

# 000 001 002 003 004 005 006 007 008 009 010 011 012 013 THE POLAR EXPRESS: OPTIMAL MATRIX SIGN METHODS AND THEIR APPLICATION TO THE MUON ALGORITHM

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## 011 ABSTRACT

014 Computing the polar decomposition and the related matrix sign function has  
 015 been a well-studied problem in numerical analysis for decades. Recently, it has  
 016 emerged as an important subroutine within the Muon optimizer for training deep  
 017 neural networks. However, the requirements of this application differ sharply from  
 018 classical settings: deep learning demands GPU-friendly algorithms that prioritize  
 019 high throughput over high precision. We introduce Polar Express, a new  
 020 method for computing the polar decomposition. Like Newton-Schulz and other  
 021 classical polynomial methods, our approach uses only matrix-matrix multiplica-  
 022 tions, making it very efficient on GPUs. Inspired by earlier work of Chen & Chow  
 023 and Nakatsukasa & Freund, Polar Express adapts the update rule at each it-  
 024 eration by solving a minimax optimization problem. We prove that this strategy  
 025 minimizes error in a worst-case sense, allowing Polar Express to converge  
 026 as rapidly as possible both in the early iterations and asymptotically. We also  
 027 address finite-precision issues, making it practical to use in bfloat16. **When**  
 028 **integrated into Muon, our method yields consistent improvements in validation**  
**loss for a GPT-2 model trained on one to ten billion tokens from the FineWeb**  
**dataset**, outperforming recent alternatives across a range of learning rates.

## 033 1 INTRODUCTION

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 035 Advanced linear algebra is making its way into deep learning. Efficient algorithms for computing  
 036 *matrix functions* have found exciting new applications in training neural networks. In particular, ap-  
 037 proximations to the matrix-inverse are used in the full Adagrad method (Duchi et al., 2011), the ma-  
 038 trix square-root and quarter-root appear as subroutines in the Shampoo and Soap optimizers (Gupta  
 039 et al., 2018; Shi et al., 2023; Vyas et al., 2025), and most recently, the matrix sign function has  
 040 become a key ingredient of the Muon optimizer (Bernstein & Newhouse, 2024b;a; Jordan et al.,  
 041 2024b). While the problem of computing these matrix functions has been studied by numerical  
 042 analysts for decades, applications in deep learning come with different requirements than those in  
 043 computational science. For deep learning, it is critical to take maximum advantage of GPU-friendly  
 044 operations like matrix-matrix products and to avoid less parallel operations. Moreover, memory  
 045 overhead must be small to handle large models. On the other hand, high accuracy is typically less  
 046 important; the gold standard of sixteen digits of accuracy is overkill in deep learning.

047 Given these considerations, there is a need to develop new matrix function methods that are tailor-  
 048 made for deep learning applications. We take on this challenge by designing a state-of-the-art,  
 049 GPU-friendly algorithm for computing the matrix sign function, or more generally, for computing  
 050 the *polar decomposition* of a rectangular matrix. We apply our new Polar Express method  
 051 ([Algorithm 1](#), [Implementation 1](#)) to compute the descent direction in the increasingly popular Muon  
 052 optimizer. In [Figure 1](#), we show that using Polar Express within Muon consistently results in  
 053 lower validation loss across all learning rates when training a GPT-2 model, as compared to other  
 matrix sign methods ([Cesista et al., 2025](#); [Jordan et al., 2024b](#)).

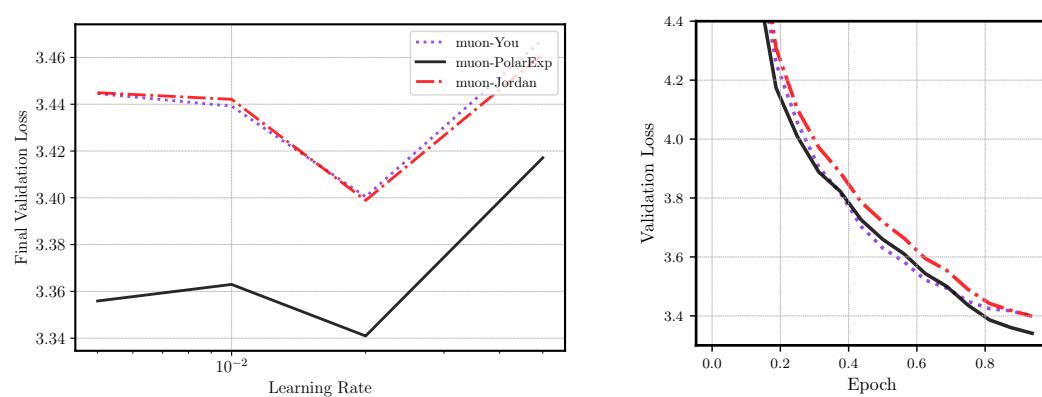


Figure 1: Training a GPT-2-Large model (774M params) on 1 billion tokens from the FineWeb dataset (Aroca-Ouellette et al., 2023). The label muon-<name> refers to implementing Muon using <name> to compute the polar factor. Left: final validation loss across learning rates. Right: validation loss across epochs using the best learning rate. The best learning rate ( $lr$ ) and final validation loss for each method were muon-You ( $lr = 0.02$ ): 3.399, muon-Jordan ( $lr = 0.02$ ): 3.398 and muon-PolarExp ( $lr = 0.02$ ): 3.340.

### 1.1 THE MUON METHOD

The Muon optimizer has recently gained popularity for training large language models, often outperforming state-of-the-art adaptive gradient methods like Adam and AdamW (Kingma & Ba, 2015; Loshchilov & Hutter, 2019). Muon has been used to set records for the NanoGPT speedrun (Jordan et al., 2024b), to expand the Pareto frontier of performance versus training FLOPs for large language models (Liu et al., 2025; Shah et al., 2025), and even to train a 1 trillion parameter frontier LLM (Kimi Team et al., 2025).

The Muon update rule (Bernstein & Newhouse, 2024b) is defined as follows. Let  $\lambda, \beta > 0$  be the learning rate and momentum coefficient hyperparameters. (By default,  $\beta = 0.9$ .) Let  $\mathbf{W}_t \in \mathbb{R}^{m \times n}$  be the weight matrix of a given neural network layer at iteration  $t$ , and let  $\mathbf{G}_t \in \mathbb{R}^{m \times n}$  be its (stochastic) gradient. Let  $\mathbf{M}_t \in \mathbb{R}^{m \times n}$  be the running momentum estimate of the gradient, where  $\mathbf{M}_0 = \mathbf{0}$ . The Muon update is given by

$$\mathbf{M}_t = \beta \mathbf{M}_{t-1} + (1 - \beta) \mathbf{G}_t, \quad \mathbf{W}_{t+1} = \mathbf{W}_t - \lambda \text{polar}(\mathbf{M}_t).$$

Whereas standard stochastic gradient descent (SGD) with momentum updates the weight matrix by taking a step in the direction  $-\mathbf{M}_t$ , the Muon method steps in the direction  $-\text{polar}(\mathbf{M}_t)$ , where  $\text{polar}(\mathbf{M})$  denotes the closest semi-orthogonal matrix to  $\mathbf{M}$  (Higham, 2008, Chapter 8). Concretely, if  $\mathbf{M} = \mathbf{U}\Sigma\mathbf{V}^\top$  is the singular value decomposition (SVD) of  $\mathbf{M}$ , then

$$\text{polar}(\mathbf{M}) := \mathbf{U}\mathbf{V}^\top. \quad (1)$$

The matrix  $\text{polar}(\mathbf{M})$  can be seen as a generalization of the matrix sign function to rectangular matrices (Benzi & Huang, 2019). Indeed, when  $\mathbf{M}$  is square symmetric with eigendecomposition  $\mathbf{M} = \mathbf{V}\Lambda\mathbf{V}^\top$ ,  $\text{polar}(\mathbf{M})$  exactly coincides with the matrix sign function  $\text{sign}(\mathbf{M}) = \mathbf{V}\text{sign}(\Lambda)\mathbf{V}^\top$  (Higham, 2008, Chapter 5). Equivalently,  $\text{polar}(\mathbf{M})$  is the left orthogonal factor of the polar decomposition of  $\mathbf{M}$  (Higham, 2008, Chapter 8). The motivation for Muon is that  $-\text{polar}(\mathbf{M})$  gives the steepest-descent direction with respect to the *spectral norm* (instead of the Frobenius norm, as in standard SGD). For analysis and further discussion on Muon we refer the reader to (Jordan et al., 2024b; Bernstein & Newhouse, 2024b; Pethick et al., 2025; Riabinin et al., 2025; Carlson et al., 2015a,b). In this paper, we take the Muon update rule as given and focus on the problem of efficiently computing the polar decomposition  $\text{polar}(\mathbf{M})$ .

### 1.2 COMPUTING THE POLAR FACTOR

Although  $\text{polar}(\mathbf{M})$  can be computed directly via an SVD in  $O(mn \min(m, n))$  time, doing so is prohibitively expensive in deep learning applications, especially as standard SVD algorithms fail to

108 take full advantage of the parallelism available on GPUs. There has been significant work on highly-  
 109 parallel methods for the SVD, but the most common approaches actually require computing the  
 110 matrix-sign function as a subroutine (Nakatsukasa & Freund, 2016; Nakatsukasa & Higham, 2013).  
 111 Numerical analysts have spent decades developing iterative methods for computing  $\text{polar}(\mathbf{M})$ . This  
 112 rich line of work includes Newton-Schulz (Higham, 2008, Chapter 8), Padé iteration (Kenney &  
 113 Laub, 1991; Higham, 1986), the Newton and scaled Newton iterations (Higham, 2008, Chapter 8),  
 114 the QDWH iteration (Nakatsukasa et al., 2010; Nakatsukasa & Higham, 2013), and Zolo-pd (Nakat-  
 115 sukasa & Freund, 2016). Unfortunately, as discussed in [Appendix B](#), most of these methods are  
 116 based on rational approximations to the function  $\text{sign}(x)$  and require computing matrix inverses  
 117 or QR decompositions. Such methods are ill-suited to GPU acceleration and deep learning applica-  
 118 tions. In contrast, the older Newton-Schulz method is based on *polynomial* approximation of  $\text{sign}(x)$   
 119 and uses only matrix-matrix products. Thus, Muon initially used Newton-Schulz (Bernstein & New-  
 120 house, 2024a). Indeed, Muon stands for “MomentUm Orthogonalized by Newton-Schulz” (Jordan  
 121 et al., 2024b). For a more comprehensive discussion on prior work, see [Appendix B](#).

122 **The Newton-Schulz methods.** Newton-Schulz constructs a sequence of approximations  $\mathbf{X}_t \approx$   
 123  $\text{polar}(\mathbf{M})$  as follows:

$$\mathbf{X}_0 = \mathbf{M} / \|\mathbf{M}\|_{\text{F}}, \quad \mathbf{X}_{t+1} = \frac{3}{2}\mathbf{X}_t - \frac{1}{2}\mathbf{X}_t \mathbf{X}_t^{\top} \mathbf{X}_t. \quad (2)$$

124 At each iteration, this rule effectively applies the cubic polynomial  $p(x) = \frac{3}{2}x - \frac{1}{2}x^3$  to each sin-  
 125 gular value of  $\mathbf{X}_t$ . The scalar fixed-point iteration  $x_{t+1} = p(x_t)$  converges to  $\text{sign}(x_0)$  as  $t \rightarrow \infty$ ,  
 126 provided  $|x_0| \leq 1$ . As a result, the matrix iteration satisfies  $\lim_{t \rightarrow \infty} \mathbf{X}_t = \mathbf{U} \mathbf{V}^{\top} = \text{polar}(\mathbf{X}_0)$ .

127 Higher-degree versions of Newton-Schulz follow the same principle. For example, the degree-5  
 128 polynomial  $p(x) = (15x - 10x^3 + 3x^5)/8$  converges even faster. The Newton-Schulz iterations  
 129 converge super-exponentially when  $\mathbf{X}_t$  is sufficiently close to  $\text{polar}(\mathbf{M})$ , but they suffer from slow  
 130 initial convergence; when  $\mathbf{X}_0$  is far from  $\text{polar}(\mathbf{M})$ , the approximation improves slowly over the  
 131 first few iterations. Due to the slow initial convergence of Newton-Schulz, Chen & Chow (2014)  
 132 developed a version of the Newton-Schulz iteration, which adapts the polynomial at each iteration.  
 133 The resulting method achieves a faster initial convergence, while retaining super-exponential con-  
 134 vergence in later iterations. Polar Express is inspired by their method.

135 **The Jordan and You methods.** In Muon, high accuracy approximations to  $\text{polar}(\mathbf{M})$  are usually  
 136 not necessary. The primary goal is instead to compute a coarse approximation in as few iterations  
 137 as possible. To accelerate convergence in the low-accuracy regime, Jordan recently proposed a  
 138 fixed-point iteration based on the polynomial  $p(x) = 3.4445x - 4.7750x^3 + 2.0315x^5$ , which was  
 139 found using a heuristic numerical search (Jordan et al., 2024b). Unlike Newton-Schulz, the scheme  
 140 that Jordan proposed does not converge to  $\text{polar}(\mathbf{M})$ , but plateaus at an error of  $\approx 0.3$ . However,  
 141 it reaches this level of accuracy rapidly and outperforms the Newton-Schulz when only a small  
 142 number of iterations are performed. Building on this idea, You proposed a method that applies  
 143 six different polynomial updates in succession, which were again found by heuristic search. This  
 144 method achieves better accuracy than Jordan’s but still fails to converge (Cesista et al., 2025).

### 145 1.3 CONTRIBUTIONS

146 We present Polar Express ([Algorithm 1](#)), an iterative method for approximating  $\text{polar}(\mathbf{M})$ .  
 147 Our method dynamically adapts the polynomial update rule at each iteration, prioritizing rapid  
 148 progress in the initial stage and high accuracy in the later stage. Polar Express constructs  
 149 polynomials  $p_1, \dots, p_T$  so that the resulting composition is the optimal approximation to the sign  
 150 function with respect to the supremum ( $L^{\infty}$ ) norm ([Theorem 3.1](#)). By iteratively applying these  
 151 polynomials to  $\mathbf{M}$ , Polar Express computes an approximation to  $\text{polar}(\mathbf{M})$  that is optimal  
 152 in the worst-case. Our method converges to  $\text{polar}(\mathbf{M})$  super-exponentially ([Theorem 3.3](#)), and it  
 153 quickly reaches a good approximation within just five to ten iterations. This early-stage acceleration  
 154 is especially valuable in deep learning applications, where runtime efficiency takes precedence over  
 155 high accuracy. In contrast, classical methods like Newton-Schulz suffer from a slow initial conver-  
 156 gence, while recent heuristic proposals (Jordan et al., 2024b; Cesista et al., 2025) fail to converge.

