GB,
 1180 and 512^3 uses ~ 21.0 GB. This indicates that memory usage remains practical for moderately
 1181 sized scenes, thanks to Taichi’s efficient memory handling. For larger-scale scenarios, multi-GPU
 1182 parallelization and other scalability strategies can be employed.
 1183

1184 F MORE DETAILS ON PROMPT DESIGN

1185 We present the prompt template for precise 3D segmentation, as shown in Fig. 13. This template is
 1186 designed to generate a brief yet accurate segmentation caption when provided with a given image.
 1187 Next, we design a physical agent based on MLLMs for open-vocabulary semantic reasoning about
 1188 materials and their physical properties. For each segmented object part, we combine its mask with a



1189 Figure 14: PhysDreamer and PhysFlow fail
 1190 to reconstruct dynamics for objects with com-
 1191 plex wind-induced motion, with the GT motion
 1192 (red) shown underneath for reference, while
 1193 our method faithfully captures the motion.

1188
 1189 **garden:** A bouquet composed mainly of
 1190 eucalyptus leaves with sprigs of small pink
 1191 and white flowers, contained in a beige
 1192 ceramic pitcher, sits on a dark weathered
 1193 wooden table. Next to it rests a small
 1194 vintage-style gramophone accent piece.
 1195 The camera remains still, no zooming, no
 1196 closing in, no direction-changing, just
 1197 remains still. A strong wind gusts from the
 1198 right, causing the flowers and foliage to
 1199 sway and dance dramatically and naturally
 1200 to the left with significant motion.
 1201
 1202

1203 **vase:** A vibrant bouquet of white lilies,
 1204 purple alstroemeria, and green
 1205 chrysanthemums in a clear glass vase sits
 1206 on weathered wooden planks. The camera
 1207 remains still as a strong wind gusts from the
 1208 right, causing the flowers and foliage to
 1209 sway and dance dramatically to the left
 1210 with significant motion.
 1211
 1212

1213 **sock:** A large, festive Christmas stocking, primarily red and olive green,
 1214 featuring white "Merry Christmas" lettering and an attached Santa Claus
 1215 figurine with a fluffy beard, is clipped by several white clothespins to a dark
 1216 brown wooden stand. The scene is set indoors against a plain off-white wall
 1217 and a light-colored tiled floor. The camera remains perfectly static,
 1218 maintaining a consistent medium shot with no zoom, pan, or tilt. A steady,
 1219 gentle breeze blows from the right, causing the stocking to sway softly and
 1220 naturally to the left. The fabric of the stocking should show slight ripples,
 1221 should exhibit realistic, subtle movements in response to the airflow.
 1222
 1223

1224 **cloth:** A light beige jacket hangs from a white plastic hanger. This hanger is
 1225 hooked onto the side of a small, light-colored wooden shelf-like structure
 1226 that rests atop a stack of three brown cardboard boxes. The background
 1227 features dark, plush armchairs and additional cardboard boxes, indicating
 1228 an indoor setting. The camera maintains a completely static medium shot,
 1229 with no zoom, pan, or tilt. A strong and steady breeze emanates from the
 1230 right, causing the jacket to sway violently and naturally towards the left.
 1231 The fabric of the jacket, especially its sleeves and lower portions, should
 1232 exhibit realistic ripples and movements due to the wind.
 1233
 1234

Figure 15: Prompt for video diffusion models.

1235 canonical reference image and concatenate it with the reference image to form a prompt image. The
 1236 agent is then guided to identify the object part within the marked area and reason about all possible
 1237 material types and their associated physical properties. The prompt template is shown in Fig. 17.
 1238 Note that the MLLM we used here is GPT5. For wind source inference, given the observed RGB
 1239 video, we first generate the monocular depth video through Video Depth AnythingChen et al. (2025),
 1240 and then infer the wind source direction using the prompt as illustrated in Fig. 16.