Our method is efficient to run on GPUs, using only a few matrix-matrix products per iteration.<sup>1</sup> We give an explicit instantiation of `Polar Express` in [Implementation 1](#), which incorporates minor modifications to make it compatible with half-precision arithmetic (see [Section 3.4](#)). [Implementation 1](#) is very short and easy to use, with no dependencies except PyTorch. It serves as a drop-in replacement for previous methods. In numerical experiments, `Polar Express` outperforms previous methods on synthetic matrices and gradient matrices from a GPT-2 transformer ([Figure 3](#)). We demonstrate the effectiveness of using `Polar Express` within the `Muon` optimizer in [Figure 1](#), showing that it consistently improves the training of GPT-2 language models on 1 billion tokens of the FineWeb dataset ([Aroca-Ouellette et al., 2023](#)). Our method has been adopted into the `NanoGPT speedrun` ([Jordan et al., 2024a](#)), a heavily optimized implementation that serves as a benchmark for LLM training efficiency.

**Notation.** We let  $\|\mathbf{M}\|_F$  and  $\|\mathbf{M}\|_2$  denote the Frobenius norm and spectral norm (largest singular value) of a matrix  $\mathbf{M}$ , respectively. We denote the spectrum (set of singular values) by  $\sigma(\mathbf{M})$ . Let  $\mathbb{P}_d$  be the set of polynomials of degree at most  $d$ . For odd  $d$ ,  $\mathbb{P}_d^{\text{odd}}$  denotes the set of polynomials of degree at most  $d$  containing only odd-degree monomials. For a polynomial  $p$ ,  $\deg(p)$  is its degree. Let  $\text{sign}(x)$  be the scalar sign function, which satisfies  $\text{sign}(0) = 0$ ,  $\text{sign}(x) = 1$  if  $x > 0$  and  $\text{sign}(x) = -1$  if  $x < 0$ . For a polynomial  $p \in \mathbb{P}_d^{\text{odd}}$  and a matrix  $\mathbf{M}$  with rank reduced SVD given by  $\mathbf{M} = \mathbf{U}\Sigma\mathbf{V}^T$  and positive singular values  $\sigma_1 \geq \dots \geq \sigma_{\text{rank}(\mathbf{M})} > 0$ , we define  $p(\mathbf{M}) := \mathbf{U}p(\Sigma)\mathbf{V}^T$ , where  $p(\Sigma)$  is the diagonal matrix with diagonal entries  $p(\sigma_i)$  for  $i = 1, \dots, \text{rank}(\mathbf{M})$ .

## 2 APPROXIMATIONS BY COMPOSITIONS OF POLYNOMIALS

To design a GPU-friendly method for computing  $\text{polar}(\mathbf{M})$ , we limit ourselves to the following GPU-friendly operations: (i) linear combinations of matrices (given scalars  $\beta, \gamma \in \mathbb{R}$  and matrices  $\mathbf{B}$  and  $\mathbf{C}$ , compute  $\beta\mathbf{B} + \gamma\mathbf{C}$ ) and (ii) matrix-matrix products (compute  $\mathbf{B}\mathbf{C}$ ). While both these computational primitives are well-suited for parallel computing environments, matrix-matrix products come at a higher computational cost than linear combinations. Therefore, our method attempts to minimize the number of matrix-matrix products. A key observation is that we can compute *odd* monomials of  $\mathbf{M} = \mathbf{U}\Sigma\mathbf{V}^T$  using the following formula:  $\mathbf{M}^{2q+1} := \mathbf{U}\Sigma^{2q+1}\mathbf{V}^T = \mathbf{M}(\mathbf{M}^T\mathbf{M})^q$ .<sup>2</sup> Hence, for an odd polynomial  $p(x) = a_0x + a_1x^3 + \dots + a_qx^{2q+1}$  we can compute

$$p(\mathbf{M}) := a_0\mathbf{M} + a_1\mathbf{M}(\mathbf{M}^T\mathbf{M}) + \dots + a_q\mathbf{M}(\mathbf{M}^T\mathbf{M})^q.$$

It has been shown that for an arbitrary polynomial  $p$ , one requires  $\Theta(\deg(p)^{1/2})$  products to compute  $p(\mathbf{M})$  ([Paterson & Stockmeyer, 1973](#)); see also [Jarlebring & Lorentzon \(2025\)](#) for related work. This compares favorably to the naive approach that forms all monomials in  $p$  and then sums them together, which requires  $\Omega(\deg(p))$  products. However, if  $p$  can be expressed as a composition of  $T$  polynomials, each of degree  $d$

$$p = p_T \circ p_{T-1} \circ \dots \circ p_1, \tag{3}$$

then the degree of  $p$  is  $d^T$ , and  $p(\mathbf{M})$  can be efficiently computed recursively by

$$\mathbf{X}_0 = \mathbf{M}, \quad \mathbf{X}_t = p_t(\mathbf{X}_{t-1}) \text{ for } t = 1, 2, \dots, T. \tag{4}$$

The final iterate is  $\mathbf{X}_T = p(\mathbf{M})$ , which we compute with just  $O(Td)$  matrix-matrix products. Iterative methods for  $\text{polar}(\mathbf{M})$  can be seen in this light. For instance, the degree-5 Newton-Schulz method uses the polynomial update  $p_t(x) = \frac{15}{8}x - \frac{10}{8}x^3 + \frac{3}{8}x^5$  for each  $t = 1, \dots, T$ . The composition  $p = p_T \circ \dots \circ p_1$  approximates  $\text{sign}(x)$ , and the approximation error goes to 0 as  $T$  grows. In this paper, we ask the following question: what choice of  $p_T \circ \dots \circ p_1$  gives the *best* approximation to  $\text{sign}(x)$ ?

The method we will present is optimal in the following sense: given lower and upper bounds  $\ell$  and  $u$  on the singular values of  $\mathbf{M}$ , an odd degree  $d \in \mathbb{N}$ , and the number of iterations  $T \in \mathbb{N}$ , our method computes the composition  $p^*(\mathbf{M})$  that minimizes the worst-case error in the spectral norm. That is,

$$p^* = \arg \min_{\substack{p=p_T \circ p_{T-1} \circ \dots \circ p_1 \\ p_t \in \mathbb{P}_d^{\text{odd}}}} \max_{\substack{\mathbf{M} \in \mathbb{R}^{m \times n} \\ \sigma(\mathbf{M}) \subset [\ell, u]}} \|\text{polar}(\mathbf{M}) - p(\mathbf{M})\|_2. \tag{5}$$

<sup>1</sup>In [Appendices I and J](#), we describe two further algorithmic ideas. They are not used in our `Muon` experiments but they may be beneficial in other settings, and we believe they merit further study.

<sup>2</sup>For non-symmetric matrices, e.g. rectangular matrices, we cannot compute even polynomials of the singular values without first explicitly computing the SVD. We are therefore restricted to odd polynomials.

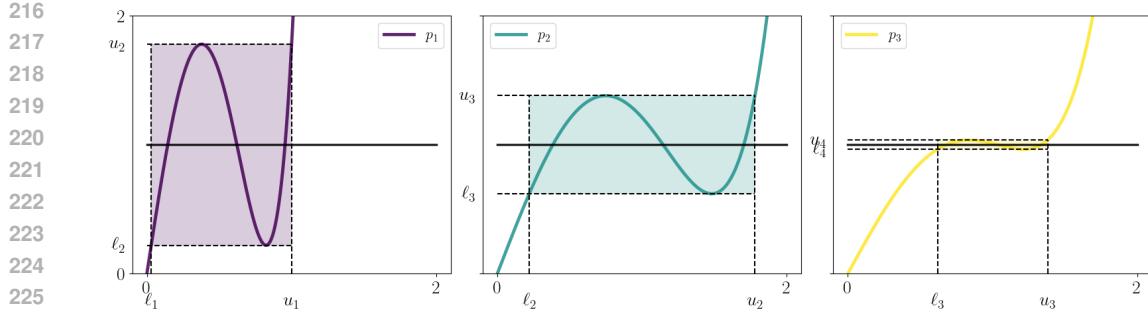


Figure 2: The evolution of the first three optimal polynomials  $p_1$ ,  $p_2$ , and  $p_3$  and the corresponding lower bounds  $\ell_{t+1} = p_t(\ell_t)$  and upper bounds  $u_{t+1} = 2 - \ell_{t+1}$ , as described in [Theorem 3.1](#). The horizontal black line shows  $y = 1$ . The polynomial degree is  $d = 5$ . We set  $\ell_1 = 0.03$  and  $u_1 = 1$ .

Given that  $\text{polar}(\mathbf{M}) - p(\mathbf{M}) = \mathbf{U}(\mathbf{I} - p(\Sigma))\mathbf{V}^\top$ , and by the unitary invariance of the spectral norm, we have that (5) is equivalent to<sup>3</sup>

$$p^* = \arg \min_{\substack{p = p_T \circ p_{T-1} \circ \dots \circ p_1 \\ p_t \in \mathbb{P}_d^{\text{odd}}}} \max_{x \in [\ell, u]} |1 - p(x)|. \quad (6)$$

In other words, the problem given in (5) reduces to that of finding a “uniform” approximation to the constant function  $x \mapsto 1$  over the interval  $[\ell, u]$ , as given in (6). Uniform approximation on an interval by polynomials or rational functions of a given degree is a central topic in approximation theory ([Trefethen, 2020](#)). Here, we seek an approximation of a particular form—a *composition* of odd polynomials of fixed degrees. In the next section, we solve the optimization problem of (6) and use the solution to create Polar Express.

### 3 THE POLAR EXPRESS

#### 3.1 GREEDY IS OPTIMAL

The key observation is that the polynomial used in each iteration can be chosen greedily, given the choice of polynomials from the previous iterations. For the first iteration, we choose  $p_1$  so as to map the interval  $[\ell, u]$  as close to 1 as possible. That is, it minimizes  $\max_{x \in [\ell, u]} |1 - p_1(x)|$ . The image of  $p_1$  will be a new interval  $[\ell_2, u_2]$ , where

$$\ell_2 = \min_{x \in [\ell, u]} p_1(x) \quad u_2 = \max_{x \in [\ell, u]} p_1(x) \quad (7)$$

We now pick  $p_2$  to map the interval  $[\ell_2, u_2]$  as close to 1 as possible, obtaining a new interval  $[\ell_3, u_3]$  that is the image of  $[\ell, u]$  through  $p_2 \circ p_1$ . We continue this process for as many iterations as desired.

The following theorem guarantees that this process finds the solution to (6), and thereby also (5). The scheme is also outlined in [Figure 2](#), which demonstrates the evolution of the lower bounds  $\ell_t$ , the upper bounds  $u_t$ , and the polynomials  $p_t$  across iterations. The proof is in [Appendix C](#).

**Theorem 3.1.** Let  $d$  be odd and define  $\ell_1 = \ell$  and  $u_1 = u$ . For  $t = 1, \dots, T$  define

$$p_t = \arg \min_{p \in \mathbb{P}_d^{\text{odd}}} \max_{x \in [\ell_t, u_t]} |1 - p(x)|, \quad \ell_{t+1} = \min_{x \in [\ell_t, u_t]} p_t(x), \quad u_{t+1} = \max_{x \in [\ell_t, u_t]} p_t(x) \quad (8)$$

The resulting composition  $p^* := p_T \circ p_{T-1} \circ \dots \circ p_1$  is optimal and the error is given by:

$$\max_{x \in [\ell, u]} |1 - p^*(x)| = \min_{\substack{p = p_T \circ p_{T-1} \circ \dots \circ p_1 \\ p_t \in \mathbb{P}_d^{\text{odd}}}} \max_{x \in [\ell, u]} |1 - p(x)| = 1 - \ell_{T+1}. \quad (9)$$

<sup>3</sup>For completeness, the equivalence between (5) and (6) is proven in [Appendix E](#).

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Furthermore the new error, lower and upper bounds can be computed through

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**Remark 3.2** (Why a fixed degree?). We note that choice of the degree of each  $p_1, p_2, \dots, p_T$  need not be the same for [Theorem 3.1](#) to hold. More generally, one may specify a sequence of degrees  $d_1, \dots, d_T$  and define each  $p_t$  as  $p_t = \arg \min_{p \in \mathbb{P}_d^{\text{odd}}} \max_{x \in [\ell_t, u_t]} |p(x) - 1|$  for  $t = 1, \dots, T$ . However, [Lee et al. \(2022, Table 2\)](#) supports setting  $d_t = 5$ , as we do.

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Fortunately, (10) shows that once  $p_t$  has been found, we can compute the new lower and upper bounds  $\ell_{t+1}$  and  $u_{t+1}$  simply by evaluating  $p_t(\ell_t)$ . Hence, for any *fixed* upper and lower bounds on the singular values of  $\mathbf{M}$ , we can *precompute* all the polynomials  $p_1, \dots, p_T$  and the bounds  $[\ell_1, u_1], \dots, [\ell_{T+1}, u_{T+1}]$ . Then, applying the iterative procedure of (4), the final iterate  $\mathbf{X}_T$  will satisfy the following error bound:

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$$\|\text{polar}(\mathbf{M}) - \mathbf{X}_T\|_2 = \|\text{polar}(\mathbf{M}) - p^*(\mathbf{M})\|_2 \leq 1 - \ell_{T+1}. \quad (11)$$

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From the optimality guarantee of [Theorem 3.1](#), we know that our method converges at least as fast as the Newton-Schulz iteration of the same degree. Combining this fact with an existing analysis of Newton-Schulz, we immediately get the following convergence guarantee showing that our method enjoys faster than exponential convergence. The proof can be found in [Appendix D](#).

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**Theorem 3.3.** Let  $\mathbf{M}$  be a matrix normalized so that  $\sigma(\mathbf{M}) \subset [\ell, 1]$ . Let  $\mathbf{X}_T = p^*(\mathbf{M})$ , where  $p^*$  is the polynomial from [Theorem 3.1](#) with  $d = 2q + 1$ . Then, we have

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$$\|\text{polar}(\mathbf{M}) - \mathbf{X}_T\|_2 \leq |1 - \ell^2|^{(q+1)^T}. \quad (12)$$

Hence, for  $d = 3$  and  $d = 5$  the method converges quadratically and cubically, respectively.

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In fact, our method is strictly faster than Newton-Schulz, even if  $\sigma_{\min}(\mathbf{M}) < \ell$ . When  $\sigma_{\min} = \ell$ , Polar Express is about twice as fast as Newton-Schulz (cf. [Chen & Chow \(2014, Section 3.1\)](#)). [Recent work has analyzed the stability and convergence of Muon when the polar factor is computed inexactly \(Shulgin et al., 2025; Refael et al., 2025\)](#). Combining these analyses with [Theorem 3.3](#) immediately yields a convergence guarantee for Muon as implemented with Polar Express.

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### 3.2 FINDING THE OPTIMAL POLYNOMIAL FOR EACH ITERATION

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[Theorem 3.1](#) shows that we can solve (6) by greedily choosing the optimal approximation  $p_t \in \mathbb{P}_d^{\text{odd}}$  for each interval  $[\ell_t, u_t]$  for  $t = 1, \dots, T$ . In this section, we show how to find each  $p_t$ . Since we are now focused on just one iteration, we drop the subscripts. Given  $\ell$  and  $u$ , we wish to solve the following optimization problem:

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$$\arg \min_{p \in \mathbb{P}_d^{\text{odd}}} \max_{x \in [\ell, u]} |1 - p(x)| \quad (13)$$

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That is, we seek a minimax or uniform approximation of the function  $x \mapsto 1$  on  $[\ell, u]$  from the set of odd polynomials. (Equivalently, we seek a minimax optimal approximation to  $\text{sign}(x)$  on  $[-u, -\ell] \cup [\ell, u]$ .) Problems of this form are well-studied in approximation theory and numerical analysis. The key mathematical insight underlying their solution is the Equioscillation Theorem, which we state formally for our setting in [Lemma C.1](#). This theorem is the basis of the Remez algorithm ([Pachón & Trefethen, 2009; Parks & McClellan, 1972](#)), a general-purpose method that finds a (nearly) optimal polynomial approximation of a given degree to *any* function on any interval. With a very minor modification to handle the constraint that  $p$  be odd, Remez can solve (13).

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However, the Remez algorithm is complicated and notoriously difficult to implement correctly.<sup>4</sup> Fortunately, we do not need the algorithm in its full generality; we seek only low-degree polynomial approximations, and the function we wish to approximate is just  $f(x) = 1$ . We use the Equioscillation Theorem to derive (17), an explicit, closed-form solution to (13) for the degree  $d = 3$  case. Up

<sup>4</sup>For implementations of the general Remez algorithm, we recommend [Chebfun](#) or [lolremez](#).

324 to rescaling, this turns out to be the same polynomial derived by different means in [Chen & Chow \(2014\)](#). For  $d = 5$ , we present [Algorithm 2](#), a simpler way of solving (13) that is mathematically  
 325 equivalent to Remez in our setting. This algorithm is implemented in its entirety in [Implementation 2](#). For more details, we refer the reader to [Appendix F](#).  
 326  
 327

### 329 3.3 UPPER AND LOWER BOUNDS ON THE SINGULAR VALUES 330

331 To instantiate our method, we need upper and lower bounds  $u$  and  $\ell$  on the singular values of the  
 332 input matrix  $M$ . A trivial upper bound is given by  $\|M\|_F$ . This can be quite loose in the worst case.  
 333 In practice, it is off only by a small constant factor because the gradient matrices of the weights  
 334 of dense linear layers in neural networks tend to have small effective rank ([Yang et al., 2024](#)). We  
 335 therefore rescale  $M$  by  $\|M\|_F$  and set  $u = 1$ . It is difficult to efficiently find a good lower bound on  
 336  $\sigma_{\min}$ , so we are forced to guess. Fortunately, the consequences of a bad guess are not severe. The  
 337 method converges for any  $\ell \in (0, u]$ , and even an order of magnitude error only delays convergence  
 338 by a few iterations. For matrices stored in floating point arithmetic, the singular values are usually  
 339 larger than machine precision  $\epsilon_{\text{mach}}$  ([Boutsikas et al., 2024](#)). We work in `bfloat16`, which has  
 340  $\epsilon_{\text{mach}} = 2^{-8} \approx 3.91 \cdot 10^{-3}$ , so we set  $\ell = 10^{-3}$ . Since we use these bounds for all input matrices,  
 341 we can pre-compute the optimal polynomials once and apply them to as many inputs as we want.  
 342  
 343

### 344 3.4 FINITE PRECISION CONSIDERATIONS 345

346 When working in finite-precision arithmetic, especially the half-precision `bfloat16` format used  
 347 in deep learning, we must take some care to avoid blowups and other problems due to numerical  
 348 error. To this end, we make a few small but crucial changes to the method in the offline stage  
 349 that stabilize it with a negligible effect on accuracy. One issue arises when numerical round-off  
 350 creates singular values that are slightly larger than our current upper bound  $u_t$ . To fix it, we replace  
 351 each polynomial  $p_t$  by  $x \mapsto p_t(x/1.01)$ , effectively increasing  $u_t$ . Another issue, identified by  
 352 [Nakatsukasa & Higham \(2013\)](#), is due to the non-monotonicity of  $p_t$ . We address it by using slightly  
 353 suboptimal (but less oscillatory) polynomials in the early iterations, as suggested by [Chen & Chow  
 354 \(2014\)](#). For a detailed discussion on the finite precision considerations, we refer to [Appendix G](#).  
 355

### 356 3.5 THE ALGORITHM 357

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#### 358 **Algorithm 1** The General Polar Express

359 **input:** Matrix  $M$ , iteration count  $T$ , degree  $d$ , approximate lower bound  $\ell$ .

360 **output:** An approximation  $X_T$  to  $\text{polar}(M)$ .

```

361 1 Offline: precompute polynomials in float64
362 2  $\ell_1 = \ell, u_1 = 1$ .
363 3 for  $t = 1, 2, \dots, T$  do
364 4   Solve using Remez (Appendix F):
365 5    $p_t = \arg \min_{p \in \mathbb{P}_d^{\text{odd}}} \max_{x \in [\max(\ell_t, u_t/10), u_t]} |1 - p(x)|$ 
366 6    $p_t \leftarrow p_t(\cdot/1.01)$ 
367 7    $\ell_{t+1} \leftarrow p_t(\ell_t), u_{t+1} \leftarrow 2 - \ell_{t+1}$ 
368 8 end for
369 9 Online: apply precomputed polynomials in bfloat16
370 10 Set  $X_0 = M / (\|M\|_F + 10^{-2})$ .
371 11 for  $t = 1, 2, \dots, T$  do
372 12    $X_t = p_t(X_{t-1})$ 
373 13 end for
374 14 return  $X_T$ .
375
```

We give the pseudocode of our proposed method for any degree in [Algorithm 1](#). We give the specific Python code of the Polar Express with degree  $d = 5$  and  $\ell = 10^{-3}$  used in our GPT experiments in [Implementations 1](#) and [2](#) in [Appendix A](#). Both incorporate the finite precision considerations discussed in [Section 3.4](#). Our algorithm precomputes the polynomials  $p_1, \dots, p_T$  of [Theorem 3.1](#) in full precision using the results of [Section 3.2](#) (or the Remez algorithm for  $d > 5$ ). This stage is offline because the coefficients of the polynomials are only computed and stored once. For every subsequent call to the algorithm, these coefficients are reused and the offline stage is skipped. For instance, in [Implementation 1](#) these polynomials have been precomputed and stored in the variable `coeffs_list`.

376 The online stage can be performed in lower precision (`bfloat16`) for greater speed on a GPU.  
 377 Horner's rule can be used to carry out each iteration. For instance, if  $p_t = ax + bx^3 + cx^5$ , then

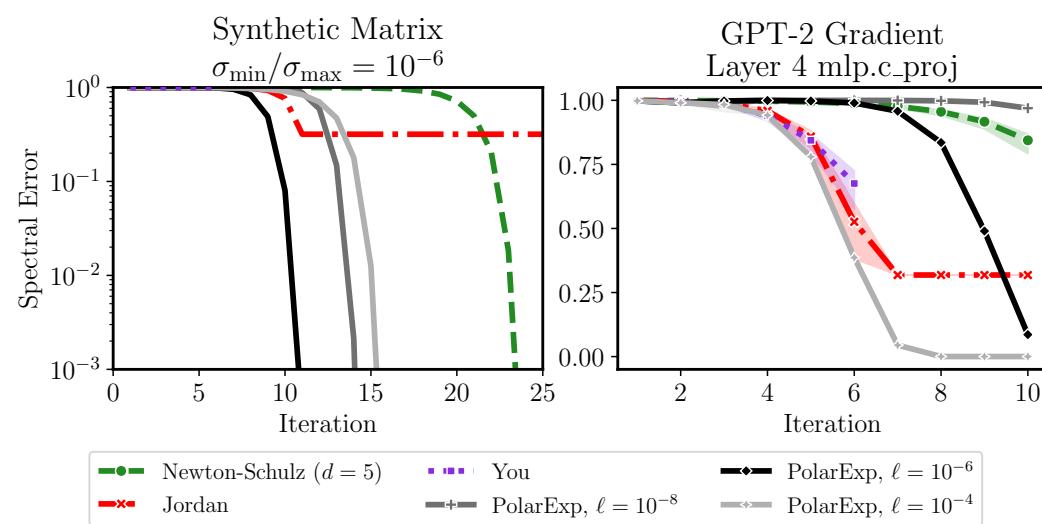


Figure 3: Convergence of degree-5 polynomial methods. Polar Express outperforms other methods at every iteration when tuned properly. Left panel: synthetic matrix with  $\sigma_{\max} = 1$ ,  $\sigma_{\min} = 10^{-6}$ . Right panel: gradient from randomly-initialized GPT-2 model on a batch of language modeling data. Shaded region shows 90% interval over 512 batches of data.

$\mathbf{X}_t = \mathbf{X}_{t-1} (a\mathbf{I} + \mathbf{Y}_{t-1} (b\mathbf{I} + c\mathbf{Y}_{t-1}))$  where  $\mathbf{Y}_{t-1} = \mathbf{X}_{t-1}^\top \mathbf{X}_{t-1}$ . A simple implementation of the offline stage of [Algorithm 1](#) is given in [Implementation 2](#). For deep learning applications, we recommend using  $d = 5$  and  $T = 5$  or  $6$  with  $\ell_1 = 10^{-3}$ . With these parameters, the offline stage as implemented in [Implementation 2](#) gives the polynomials encoded in `coeffs_list` in [Implementation 1](#). All told, our proposal for Muon is to apply the composition of these polynomials to  $\mathbf{M}/(\|\mathbf{M}\|_F + 10^{-2})$ .

## 4 NUMERICAL EXPERIMENTS

### 4.1 CONVERGENCE OF POLAR EXPRESS

We compare Polar Express against degree-5 Newton-Schulz and the methods of [Jordan et al. \(2024b\)](#) and [Cesista et al. \(2025\)](#). We first generate a random matrix whose singular values are evenly spaced on a logarithmic scale between  $10^{-6}$  and  $1$ , with singular vectors chosen randomly. The left panel of [Figure 3](#) shows the results. Since all the methods in this plot use degree-5 polynomials, their computational and runtime costs are all proportional to the number of iterations. As expected, Newton-Schulz converges but makes almost no progress for the first 17 iterations. Jordan’s method rapidly achieves an error of  $\approx 0.3$  after just 11 iterations, but ceases to converge further. You’s method, which is only defined for six iterations, converges at a similar rate as Jordan’s method. When Polar Express is instantiated with  $\ell = \sigma_{\min}$ , it dominates the other methods at every iteration, achieving excellent accuracy after just 11 iterations and converging about twice as fast as Newton-Schulz to any given error. Even when  $\ell$  is wrong by two orders of magnitude in either direction, the method remains competitive, though it does not outperform Jordan’s method until iteration 13 or 14. We also test convergence on a non-synthetic matrix: the gradient of a weight matrix from the fourth transformer block of a GPT-2 model ([Figure 3](#), right). Again, the best-tuned version of Polar Express outperforms the other methods, but setting  $\ell$  to be many orders of magnitude too small can delay convergence. Note that [Figure 3](#) measures error in the spectral norm. For many applications we may be satisfied with a looser measure of error, like the Frobenius norm.

### 4.2 TRAINING GPT-2

We compare the performance of using Polar Express ([Implementation 1](#)) inside Muon against Jordan’s ([Jordan et al., 2024b](#)) and You’s ([Cesista et al., 2025](#)) methods. We train two architectures:

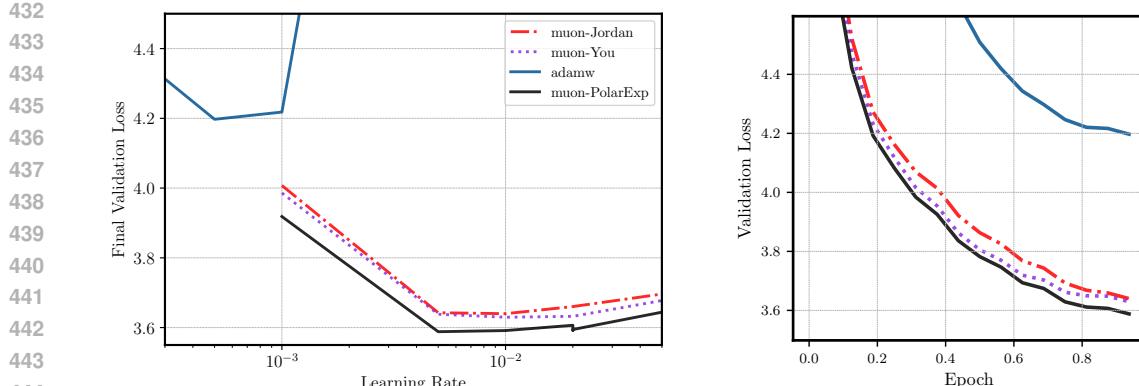


Figure 4: Training a GPT-2-Small (124M) model on 1 Billion tokens of the [FineWeb](#) data set (Aroca-Ouellette et al., 2023). muon-<method> denotes Muon with 5 iterations of <method> to compute  $\text{polar}(\mathbf{M})$ . No weight decay is used. Left: final validation loss vs. learning rate. The best final validation losses for each method were adamw( $\text{lr} = 0.0005$ ): 4.197, muon-Jordan( $\text{lr} = 0.01$ ): 3.639, muon-You( $\text{lr} = 0.01$ ): 3.629 and muon-PolarExp( $\text{lr} = 0.005$ ): 3.588. Right: Validation loss vs. training iteration.