1241 In addition, we utilize the prompt design shown in Fig. 15 to guide video generation models in
 1242 synthesizing scenes of wind-induced object motion. These synthesized videos are qualitatively
 1243 compared with our physically-based forward simulation results to evaluate visual realism and motion
 1244 plausibility, as presented in the main paper.

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1252 You are an AI model specialized in analyzing RGB video and depth video.
1253 Focus exclusively on the following task and do not generate unrelated content.
1254 Wait until I provide the video input before producing any result.

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Task Description:

You will be given a monocular RGB video and its corresponding depth video that together capture an object being blown by the wind.

Your goal is to estimate the direction of the wind source in 3D space.

⚠ It is essential that you use **both the RGB video and the depth video together**. Do not rely only on RGB appearance or 2D motion — you must incorporate depth information to reconstruct the object's **3D motion trajectory** before inferring the wind source direction.

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Coordinate System:

- **x-axis:** horizontal direction (positive to the right)
- **y-axis:** vertical direction (positive upward)
- **z-axis:** perpendicular to the screen, pointing outward toward the viewer
- The origin is at the **bottom-left corner** of the image.

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Instructions:

1. When the videos are provided, jointly analyze the object's motion using **RGB appearance + Depth structure**.
2. Explicitly use the depth data to lift the motion from 2D image space **into 3D motion space**.
3. Infer the most likely **wind source direction** that causes this motion.
4. Output the result strictly in the following format:

```
'''json
{
  "wind_source_direction": [x, y, z],
  "explanation": "concise explanation showing how both RGB and Depth information together indicate this direction"
}
```

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Figure 16: GPT Prompt for Wind Source Inference.

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You are ChatGPT, a large language model trained by OpenAI, based on the GPT-5 architecture.

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You need to complete a task related to materials in computer graphics. The task here is object recognition and reasoning about the possible material types and material properties (such as Young's modulus and Poisson's ratio) of the object based on the object information you recognized. I will provide you with a image of a [object_name]. This image is divided into two sub-images. The left one is the original image, and the right one uses blue color to mask the object. I need you to: Identify what object the blue masked area corresponds to. Use your identification results to infer the material type of the object. For the identified material, provide the corresponding Young's modulus and Poisson's ratio. # For instance, I give you an image of a potted plant where the area masked in blue corresponds to the pot. You should identify it as the pot, and then infer the material of the pot, such as ceramic, plastic, or metal, based on the characteristics of the pot. Then, provide the Young's modulus and Poisson's ratio for the inferred material (e.g., ceramic: Young's modulus = 70 GPa, Poisson's ratio = 0.25).

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The right sub-image indicates the mask area, and then you could judge based on the left sub-image which is the original image. You must zoom in the corresponding area in the left original image for better recognition. VERY IMPORTANT: For object recognition, you must make your judgment based on this principle: The mask area is all of the object and the area outside the mask does not belong to the object. Let's think step by step for object recognition.

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IMPORTANT: Please also combine image information when reasoning about materials, and sort the results according to likelihood. Let's think step by step for material reasoning.

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To provide an answer, please provide a short analysis for object recognition and material reasoning. Analysis: - Object recognition: The object in the masked area is [object name], because [reasoning for identification]. - Material reasoning: The possible material types are [material 1], [material 2], ..., because [reasoning for material types]. Material properties: - Young's modulus: [value] GPa (for identified material) - Poisson's ratio: [value] (for identified material) Final answer: [material 1] [material 2] [material 3] ... (sorted by likelihood) Example Output:

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Analysis:

- Object recognition: The object in the masked area is the pot, because its shape and appearance match typical pot structures.
- Material reasoning: The possible material types are ceramic, plastic, metal, because pots are commonly made from these materials, based on the visual texture and color.

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Material properties:

- Young's modulus: 70 GPa (for ceramic)
- Poisson's ratio: 0.25 (for ceramic)

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Final answer: ceramic plastic metal

Figure 17: GPT Prompt for Physical Properties Reasoning.