GPT-2-Small ( $n_{\text{embd}} = 768$ ,  $n_{\text{layer}} = 12$ ,  $n_{\text{head}} = 12$ ) and GPT-2-Large ( $n_{\text{embd}} = 1280$ ,  $n_{\text{layer}} = 36$ ,  $n_{\text{head}} = 20$ ), both with a vocabulary size of 50,257 and a context length of 1024. We train on 1B tokens of the [FineWeb](#) dataset (Aroca-Ouellette et al., 2023) for one epoch with batch size 32. All runs use mixed precision (bf16) on 4 H100 GPUs with the learning rate schedule proposed in Jordan et al. (2024a)—a constant phase for the first 40% of training steps followed by linear decay. All methods for the matrix sign computations are performed in bf16 precision and use five iterations. Following nano-gpt (Jordan et al., 2024a), we assign Muon to all parameters with at least two dimensions (e.g., excluding RMS norm parameters), except for embeddings, unembeddings, and positional encodings. These excluded parameters are optimized with AdamW.

Figures 1 and 4 show the resulting in terms of validation loss for the GPT-Large and GPT-Small models, respectively. In both cases, muon-PolarExp achieves a better validation loss than muon-Jordan or muon-You. The advantage is remarkably consistent across all learning rates and epochs. While not shown in Figures 1 and 4, muon-PolarExp also achieves a better training loss than the baselines, and the improvements in training loss are nearly identical to the improvements in validation loss. Furthermore, since all three of these matrix sign methods are equally expensive (they all apply a degree 5 polynomial at each iteration), improved validation loss in terms of training steps also implies improved loss in terms of wall clock time. For figures displaying the improvements in training loss and wall-clock time, see Appendix H.2, Figure 11.

### 4.3 ABLATIONS

**Accuracy of polar approximation** We now explore how the accuracy of approximating  $\text{polar}(\mathbf{M})$  affects the optimization quality of Muon. Our main experiments with GPT-2 use 5 iterations. We trained GPT-2 Small with Muon using between 2 and 30 iterations of Polar Express instead. For comparison, we also implemented Muon with the *exact* polar factor, computed using `torch.linalg.svd`. Figure 5 shows the results. The left plot shows that when using only 2 or 3 iterations of Polar Express, the final validation loss is worse than when using 5 or 6 iterations. However, increasing the accuracy of the polar approximation further—even computing it exactly with the SVD—does not improve the optimization quality. The right plot shows that changing the number of iterations does not meaningfully change the runtime of Muon; in our setting, the runtime of computing  $\text{polar}(\mathbf{M})$  is dominated by the forward and backward passes. However, the SVD is so costly that using it *doubles* the runtime of each training step. These results validate the standard way of implementing Muon: using 5 or 6 iterations of an iterative approximation like Polar Express

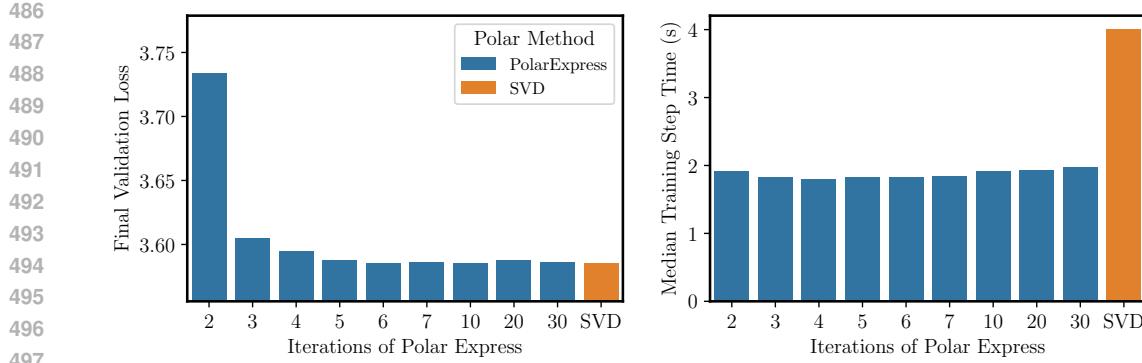


Figure 5: **Ablating the number of iterations of Polar Express used to implement Muon, and comparing to computing polar( $M$ ) exactly via an SVD.** Left: using  $> 6$  iterations or the SVD does not improve final validation loss. Right: Runtime of Muon is not sensitive to the number of iterations of Polar Express, but the SVD makes it significantly slower. All runs use GPT-2-Small with 1 Billion tokens of FineWeb data, learning rate 0.05, and weight decay 0.1.

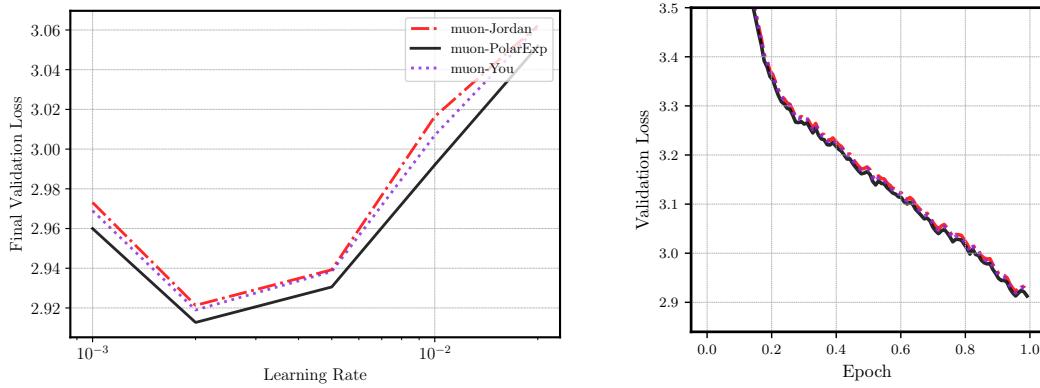


Figure 6: **Training GPT-2-Large on 10 billion tokens of FineWeb with weight decay 0.1.** Best final validation losses were muon-Jordan ( $lr = 0.002$ ): 2.921, muon-You ( $lr = 0.002$ ): 2.919 and muon-PolarExp ( $lr = 0.002$ ): 2.913.

rather than computing polar( $M$ ) exactly. For further experiments supporting this conclusion, see Appendix H.1, Figure 9.

**Weight decay** We also experimented with adding weight decay of 0.1 to the GPT-2 training runs, keeping all else the same. The results are presented in Appendix H.2, Figure 12. They are quite similar to Figures 1 and 4. We again find that muon-PolarExp outperforms the other methods.

**Number of Training Tokens** Our main experiments with GPT-2 use 1 billion tokens of training data from FineWeb (Aroca-Ouellette et al., 2023). We now select a subset of our training runs and extend them to 10 billion tokens. 10 billion tokens roughly matches the Chinchilla scaling rule for GPT-2-Large (774M params) and exceeds it for GPT-2-Small, as per Table 3 in Hoffmann et al. (2022). Figure 6 shows the results for GPT-2-Large with weight decay. (For GPT-2-Small, see Appendix H.2, Figure 13b). Polar Express still outperforms the baselines by a small but consistent margin.

540     **Reproducibility statement** A complete Pytorch implementation of our method is given in [Appendix A](#). Details of our experiments, including hyperparameters, are given in [Sections 4.1](#) and [4.2](#).  
 541     Source code to reproduce our experiments is given in the supplementary materials. Proofs of all  
 542     theoretical claims can be found in the appendices.  
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756  
757

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810 A CODE FOR POLAR EXPRESS  
811

812 **Implementation 1** gives a Python implementation of the online stage of [Algorithm 1](#) for degree  
813 = 5, which we use in our numerical experiments. It uses hard-coded polynomials generated from  
814 [Implementation 2](#) and incorporates a numerical safety factor of 1.01 as described in [Section 3.4](#).  
815 This implementation is designed for ease of use. It is short, it has no dependencies besides PyTorch,  
816 and it is a drop-in replacement for previous implementations of matrix sign methods ([Cesista et al.,](#)  
817 [2025](#); [Jordan et al., 2024b](#)), such as [Modula \(2024\)](#).

818 **Implementation 1** Python code for Polar Express of degree = 5.  
819

```
820 from itertools import repeat
821 import torch
822
823 coeffs_list = [
824     (8.28721201814563, -23.595886519098837, 17.300387312530933),
825     (4.107059111542203, -2.9478499167379106, 0.5448431082926601),
826     (3.9486908534822946, -2.908902115962949, 0.5518191394370137),
827     (3.3184196573706015, -2.488488024314874, 0.51004894012372),
828     (2.300652019954817, -1.6689039845747493, 0.4188073119525673),
829     (1.891301407787398, -1.2679958271945868, 0.37680408948524835),
830     (1.8750014808534479, -1.2500016453999487, 0.3750001645474248),
831     (1.875, -1.25, 0.375), # subsequent coeffs equal this numerically
832 ]
833 # safety factor for numerical stability (but exclude last polynomial)
834 coeffs_list = [(a / 1.01, b / 1.01**3, c / 1.01**5)
835             for (a, b, c) in coeffs_list[:-1]] + [coeffs_list[-1]]
836
837 @torch.compile
838 def PolarExpress(G: torch.Tensor, steps: int) -> torch.Tensor:
839     assert G.ndim >= 2
840     X = G.bfloat16() # for speed
841     if G.size(-2) > G.size(-1): X = X.mT # this reduces FLOPs
842     X = X / (X.norm(dim=(-2, -1), keepdim=True) * 1.01 + 1e-7)
843     hs = coeffs_list[:steps] + list(
844         repeat(coeffs_list[-1], steps - len(coeffs_list)))
845     for a, b, c in hs:
846         A = X @ X.mT
847         B = b * A + c * A @ A
848         X = a * X + B @ X # X <- aX + bX^3 + cX^5
849     if G.size(-2) > G.size(-1): X = X.mT
850
851 return X
```

846 **Implementation 2** gives a Python implementation of the offline stage of [Algorithm 1](#). This code was  
847 used to construct the coefficients of the polynomials given in [Implementation 1](#), which in turn were  
848 used in our Muon experiments ([Section 4.2](#)). It uses  $\ell = 10^{-3}$  and  $u = 1$  by default. It incorporates  
849 [Algorithm 2](#) and the finite precision modifications described in [Section 3.4](#).  
850

851 **Implementation 2** Polar Express, Offline Stage

```
852 from math import inf, sqrt
853 import numpy as np
854
855
856 def optimal_quintic(l, u):
857     assert 0 <= l <= u
858     if 1 - 5e-6 <= l / u:
859         # Above this threshold, the equioscillating polynomials
860         # is numerically equal to...
861         return (15/8)/u, (-10/8)/(u**3), (3/8)/(u**5)
862     # This initialization becomes exact as l -> u
863     q = (3*l + 1) / 4
864     r = (l + 3) / 4
865     E, old_E = inf, None
866     while not old_E or abs(old_E - E) > 1e-15:
```

```

864     old_E = E
865     LHS = np.array([
866         [1, 1**3, 1**5, 1],
867         [q, q**3, q**5, -1],
868         [r, r**3, r**5, 1],
869         [u, u**3, u**5, -1],
870     ])
871     a, b, c, E = np.linalg.solve(LHS, np.ones(4))
872     q, r = np.sqrt((-3*b + np.array([-1, 1]) *
873                     sqrt(9*b**2 - 20*a*c)) / (10*c))
874     return float(a), float(b), float(c)
875
876 def optimal_composition(l, num_iters, cushion=0.02407327424182761):
877     u = 1
878     coefficients = []
879     for _ in range(num_iters):
880         a, b, c = optimal_quintic(max(l, cushion*u), u)
881         # Due to cushioning, this may be centered around 1 with
882         # respect to 0.024*u, u. Recenter it around 1 with respect
883         # to l, u, meaning find c so that 1 - c*p(l) = c*p(u) - 1:
884         pl = a*l + b*l**3 + c*l**5
885         pu = a*u + b*u**3 + c*u**5
886         rescalar = 2/(pl + pu)
887         a *= rescalar; b *= rescalar; c *= rescalar
888         # Optionally incorporate safety factor here:
889         # a /= 1.01; b /= 1.01**3; c /= 1.01**5
890         coefficients.append((a, b, c))
891         l = a*l + b*l**3 + c*l**5
892         u = 2 - 1
893     return coefficients
894
895 print(*optimal_composition(1e-3, 10), sep="\n")
896
897

```

## B RELATED WORK

Computing  $\text{polar}(\mathbf{M})$  is an important and longstanding problem in numerical linear algebra, with applications spanning electronic structure calculations, lattice quantum chromodynamics, orthogonal Procrustes analysis, parallel algorithms for computing the SVD, and beyond; see e.g. (Higham, 1986; Kaneko et al., 2013; Douglas Carroll & Arabie, 1998; Gower & Dijksterhuis, 2004; Neuberger, 1998; Szabo & Ostlund, 1996).

**Newton-Schulz and polynomial Padé methods.** The earliest methods in the literature are polynomial iterations like (2). Several nearly simultaneous papers introduced the family of polynomial Padé iterations, comprising Newton-Schulz and its higher-degree analogues (Kovářík, 1970; Björck & Bowie, 1971; Higham, 1986; Leipnik, 1971). These higher-degree methods are also sometimes called “Newton-Schulz”; when doing so, we will specify the degree for clarity. In these methods, each iteration refines the current approximation  $\mathbf{X}_t$  by applying a low-degree odd matrix polynomial, where any odd monomial  $x \mapsto x^{2q+1}$  is defined for rectangular matrices by the formula  $\mathbf{X}_t \mapsto \mathbf{X}_t (\mathbf{X}_t^\top \mathbf{X}_t)^q$ . Our Polar Express method also takes this form, though unlike Newton-Schulz, it changes the polynomial at each iteration.

The polynomials used in Padé methods are chosen to match the value and first few derivatives of  $\text{sign}(x)$  at the points  $x = \pm 1$ . For instance, the update rule of the third method in this family is defined by  $p(x) = \frac{1}{16} (35x - 35x^3 + 21x^5 - 5x^7)$ , which is the unique degree-7 polynomial satisfying  $p(\pm 1) = \pm 1$  and  $p'(\pm 1) = p''(\pm 1) = p'''(\pm 1) = 0$ . These methods converge so long as all singular values of  $\mathbf{X}_0$  lie in  $(0, 1]$ , a condition guaranteed by the initialization of (2). Furthermore, the order of convergence of the degree  $2q + 1$  method is  $q + 1$  (Björck & Bowie, 1971). In particular, the Newton-Schulz method ( $q = 1$ ) converges quadratically.

918 **Newton’s method and rational Padé.** In the numerical analysis literature, polynomial methods  
 919 were succeeded by rational iterations like Newton’s method (Higham, 1986), defined as follows<sup>5</sup>:  
 920

$$921 \quad \mathbf{X}_0 = \mathbf{M} \quad \mathbf{X}_{t+1} = \frac{1}{2} (\mathbf{X}_t + \mathbf{X}_t^{-\top}) \quad (14)$$

922 Newton’s method also converges quadratically. Like Newton-Schulz, it works because the rational  
 923 function  $r(x) = \frac{1}{2}(x + x^{-1})$  has a stable fixed point at 1; unlike for Newton-Schulz, this point is a  
 924 global attractor for the whole positive real line. At first glance, Newton’s method has nothing to do  
 925 with the Padé iterations discussed above. However, after a change of variables  $\mathbf{Y}_t = \mathbf{X}_t^{-1}$ , it can be  
 926 reinterpreted as  $\mathbf{Y}_{t+1} = 2\mathbf{Y}_t(\mathbf{I} + \mathbf{Y}_t^\top \mathbf{Y}_t)^{-1}$ , which is sometimes called inverse Newton. Observing  
 927 that  $r(x) = \frac{2x}{1+x^2}$  satisfies  $r(\pm 1) = \pm 1$  and  $r'(\pm 1) = 0$ , we see that (inverse) Newton is also a  
 928 Padé method, though a rational rather than polynomial one. In fact, given a odd degree  $2q_n + 1$   
 929 for the numerator and an even degree  $2q_d$  for the denominator, there is a unique rational function  
 930 that matches the value and first  $q_n + q_d$  derivatives of  $\text{sign}(x)$  at  $x = \pm 1$ . This directly yields a  
 931 Padé method for computing polar( $\mathbf{M}$ ) whose order of convergence is  $q_n + q_d + 1$ . For instance,  
 932  $r(x) = \frac{3x+x^3}{1+3x^2}$  is called Halley’s method, which converges cubically. When  $q_d = 0$ , we recover the  
 933 polynomial Padé methods.

934 There are two main weakness of Newton’s method and the Padé iterations: slow convergence in the  
 935 initial phase and the need to compute explicit inverses. To accelerate initial convergence, Higham  
 936 popularized the technique of rescaling the matrix after every Newton iteration (Higham, 1986).  
 937 Intuitively, rescaling  $\mathbf{X}_t$  so that  $\sigma_{\max} = 1/\sigma_{\min}$  centers the spectrum around 1, where convergence  
 938 is fastest. Several easily-computable choices of scaling factor exist to accomplish this approximately.  
 939 Note that this rescaling scheme would fail for Newton-Schulz, which likewise suffers from slow  
 940 initial convergence but which would diverge if  $\sigma_{\max} \gg 1$ .

941 Computing matrix inverses is difficult to parallelize and to implement stably in low precision arith-  
 942 metic. However, a trick was developed for stably computing many rational methods *without* explicit  
 943 inverses; QR decompositions can be used instead (Nakatsukasa et al., 2010; Zhang et al., 2007).  
 944 Applying this trick to Halley’s method and combining with a special rescaling scheme yields the  
 945 QDWH (QR-based dynamically weighted Halley) method, which converges in just six iterations for  
 946 any reasonably conditioned matrix (Nakatsukasa et al., 2010).

947 **Adaptive rational methods from optimal approximations.** A landmark 2016 paper introduced a  
 948 new paradigm to design iterative methods for computing polar( $\mathbf{M}$ ) (Nakatsukasa & Freund, 2016).  
 949 The main insight is as follows. Padé methods choose the update rule to be an approximation to  
 950  $\text{sign}(x)$  of a given degree that is optimally accurate in the neighborhood of  $x = 1$ . Instead, we should  
 951 choose the approximation to  $\text{sign}(x)$  that is optimal over an *interval*  $[\ell, 1] \subset \mathbb{R}_{\geq 0}$  that contains the  
 952 singular values. Moreover, after each step of the algorithm, the range of the singular values changes;  
 953 therefore, we adapt the update rule at each iteration to match the new interval. When the range of the  
 954 singular values is large, this approach ensures that the update rule shrinks it as quickly as possible.  
 955 As the algorithm proceeds and the interval shrinks to a small neighborhood of 1, the update rule  
 956 approaches that of a Padé method, maintaining the same high order of convergence as it has.

957 Within the class of odd rational functions whose numerators and denominators have degree  $2q + 1$   
 958 and  $2q$ , respectively, an explicit formula for this optimal approximation to  $\text{sign}(x)$  on any interval  
 959  $[\ell, 1]$  was found by Zolotarev. It was shown that these rationals have remarkable convergence prop-  
 960 erties for any  $q$  (Nakatsukasa & Freund, 2016). For  $q = 1$ , this optimal approximation coincides  
 961 exactly with the dynamically weighted Halley’s method (QDWH) referenced above. For even faster  
 962 convergence than QDWH, (Nakatsukasa & Freund, 2016) proposed the Zolo-pd method, which uses  
 963  $q = 17$ . Finally, these methods all admit the same QR-based implementation trick as QDWH.

964 **Adaptive polynomial methods.** In this paper, we adopt the paradigm of Zolo-pd (Nakatsukasa  
 965 & Freund, 2016) but with polynomials rather than rationals of degree  $(2q + 1, 2q)$ . This choice  
 966 avoids the need for QR factorizations, relying solely on GPU-friendly matrix-matrix multiplications  
 967 in low-precision arithmetic. While this class of methods has not been fully developed in the numer-  
 968 ical analysis literature, similar ideas have been rediscovered in different guises. In an unpublished

970 <sup>5</sup>Our description of Newton’s method and other rational methods assumes square non-singular  $\mathbf{M}$ . Non-  
 971 square problems can be reduced to the square case by an initial QR decomposition, but this is not an option for  
 972 purely polynomial methods like ours.

manuscript that predates Zolo-pd, [Chen & Chow \(2014\)](#) describe a rescaling strategy for Newton-Schulz. Though motivated differently, their method is equivalent to ours for degree-3 polynomials (unlike our work, they do not consider general odd degree). They also observe numerical instability that prevents the method from converging to all the way to machine precision. Using the insights of [Nakatsukasa & Higham \(2012\)](#), they propose a simple mitigation for this issue that we adopt in [Section 3.4](#). Our work gives the approach from [Nakatsukasa & Higham \(2012\)](#) a stronger theoretical foundation that connects to the paradigm of Zolo-pd. Concretely, we prove that choosing an optimal polynomial at each iteration leads to a composed polynomial that is *globally* optimal in the sense of [\(5\)](#).

Independently, a group of cryptographers developed a similar method for approximating the scalar function  $\text{sign}(x)$  in the context of homomorphic encryption schemes ([Lee et al., 2022](#)). Their focus is mainly on tuning the analogues in their setting of the polynomial degree and number of iterations, whereas we focus on demonstrating optimality and efficiently constructing the update polynomials for degree 3 and 5. In addition, we consider matrix-valued inputs in low-precision arithmetic—not scalars in exact arithmetic—and we demonstrate our method’s effectiveness within the Muon algorithm for training deep neural networks.

**Application within Muon.** The designers of Muon realized that, due to the extreme efficiency requirements and lax accuracy requirements of their setting, rational-based methods from the numerical analysis literature are inapplicable. However, polynomial-based iteration schemes can take full advantage of GPUs because they use only matrix-matrix products in half-precision arithmetic, not inverses or QR decompositions. The preference for speed over accuracy motivates methods that aim to quickly produce coarse approximations, even at the cost of asymptotic convergence. Examples include the proposals of Jordan ([Jordan et al., 2024b](#)) and You ([Cesista et al., 2025](#)), as discussed in [Section 1.2](#). Like [Chen & Chow \(2014\)](#), Jordan found that convergence in the initial phase can be accelerated by choosing update rules that have a large derivative near zero, so as to increase the small singular values as much as possible at each iteration. You furthermore chose to use different update rules at each iteration, allowing extra flexibility to tune the trade-off between speed and accuracy. Both used degree-5 polynomials that were found through gradient descent on heuristic objective functions. These proposals were previously compared to Newton-Schulz<sup>6</sup>, but never to [Nakatsukasa & Higham \(2012\)](#). We find that our method (which generalizes [Nakatsukasa & Higham \(2012\)](#)) outperforms them all.

Finally, we remark that concurrent work of Grishina, Smirnov, and Rakhaba also proposes an adaptive polynomial method that generalizes [Nakatsukasa & Higham \(2012\)](#) and applies it to accelerating Muon ([Grishina et al., 2025](#)). Like [Nakatsukasa & Higham \(2012\)](#), this work does not establish global optimality of the composed polynomial as we do in [Section 3](#) or address finite precision considerations.

## C PROOF OF THEOREM 3.1

The aim of this section is to prove [Theorem 3.1](#). We begin with a result that provides a few essential properties for the the polynomial solving [\(6\)](#) when  $T = 1$ . This result is known as Chebyshev’s theorem ([Chebyshev, 1947](#)) or the equioscillation theorem ([Trefethen, 2020](#), Chapter 10).

**Lemma C.1.** Let  $d = 2q + 1$  and  $u, \ell > 0$ . Consider the problem

$$\min_{p \in \mathbb{P}_d^{\text{odd}}} \max_{x \in [\ell, u]} |1 - p(x)|. \quad (15)$$

There exists a unique polynomial  $p^* \in \mathbb{P}_d^{\text{odd}}$  solving [\(15\)](#). Furthermore,  $p^*$  is the unique solution to the above problem if and only if there exist  $q + 2$  distinct points  $\{x_0, \dots, x_{q+1}\} \subset [\ell, u]$  such

<sup>6</sup>[Jordan et al. \(2024b\)](#) actually compares to  $2x - \frac{3}{2}x^3 + \frac{1}{2}x^5$ , whereas the true degree-5 Newton-Schulz polynomial is  $(15x - 10x^3 + 3x^5)/8$ . However, the difference in performance is negligible for the first few iterations.

1026

that

$$1 - p^*(x_i) = \eta(-1)^i \max_{x \in [\ell, u]} |1 - p^*(x)|, \quad \text{for } i = 0, \dots, q+1,$$

for  $\eta = 1$  or  $\eta = -1$ .

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*Proof.* A discussion can be found in [Eremenko & Yuditskii \(2007\)](#). Here we include a formal proof for completeness.

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By Chebyshev's Theorem ([Achieser, 1992](#); [Chebyshev, 1947](#); [Cheney, 1966](#)) it is sufficient to show that  $\mathbb{P}_d^{\text{odd}}$  satisfies the Haar condition: any non-zero  $p \in \mathbb{P}_d^{\text{odd}} = \text{span}\{x, \dots, x^3, \dots, x^{2q+1}\}$  can have at most  $q$  roots in  $[\ell, u]$ .

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Since  $\deg(p) = d = 2q+1$  we know that  $p$  can have at most  $2q+1$  roots in  $\mathbb{R}$ . However, since  $p(0) = 0$  and  $p(x) = -p(-x)$  we know that  $p$  has one root at zero, and the remaining roots come in symmetric pairs  $(x, -x)$ . Because of this,  $p$  can have at most  $q$  roots in the positive orthant, and thus it can have at most  $q$  roots in  $[\ell, u] \subset (0, \infty)$ . Hence,  $\mathbb{P}_d^{\text{odd}}$  satisfies the Haar condition, which yields the desired result.  $\square$

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The proof of [Theorem 3.1](#) will be by induction on  $T$ . We begin by establishing the base case,  $T = 1$ , which is handled by the following result.

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**Lemma C.2.** Let  $u, \ell > 0$  and define

$$p^* := \arg \min_{p \in \mathbb{P}_d^*} \max_{x \in [\ell, u]} |1 - p(x)|.$$

Then

$$p^*(\ell) = \min_{x \in [\ell, u]} p^*(x), \quad \max_{x \in [\ell, u]} p^*(x) = 2 - p^*(\ell), \quad \text{and} \quad \max_{x \in [\ell, u]} |1 - p^*(x)| = 1 - p^*(\ell).$$

1055

1056

*Proof.* Throughout the proof we assume  $d = 2q+1$ . We begin with proving

1057

1058

$$p^*(\ell) = \min_{x \in [\ell, u]} p^*(x).$$

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Consider the polynomial  $e(x) := 1 - p^*(x)$ . The proof will contain three steps. We first rule out the trivial case that  $p^* \neq 0$ , since  $p(x) = \frac{2}{\ell+u}x$  would then be a better approximation. Hence,  $p^*$  cannot be the zero polynomial.

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1063

Step 1:  $e(x)$  has exactly  $q$  stationary points inside the open interval  $(\ell, u)$ .

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Note that  $e(x)$  has at most  $2q$  stationary points in  $\mathbb{R}$ , since its derivative  $e'(x)$  is a polynomial of degree  $2q$ . Furthermore, since  $p^*$  is odd, we have that  $e'(x) = -p'(x)$  is even of degree  $2q$ , and thus can have at most  $q$  stationary points contained in  $(0, +\infty)$ . Hence, there can be at most  $q$  stationary points of  $e(x)$  inside the interval  $[\ell, u]$ .

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By [Lemma C.1](#) there are  $q+2$  points  $x_0, \dots, x_{q+1} \in [\ell, u]$  where  $e(x)$  is maximized or minimized in  $[\ell, u]$ . These points are either stationary points or they are endpoints of the interval  $[\ell, u]$ . Let  $n_{\text{ext}}$  be the number of stationary points and  $n_{\text{stat}}$  be the number of endpoints in the set  $\{x_0, \dots, x_{q+1}\}$ . Since a point can be both a stationary point and an endpoint we have  $q+2 \leq n_{\text{end}} + n_{\text{stat}}$ . However,  $n_{\text{end}} \leq 2$  and  $n_{\text{stat}} \leq q$ , which follows from the previous paragraph where we showed that there are at most  $q$  stationary points of  $e(x)$  in  $[\ell, u]$ . So  $n_{\text{end}} + n_{\text{stat}} \leq q+2$ , and consequently we must have  $n_{\text{end}} = 2$  and  $n_{\text{stat}} = q$ , as required.

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Step 2:  $x = \ell$  is a maximum of  $e(x)$  on the interval  $[\ell, u]$

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By [Lemma C.1](#) and the discussion from Step 1, we know that  $|e(x)|$  is maximized at  $q+2$  points inside  $[\ell, u]$  and  $q$  of these points are contained inside the open interval  $(\ell, u)$ . Hence,  $x = \ell$  must either be a maximum or a minimum of  $e(x)$ . We will show that  $x = \ell$  must be a maximum by contradiction.

1080 Suppose  $x = \ell$  was a minimum of  $e(x)$  on  $[\ell, u]$ . First note that  $p^*$  is trivially non-negative on  $[\ell, u]$ ,  
 1081 or else  $p(x) = 0$  would be a better polynomial. Hence, since  $p^*(0) = 0$  we must have  $p^{*'}(\delta) > 0$   
 1082 for some  $\delta \in [0, \ell]$ , or else the zero polynomial  $p(x) = 0$  would be a better approximation. Hence,  
 1083 for some  $\delta \in [0, \ell]$  we have  $e'(\delta) < 0$ .

1084 We must also have  $e'(\ell) \geq 0$  or else  $x = \ell$  is not a minimum of  $e(x)$ . Since  $e'(\delta) < 0$  for some  
 1085  $\delta \in [0, \ell]$  and  $e'(\ell) \geq 0$ , by the intermediate value theorem there exists a point  $x^* \in [0, \ell]$  such that  
 1086  $e'(x^*) = 0$ . However, by the discussion above we know that all stationary points of  $e$  are contained  
 1087 inside the open interval  $(\ell, u)$ . Hence,  $x = \ell$  cannot be a minimum of  $e(x)$  on  $[\ell, u]$ . However, by  
 1088 Step 1 we know that the endpoints of  $[\ell, u]$  must be either minima or maxima of  $e(x)$ . Hence,  $x = \ell$   
 1089 is a maximum of  $e(x)$  on  $[\ell, u]$ .

1090 Step 3: Obtaining the desired equalities

1091 Since  $e(x)$  has a maximum in  $[\ell, u]$  at  $x = \ell$ , we have  $p^*(\ell) = \min_{x \in [\ell, u]} p^*(x)$ . The other two equalities  
 1092 are immediate consequences of the equioscillation property of  $p^*$  [Lemma C.1](#) and that  $x = \ell$  is a  
 1093 minimum of  $p^*$  over the set  $[\ell, u]$ .  $\square$

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1099 With the above-mentioned result in hand, we are ready to prove [Theorem 3.1](#).

1100

1101 **Theorem 3.1.** Let  $d$  be odd and define  $\ell_1 = \ell$  and  $u_1 = u$ . For  $t = 1, \dots, T$  define

$$p_t = \arg \min_{p \in \mathbb{P}_d^{\text{odd}}} \max_{x \in [\ell_t, u_t]} |1 - p(x)|, \quad \ell_{t+1} = \min_{x \in [\ell_t, u_t]} p_t(x), \quad u_{t+1} = \max_{x \in [\ell_t, u_t]} p_t(x) \quad (8)$$

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1103 The resulting composition  $p^* := p_T \circ p_{T-1} \circ \dots \circ p_1$  is optimal and the error is given by:

$$\max_{x \in [\ell, u]} |1 - p^*(x)| = \min_{\substack{p = p_T \circ p_{T-1} \circ \dots \circ p_1 \\ p \in \mathbb{P}_d^{\text{odd}}}} \max_{x \in [\ell, u]} |1 - p(x)| = 1 - \ell_{T+1}. \quad (9)$$

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1105 Furthermore the new error, lower and upper bounds can be computed through

$$\ell_{t+1} = p_t(\ell_t), \quad u_{t+1} = 2 - \ell_{t+1}, \quad \text{and} \quad \max_{x \in [\ell_t, u_t]} |1 - p_t(x)| = 1 - \ell_{t+1}. \quad (10)$$

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1116 *Proof.* The proof of (10) is an immediate consequence of [Lemma C.2](#), since for each  $t = 1, \dots, T$ ,  
 1117  $p_t$  is the optimal approximation in  $\mathbb{P}_d^{\text{odd}}$  to  $x \mapsto 1$ .  
 1118

1119 We now proceed with the proof of (9), which will be by induction. The proof for  $T = 1$  is an  
 1120 immediate consequence of [Lemma C.2](#) and we also have  $p^*(\ell) = \ell_2$  by (10). Now suppose the  
 1121 result is true for all  $t \leq T - 1$ . For  $t = 1, \dots, T - 1$ , note that the image of  $p_t$  on  $[\ell_t, u_t]$  is exactly  
 1122  $[\ell_{t+1}, u_{t+1}]$  by i). Hence, if we define  $g(x) := p_{T-1} \circ \dots \circ p_1(x)$ , then the image of  $g$  on  $[\ell, u]$  is  
 1123  $[\ell_T, u_T]$ . Furthermore, by i) we also have  $g(\ell) = \ell_T$ . Pick any  $f$  such that  $f \neq g$  and

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1127  $f = \tilde{p}_{T-1} \circ \dots \circ \tilde{p}_1$ ,  
 1128 for some  $\tilde{p}_1, \dots, \tilde{p}_{T-1} \in \mathbb{P}_d^{\text{odd}}$ . Let the image of  $f$  on  $[\ell, u]$  be  $[a, b]$ . We will prove that  $\frac{a}{b} \leq \frac{\ell_T}{u_T}$  by  
 1129 contradiction.

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1134 Suppose  $\frac{a}{b} > \frac{\ell_T}{u_T}$ . Define  $c = \frac{2}{a+b}$ . Then, the image of the scaled function  $cf$  on  $[\ell, u]$  is  $[ca, cb]$  and  $cf$  satisfies

$$\max_{x \in [\ell, u]} |1 - cf(x)| = \max \{1 - ca, cb - 1\} = \frac{b - a}{a + b}.$$

1134 Recall by our inductive hypothesis, we have  $\max_{x \in [\ell, u]} |1 - g(x)| = 1 - \ell_T = u_T - 1$  where the second  
 1135 equality holds by (10). It follows that  
 1136

$$\begin{aligned} & \frac{a}{b} > \frac{\ell_T}{u_T} \\ \Leftrightarrow & \frac{a}{b} > \frac{\ell_T}{2 - \ell_T} \\ \Leftrightarrow & \ell_T < \frac{2a}{a + b} \\ \Leftrightarrow & 1 - \ell_T > \frac{b - a}{a + b} \\ \Leftrightarrow & \max_{x \in [\ell, u]} |1 - g(x)| > \max_{x \in [\ell, u]} |1 - cf(x)|, \end{aligned}$$

1147 which leads to a contradiction to our inductive hypothesis that  $g$  is optimal. Hence, we must have  
 1148  $\frac{a}{b} \leq \frac{\ell_T}{u_T}$ .  
 1149

1150 Consequently, using that  $\frac{a}{b} \leq \frac{\ell_T}{u_T}$ , we will show that for any  $\tilde{p}_T \in \mathbb{P}_d^{\text{odd}}$  and for any  $f = \tilde{p}_{T-1} \circ \dots \circ \tilde{p}_1$ ,  $\tilde{p}_T \circ f$  cannot be a better approximation than  $p_T \circ g$ . In particular, we have  
 1151

$$\begin{aligned} \max_{x \in [\ell, u]} |1 - \tilde{p}_T(f(x))| & \geq \min_{p \in \mathbb{P}_d^*} \max_{x \in [\ell, u]} |1 - p(f(x))| \\ & = \min_{p \in \mathbb{P}_d^*} \max_{x \in [a, b]} |1 - p(x)| \\ & = \min_{p \in \mathbb{P}_d^*} \max_{x \in [a/b, 1]} |1 - p(x)| \\ & \geq \min_{p \in \mathbb{P}_d^*} \max_{x \in [\ell_T/u_T, 1]} |1 - p(x)| \\ & = \min_{p \in \mathbb{P}_d^*} \max_{x \in [\ell_T, u_T]} |1 - p(x)| \\ & = \min_{p \in \mathbb{P}_d^*} \max_{x \in [\ell, u]} |1 - p(g(x))| \\ & = \max_{x \in [\ell_T, u_T]} |1 - p_T(g(x))| = 1 - p_T(\ell_T) = 1 - \ell_{T+1}, \end{aligned}$$

1167 where the second and third equality follow by changing variables  $y = x/b$  so that  
 1168

$$\min_{p \in \mathbb{P}_d^*} \max_{x \in [a, b]} |1 - p(x)| = \min_{p \in \mathbb{P}_d^*} \max_{y \in [a/b, 1]} |1 - p(by)| = \min_{p \in \mathbb{P}_d^*} \max_{y \in [a/b, 1]} |1 - p(y)|$$

1169 and this last equality follows because the space  $\mathbb{P}_d^*$  is invariant under input rescaling; that is, for any  
 1170  $b \neq 0$ , the map  $x \mapsto bx$  preserves the space  $\text{span}\{x, x^3, \dots, x^d\}$ . This concludes the proof.  $\square$   
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## D PROOF OF THEOREM 3.3

1177 In this section we provide the proof of the convergence guarantee stated in Theorem 3.3.  
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1179 **Theorem 3.3.** Let  $\mathbf{M}$  be a matrix normalized so that  $\sigma(\mathbf{M}) \subset [\ell, 1]$ . Let  $\mathbf{X}_T = p^*(\mathbf{M})$ , where  
 1180  $p^*$  is the polynomial from Theorem 3.1 with  $d = 2q + 1$ . Then, we have

$$\|\text{polar}(\mathbf{M}) - \mathbf{X}_T\|_2 \leq |1 - \ell^2|^{(q+1)^T}. \quad (12)$$

1183 Hence, for  $d = 3$  and  $d = 5$  the method converges quadratically and cubically, respectively.  
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1185 *Proof.* Define  
 1186

$$p^* = \arg \min_{\substack{p = p_T \circ p_{T-1} \circ \dots \circ p_1 \\ p_t \in \mathbb{P}_d^*}} \max_{x \in [\ell, u]} |1 - p(x)|.$$

1188 Then [Algorithm 1](#) returns  $\mathbf{X}_T = p^*(\mathbf{M})$ . Let  $h \in \mathbb{P}_q$  be the  $[q/0]$  Padé-approximant to  $(1-x)^{-1/2}$   
 1189 ([Kenney & Laub, 1991](#), Section 3) and define  $p(x) = xh(1-x^2) \in \mathbb{P}_d^{\text{odd}}$ . Define  $f = p \circ \dots \circ p$  as  
 1190 the composition of  $p$  with itself  $T$  times. Then, by [Theorem 3.1](#), ([Kenney & Laub, 1991](#), Theorem  
 1191 3.1), and  $f(x) \geq 0$  for  $x \geq 0$  we have

$$\begin{aligned} 1193 \quad \|\text{sign}(\mathbf{M}) - \mathbf{X}_T\|_2 &\leq \max_{x \in [\ell, 1]} |1 - p^*(x)| \\ 1194 \quad &\leq \max_{x \in [\ell, 1]} |1 - f(x)| \\ 1196 \quad &\leq \max_{x \in [\ell, 1]} \left[ \frac{|1 - x^2|^{(d+1)^T}}{1 + f(x)} \right] \\ 1197 \quad &\leq |1 - \ell^2|^{(d+1)^T}, \\ 1199 \quad & \end{aligned}$$

1200 as required.  $\square$

## E PROOF OF EQUIVALENCE BETWEEN (5) AND (6)

1205 In this section we provide a proof for the equivalence between (5) and (6). It is sufficient to show  
 1206 that for any fixed polynomial  $p$  we have

$$1208 \quad \varepsilon_1 := \max_{\substack{\mathbf{M} \in \mathbb{R}^{m \times n} \\ \sigma(\mathbf{M}) \subset [\ell, u]}} \|\text{polar}(\mathbf{M}) - p(\mathbf{M})\|_2 = \max_{x \in [\ell, u]} |1 - p(x)| := \varepsilon_2.$$

1210 For any fixed  $\mathbf{M}$ , by the unitary invariance of the spectral norm we immediately have

$$1212 \quad \|\text{polar}(\mathbf{M}) - p(\mathbf{M})\|_2 = \max_{\sigma_i \in \sigma(\mathbf{M})} |1 - p(\sigma_i)| \leq \max_{x \in [\ell, u]} |1 - p(x)|.$$

1214 Consequently,  $\varepsilon_1 \leq \varepsilon_2$ .

1215 Suppose that  $x^* \in [\ell, u]$  is chosen so that  $|1 - p(x^*)| = \max_{x \in [\ell, u]} |1 - p(x)|$ . Without loss of  
 1216 generality, assume  $m \geq n$ . Letting  $\mathbf{M} = x^* \mathbf{U} \mathbf{V}^\top$ , for any matrix  $\mathbf{U} \in \mathbb{R}^{m \times n}$  and  $\mathbf{V} \in \mathbb{R}^{n \times n}$  with  
 1217 orthonormal columns, and noting  $\text{polar}(\mathbf{M}) = \mathbf{U} \mathbf{V}^\top$  yields

$$\begin{aligned} 1219 \quad \varepsilon_1 &\geq \|\text{polar}(\mathbf{M}) - p(\mathbf{M})\|_2 \\ 1220 \quad &= \|\mathbf{I}_n - p(x^*) \mathbf{I}_n\|_2 \\ 1221 \quad &= |1 - p(x^*)| \\ 1222 \quad &= \max_{x \in [\ell, u]} |1 - p(x)| = \varepsilon_2 \\ 1223 \quad & \end{aligned}$$

1224 Consequently,  $\varepsilon_1 \geq \varepsilon_2$ . Hence,  $\varepsilon_1 = \varepsilon_2$ , as desired.

## F REMEZ ALGORITHM

1229 In this section, we show in detail how to solve (13). By [Theorem 3.1](#), these solutions give the update  
 1230 rule for a single step of Polar Express. We give a closed form solution for  $d = 3$ . We then  
 1231 describe how the Remez algorithm ([Pachón & Trefethen, 2009](#); [Parks & McClellan, 1972](#)) can be  
 1232 used to approximate  $p_t$  for arbitrary  $d$ . We then present [Algorithm 2](#), a simplified version of Remez  
 1233 for solving (13) with  $d = 5$ . Recall (13):

$$1234 \quad \arg \min_{p \in \mathbb{P}_d^{\text{odd}}} \max_{x \in [\ell, u]} |1 - p(x)|$$

1237 We begin with the case when  $d = 3$ . We seek a polynomial of the form  $p(x) = ax + bx^3$ . The  
 1238 Equioscillation Theorem ([Lemma C.1](#)) stipulates that  $p$  must have an equioscillating set of size 3.  
 1239 For  $p$  to achieve its maximum error at a point  $x$ ,  $x$  must be a local extremum of  $p(x) - 1$  on the  
 1240 interval  $[\ell, u]$ . Thus, for  $x$  to be eligible for membership in the equioscillating set, it must either  
 1241 be a true local extremum of  $p(x) - 1$  that happens to lie in  $[\ell, u]$ , or else one of the endpoints  $\ell, u$ .  
 1242 However, because  $p$  is an odd cubic, it has at most one true local extremum on  $\mathbb{R}_{\geq 0}$ . Thus, to build

1242 an equioscillating set of three points, we must include  $p$ 's unique positive local extremum *and* both  
 1243 endpoints. This local extremum of  $p$  occurs at  $\sqrt{\frac{-a}{3b}}$ . Therefore, we seek  $a, b$  such that  
 1244

$$1245 \quad p(\ell) = 1 - E, \quad p\left(\sqrt{\frac{-a}{3b}}\right) = 1 + E, \quad p(u) = 1 - E \quad (16)$$

1246 for some  $E$ . This is a system of three equations in three variables. The solution  $p(x) = ax + bx^3$  is  
 1247 most easily expressed as follows. Let  $p_{\text{NS}}(x) = \frac{3}{2}x - \frac{1}{2}x^3$ . Then  
 1248

$$1249 \quad p(x) = \beta p_{\text{NS}}(\alpha x), \quad \text{where } \alpha = \sqrt{\frac{3}{u^2 + lu + \ell^2}} \quad \text{and} \quad \beta = \frac{4}{2 + \ell u (\ell + u) \alpha^3}. \quad (17)$$

1250 One can verify that this polynomial satisfies the equioscillation condition of (16), with  $\sqrt{\frac{-a}{3b}} = \frac{1}{\alpha}$   
 1251 and  $E = \beta - 1$ . Therefore, it must necessarily be the optimal approximation from  $\mathbb{P}_3^{\text{odd}}$ . Note that  
 1252 for  $u = 1$ ,  $x \mapsto p_{\text{NS}}(\alpha x)$  is the same polynomial derived in [Chen & Chow \(2014\)](#).

1253 Unfortunately, for larger  $d$ , finding closed form expressions for optimal approximations from  $\mathbb{P}_d^{\text{odd}}$   
 1254 becomes challenging, and we know of no closed form solution. However, we can approximate the  
 1255 optimal polynomial using the Remez algorithm. Let  $d = 2q + 1$ . Again recalling [Lemma C.1](#),  
 1256 the optimal polynomial must satisfy the equioscillation property at a set of  $q + 2$  points, as in (16).  
 1257 The Remez algorithm finds the equioscillation points  $A = \{x_0, \dots, x_{q+1}\}$  from [Lemma C.1](#) by  
 1258 iteratively refining a sequence of trial points  $A^{(k)} = \{x_0^{(k)}, \dots, x_{q+1}^{(k)}\}$  so that  $A^{(k)}$  converges to  $A$ .  
 1259 From the sequence of trial points  $A^{(k)}$  the algorithm also finds a sequence of polynomials  $p^{(k)}$  so  
 1260 that  $p^{(k)}$  converges to the optimal polynomial. The convergence is very fast, and usually 10 iterations  
 1261 is sufficient to converge to the optimal polynomial up to double precision machine epsilon ([Pachón & Trefethen, 2009](#)).  
 1262 More commonly, the Remez algorithm is used to find optimal polynomial  
 1263 approximations to general continuous functions where  $d \approx 100$  or even  $d \approx 1000$ . However,  
 1264 because the polynomial we build to approximate  $\text{sign}(x)$  is a composition of polynomials, each of  
 1265 which has a low degree, in our setting the degree  $d$  is small, usually  $d = 5$ . For  $d = 5$  the Remez  
 1266 algorithm simplifies significantly. We now describe this simplified algorithm.

1267 We first choose an initial set of trial points  $A^{(1)}$ , which ideally should come close to satisfying the  
 1268 equioscillation property. From [Lemma C.1](#), the unique optimal approximation  $p^* \in \mathbb{P}_5^{\text{odd}}$  satisfies  
 1269 the equioscillation property at four points in  $[\ell, u]$ . Since the function we wish to approximate is  
 1270 constant, the equioscillation points must be extrema of  $p^*$  on  $[\ell, u]$ . Because  $p^*$  is a odd quintic,  
 1271 it can have at most two local extrema on the positive real line, and thus at most two local extrema  
 1272 on  $[\ell, u]$ . The other two equioscillation points must therefore be the endpoints  $\ell$  and  $u$ . Since we  
 1273 know that  $\ell$  and  $u$  must be equioscillation points we always set  $x_0^{(k)} = \ell$  and  $x_3^{(k)} = u$  for all  $k$ .  
 1274 We initialize  $x_1^{(1)}$  and  $x_2^{(1)}$  to  $\frac{1}{4}\ell + \frac{3}{4}u$  and  $\frac{3}{4}\ell + \frac{1}{4}u$ , since we observe that as  $\ell \rightarrow u$  these are  
 1275 approximately the other two equioscillation points.

1276 We now show how to refine a candidate set of trial points  $A^{(k)}$  to produce  $A^{(k+1)}$  as well as an  
 1277 approximately equioscillating polynomial  $p_k$ . For any fixed set of trial points  $\{\ell, x_1^{(k)}, x_2^{(k)}, u\}$ , we  
 1278 can find a degree-5 odd polynomial  $p_k(x) = a_k x + b_k x^3 + c_k x^5$  that satisfies

$$1279 \quad p_k(\ell) = 1 - E_k, \quad p_k(x_1^{(k)}) = 1 + E_k, \quad p_k(x_2^{(k)}) = 1 - E_k, \quad p_k(u) = 1 + E_k \quad (18)$$

1280 for some  $E_k$  by solving a linear system in  $a_k, b_k, c_k$  and  $E_k$ . This can be rewritten as follows:

$$1281 \quad \begin{bmatrix} \ell & \ell^3 & \ell^5 & 1 \\ x_1^{(k)} & (x_1^{(k)})^3 & (x_1^{(k)})^5 & -1 \\ x_2^{(k)} & (x_2^{(k)})^3 & (x_2^{(k)})^5 & 1 \\ u & u^3 & u^5 & -1 \end{bmatrix} \begin{bmatrix} a_k \\ b_k \\ c_k \\ E_k \end{bmatrix} = \begin{bmatrix} 1 \\ 1 \\ 1 \\ 1 \end{bmatrix}. \quad (19)$$

1282 If  $A^{(k)}$  were the extrema of the error function  $e_k(x) = 1 - p_k(x)$  on  $[\ell, u]$ , then they would be  
 1283 an equioscillating set for  $p_k$ , and  $p_k$  would be the solution. Therefore, to refine  $A^{(k)}$ , we find the  
 1284 extrema of  $e_k(x) = 1 - p_k(x)$ . These can occur at  $\ell, u$  and the roots of  $e'_k(x)$ . Setting  $e'_k(x) = 0$

1296 yields the quartic equation  $5c_kx^4 + 3b_kx^2 + a_k = 0$ , whose two solutions are given explicitly by  
 1297 the quadratic formula after the substitution  $y = x^2$ . We set  $x_1^{(k+1)}$  and  $x_2^{(k+1)}$  to be the solutions  
 1298 to this equation and let  $A^{(k+1)} = \{\ell, x_1^{(k+1)}, x_2^{(k+1)}, u\}$ . We repeat the procedure until  $|E_k| :=$   
 1299  $\max_{x \in [\ell, u]} |1 - p_k(x)| \approx \max_{x \in [\ell, u]} |1 - p_{k+1}(x)| =: |E_{k+1}|$ .  
 1300

1301 We note that the matrix appearing in (19) is a Vandermonde matrix. Vandermonde matrices become  
 1302 notoriously ill-conditioned as the degree grows large (Golub & Van Loan, 2013, Section 4.6). How-  
 1303 ever, since in our setting we choose  $d$  to be small, there is no ill-conditioning due to large degrees.  
 1304 Instead, we observe ill-conditioning when  $\ell \approx u$ . However, as  $\ell/u \rightarrow 1$  the optimal polynomial will  
 1305 converge to the polynomial  $\frac{x/u}{8} (15 - 10(x/u)^2 + 3(x/u)^4)$ , which can be verified by noting that  
 1306 as  $\ell/u \rightarrow 1$  all equioscillation points  $x_0, x_1, x_2, x_3$  must converge to  $u$ . For general  $d = 2q + 1$ ,  
 1307 the polynomial will converge to  $(x/\ell)h(1 - (x/\ell)^2)$  where  $h \in \mathbb{P}_q$  is the  $[q/0]$  Padé approximant  
 1308 to  $(1 - x)^{1/2}$  (Kenney & Laub, 1991). In fact, this polynomial is extremely close to the optimal  
 1309 polynomial for sufficiently large  $\ell$ . To see this, let  $p^*$  be the optimal approximation from  $\mathbb{P}_5^{\text{odd}}$  and  
 1310 let  $p(x) = \frac{x/u}{8} (15 - 10(x/u)^2 + 3(x/u)^4)$ . Then,  
 1311

$$\begin{aligned} \max_{x \in [\ell, u]} |p^*(x) - p(x)| &\leq \max_{x \in [\ell, u]} |1 - p(x)| + \max_{x \in [\ell, u]} |1 - p^*(x)| \\ &\leq 2 \max_{x \in [\ell, u]} |1 - p(x)| \\ &\leq 2 (1 - \ell/u)^3. \end{aligned}$$

1312 where we invoked (Kenney & Laub, 1991, Theorem 3.1) and the fact that  $p^*$  is the optimal approx-  
 1313 imation to  $x \mapsto 1$  from  $\mathbb{P}_5^{\text{odd}}$ . Hence, when  $\ell/u \geq 1 - \epsilon_d^{1/3}$ , where  $\epsilon_{\text{double}} \approx 1.1 \times 10^{-16}$  is the  
 1314 double precision machine epsilon, then  $|p^*(x) - p(x)| \leq 2\epsilon_{\text{double}}$ . In other words, up to double  
 1315 precision machine epsilon,  $p^*$  is equal to  $p$ . Therefore, whenever  $\ell/u \geq 1 - \epsilon_{\text{double}}^{1/3}$  the algorithm  
 1316 simply returns the Padé approximant (that is, the scaled Newton-Schulz polynomial).  
 1317

1318 The full algorithm is given in [Algorithm 2](#). In our experiments, we never observed [Algorithm 2](#) tak-  
 1319 ing more than five iterations to converge. This algorithm is implemented in full in [Implementation 2](#).  
 1320

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1321 **Algorithm 2** Remez algorithm (degree 5 approximation for  $\text{sign}(x)$ )

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1322 **input:** interval  $[\ell, u]$  for  $u > \ell > 0$ .  
 1323 **output:** Approximation  $p \in \mathbb{P}_5^{\text{odd}}$  to  $p^* = \arg \min_{p \in \mathbb{P}_5^{\text{odd}}} \max_{x \in [\ell, u]} |1 - p(x)|$ .  
 1324

1325 **define**  $\epsilon_{\text{double}} = 1.11 \times 10^{-16}$   
 1326 **if**  $\ell/u \geq 1 - \epsilon_{\text{double}}^{1/3}$  **then**  
 1327     Return  $p(x) = \frac{x/u}{8} (15 - 10(x/u)^2 + 3(x/u)^4)$   
 1328 **end if**  
 1329      $x_1^{(1)} = \frac{1}{4}\ell + \frac{3}{4}u$ ,      $x_2^{(1)} = \frac{3}{4}\ell + \frac{1}{4}u$ .  
 1330      $E_0 = \infty$ ,      $E_{-1} = -\infty$   
 1331      $k \leftarrow 0$   
 1332     **while**  $||E_k| - |E_{k-1}|| > \epsilon_{\text{double}}$  **do**  
 1333          $k \leftarrow k + 1$   
 1334          $\begin{bmatrix} a_k \\ b_k \\ c_k \\ E_k \end{bmatrix} = \begin{bmatrix} \ell & \ell^3 & \ell^5 & 1 \\ x_1^{(k)} & (x_1^{(k)})^3 & (x_1^{(k)})^5 & -1 \\ x_2^{(k)} & (x_2^{(k)})^3 & (x_2^{(k)})^5 & 1 \\ u & u^3 & u^5 & -1 \end{bmatrix}^{-1} \begin{bmatrix} 1 \\ 1 \\ 1 \\ 1 \end{bmatrix}$   
 1335          $x_1^{(k+1)} = \sqrt{\frac{-3b_k - \sqrt{9b_k^2 - 20a_k c_k}}{10c_k}}$ ,      $x_2^{(k+1)} = \sqrt{\frac{-3b_k + \sqrt{9b_k^2 - 20a_k c_k}}{10c_k}}$   
 1336     **end while**  
 1337     Return  $p(x) = a_k x + b_k x^3 + c_k x^5$

---

1350 **G FINITE PRECISION CONSIDERATIONS**

1352 As highlighted in [Section 3.4](#), one must take care to implement Polar Express in finite precision.  
1353 In this section we outline modifications to our method to ensure stability in finite precision  
1354 arithmetic.

1355 The first issue arises when numerical round-off creates singular values that are slightly larger than  
1356 our current upper bound  $u_t$ . Our optimal polynomials converge only when the singular values of  $\mathbf{X}_t$   
1357 are less than  $u_t$ . In some cases we have

1359 
$$p_t(u_t + \epsilon) > u_{t+1} + \epsilon,$$

1360 so over many iterations, a singular value that is slightly larger than  $u_t$  large could grow to  $\infty$  instead  
1361 of converging to 1.

1362 To fix this issue, we simply replace each polynomial  $x \mapsto p_t(x)$  by  $x \mapsto p_t(x/1.01)$ . This safety  
1363 factor corrects for round-off errors in previous iterations while only slightly changing the behavior  
1364 of the polynomial on the interval  $[\ell_t, u_t]$ , though it does cause the singular values to converge to  
1365 0.999998 instead of to 1. To correct for this, the safety factor can be omitted in the final iteration.  
1366 This fix is reflected in line 5 of [Algorithm 1](#).

1367 The second issue was identified in [Nakatsukasa & Higham \(2012\)](#) and addressed in the context of  
1368 polynomial iterations by [Chen & Chow \(2014\)](#). In general, iterative methods for polar( $\mathbf{M}$ ) aim to  
1369 increase each singular value relative to the largest singular value; while  $\sigma_{\min}(\mathbf{X}_0) \ll \sigma_{\max}(\mathbf{X}_0)$ ,  
1370 after enough iterations,  $\sigma_{\min}(\mathbf{X}_t) \approx \sigma_{\max}(\mathbf{X}_t) \approx 1$ . However, the convergence of each singular  
1371 value to  $\sigma_{\max}$  may not be monotonic. Over the domain  $[\ell_t, u_t]$ , our optimal polynomial  $p_t$  oscillates  
1372 repeatedly between  $\ell_{t+1}$  and  $u_{t+1}$ , so some singular values that are near  $u_t$  may get mapped down  
1373 to  $\ell_{t+1}$ . It so happens that this non-monotonicity—even at a single iteration—can cause loss of  
1374 precision. That is, problems occur if

1376 
$$\frac{p_t(\sigma_i)}{\sigma_i} \ll \frac{\max_{x \in [\sigma_{\min}, \sigma_{\max}]} p_t(x)}{\sigma_{\max}},$$
1377

1378 where  $0 \leq \sigma_{\min} \leq \sigma_i \leq \sigma_{\max}$  are singular values of  $\mathbf{X}_t$  ([Nakatsukasa & Higham, 2012](#)). In the  
1379 extreme case  $p_t(\sigma_i) < 0$ , the  $i$ th singular vector will change sign, causing the method to converge  
1380 to the polar factor of the wrong matrix. Unlike Newton-Schulz, unscaled Newton, or QDWH, our  
1381 method is affected by this loss of precision.

1382 To mitigate this issue, [Chen & Chow \(2014\)](#) propose modifying their update polynomials to enforce  
1383 a lower bound on the ratio  $\frac{p_t(\sigma_i)}{\sigma_i}$ . This issue only occurs when  $\ell_t \ll u_t$ ; as  $\ell_t \rightarrow u_t$ , our optimal  
1384 polynomial approaches the Padé approximant and so  $\frac{p_t(x)}{x} \geq 1$  for all  $x \in [0, u_t]$ . We could fully  
1385 solve the problem by using the Padé approximant instead of our optimal polynomial, but this would  
1386 significantly slow down convergence. Instead we compromise. When  $\ell_t \geq u_t/10$ , we find that  
1387  $\frac{p_t(x)}{x} \geq 0.236$ . Therefore, whenever  $\ell_t < u_t/10$  we select the update rule as though  $\ell_t = u_t/10$ .  
1388 This change slows convergence, but only very slightly. (The choice of 10 is somewhat arbitrary. In  
1389 [Implementation 2](#), we use a different factor.) This fix is reflected in line 4 of [Algorithm 1](#).

1390 The third change is copied from the original Muon implementation: normalize  $\mathbf{M}$  by  $\|\mathbf{M}\|_{\text{F}} + 10^{-2}$   
1391 instead of by  $\|\mathbf{M}\|_{\text{F}}$ . As before, we set  $u_1 = 1$ . This fix is reflected in line 10 of [Algorithm 1](#).

1394 **H ADDITIONAL EXPERIMENTAL RESULTS**

1395 In this section, we present additional experimental results.

1396 **H.1 CONVERGENCE OF POLAR EXPRESS AND ITS IMPACT ON MUON**

1397 **Convergence in Frobenius Norm** In Figure 8, we plot the convergence of Polar Express and  
1398 three baselines as measured in the Frobenius norm. We also plot convergence in cosine similarity,

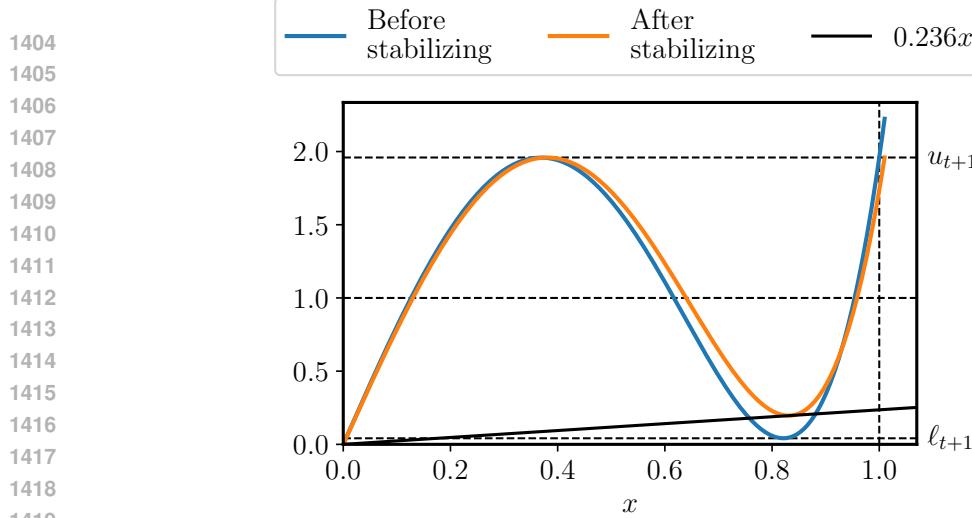


Figure 7: Effects of stabilizing the update rules with a safety factor and cushioning, as described in Appendix G. The blue curve is the optimal degree-5 polynomial for the interval  $[0.005, 1]$ . It has numerical issues because it maps singular values near 0.8 down to almost zero and maps  $1 + \epsilon$  to  $\approx u_{t+1} + 25\epsilon$ . The stabilized version is better because it ensures  $\frac{p_t(x)}{x} \geq 0.236$  and maps all  $x \leq 1.01$  to at most  $u_{t+1}$ .

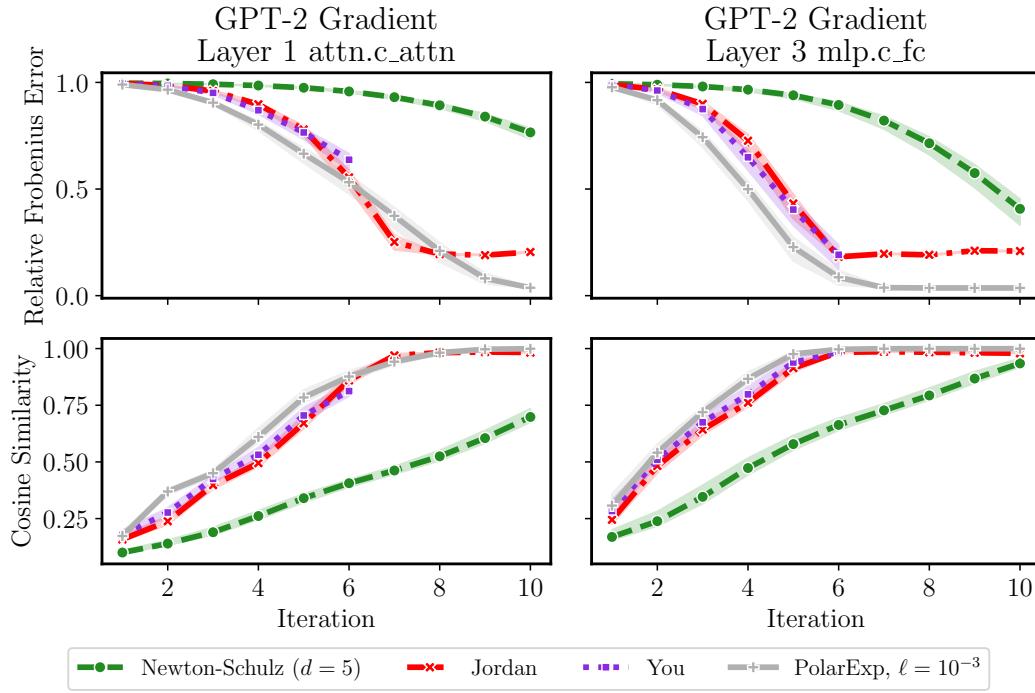


Figure 8: Convergence of degree-5 polynomial methods measured in Frobenius norm and cosine similarity. Test matrices are gradients of two layers of a randomly-initialized GPT-2 model on a batch of language modeling data. Polar Express outperforms other methods.

which is defined with respect to the Frobenius inner product  $\langle \mathbf{A}, \mathbf{B} \rangle = \text{Tr}(\mathbf{A}^\top \mathbf{B})$ . Formally, the cosine similarity between  $\mathbf{A}$  and  $\mathbf{B}$  is defined as  $\frac{\langle \mathbf{A}, \mathbf{B} \rangle}{\|\mathbf{A}\|_F \|\mathbf{B}\|_F}$ . We use gradients of GPT-2 layers as test matrices. While Polar Express is designed to minimize the spectral norm error, convergence in the Frobenius norm is similar (compare with Figure 3).

1458

1459

1460 **(In)sensitivity of Muon to Small Singular Values** Figure 5 shows that using more than five or six  
 1461 iterations of Polar Express does not improve the performance of Muon. However, Figures 3  
 1462 and 8 show that five iterations is not enough for Polar Express or any other method to converge.  
 1463 In practice, Polar Express is taking steps in directions that are meaningfully different from  
 1464 the exact polar( $M$ ) (as computed by an SVD), but still converging equally fast. One possible  
 1465 explanation for this observation is that Muon may not be sensitive to the convergence of small  
 1466 singular values of  $M$ . Intuitively, the singular vectors associated with these small singular values  
 1467 correspond to directions which have little effect on the output of the neural network; they may  
 1468 signify little more than noise in the stochastic gradients.

1469 We now conduct an experiment to test this hypothesis. We compare three ways that a Muon-like  
 1470 optimizer could handle the small singular values. Assume  $M$  has full rank, and partition the singular  
 1471 value decomposition of  $M$  into two parts

$$M = U\Sigma V^\top = [U_1 \ U_2] \begin{bmatrix} \Sigma_1 & \\ & \Sigma_2 \end{bmatrix} [V_1 \ V_2]^\top = U_1 \Sigma_1 V_1^\top + U_2 \Sigma_2 V_2^\top \quad (20)$$

1472 where  $\Sigma_1$  contains the singular values larger than some threshold  $\gamma\sigma_{\max}$  and  $\Sigma_2$  contains those  
 1473 smaller than  $\gamma\sigma_{\max}$ , where  $\sigma_{\max}$  is the largest singular value of  $M$ . Recall that

$$\text{polar}(M) := UV^\top = U_1 V_1^\top + U_2 V_2^\top \quad (21)$$

1474 is obtained by mapping each singular value of  $M$  to 1. We define the truncated polar factor by  
 1475 mapping the larger singular values to 1 and the smaller singular values to 0:

$$\text{polar}_\gamma(M) := U_1 V_1^\top. \quad (22)$$

1476 A third possibility is to map the small singular values to  $-1$ :

$$UV^\top = U_1 V_1^\top - U_2 V_2^\top \quad (23)$$

1477 Note that  $-U_2 V_2^\top$  is in the *opposite* direction as the Muon update. If the small singular values  
 1478 carry meaningful information about the loss landscape, then we expect this partly “uphill” step to  
 1479 hurt performance. Comparing the three update rules in Equations (21) to (23) can tell us how small  
 1480 singular values affect Muon.

1481 We train GPT-2 Small using each of these three update rules with learning rate 0.05 and weight  
 1482 decay 0.1. We sweep three different options for the cutoff  $\gamma$  that defines the ‘small’ singular values:  
 1483  $10^{-4}$ ,  $10^{-3}$ , and  $10^{-2}$ . The results are plotted in Figure 9. They show that the treatment of singular  
 1484 values smaller than  $10^{-4}\sigma_{\max}$  does not matter at all for the performance of Muon, and those smaller  
 1485 than  $10^{-3}\sigma_{\max}$  have a very minor effect. Notably, even *reversing* the direction of the Muon step  
 1486 in the bottom singular subspace barely worsens performance, showing that the gradient information  
 1487 in this subspace not very informative. The bottom panel of Figure 9 shows how five iterations of  
 1488 Polar Express (with  $\ell = 10^{-3}$ ) affect small singular values. Singular values greater than  $10^{-3}$   
 1489 are all mapped close to 1, while those smaller than  $10^{-4}$  are all mapped close to 0. Thus, while  
 1490 Polar Express does not fully converge after five iterations, it does converge in the ways that  
 1491 matter for Muon.

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1512 **Convergence of Top Singular Values** As discussed in the previous paragraph, we hypothesize  
 1513 that Muon may not be sensitive to the convergence of the small singular values of  $M$  when approx-  
 1514 imating polar( $M$ ). Therefore, in Figure 10, we plot the convergence of Polar Express and the  
 1515 baselines when all singular values smaller than  $10^{-3}$  are ignored. Specifically, if  $\text{alg}(M)$  denotes  
 1516 the output of an algorithm for approximating polar( $M$ ), then we compare

$$U_1 U_1^\top \cdot \text{alg}(M) \cdot V_1 V_1^\top \quad \text{to} \quad \text{polar}_{10^{-3}}(M),$$

1517 where  $\text{polar}_{10^{-3}}(M) = U_1 V_1^\top = U_1 U_1^\top \cdot \text{polar}(M) \cdot V_1 V_1^\top$  is the truncated polar factor defined  
 1518 above. The results show that Polar Express converges in just six iterations as measured in  
 1519 the relative Frobenius norm and just five iterations when measuring in cosine similarity. The other  
 1520 methods converge faster too, but Polar Express still outperforms them. These results may  
 1521 explain why the performance of Muon saturates at five or six iterations of Polar Express, as  
 1522 shown in Figure 5.

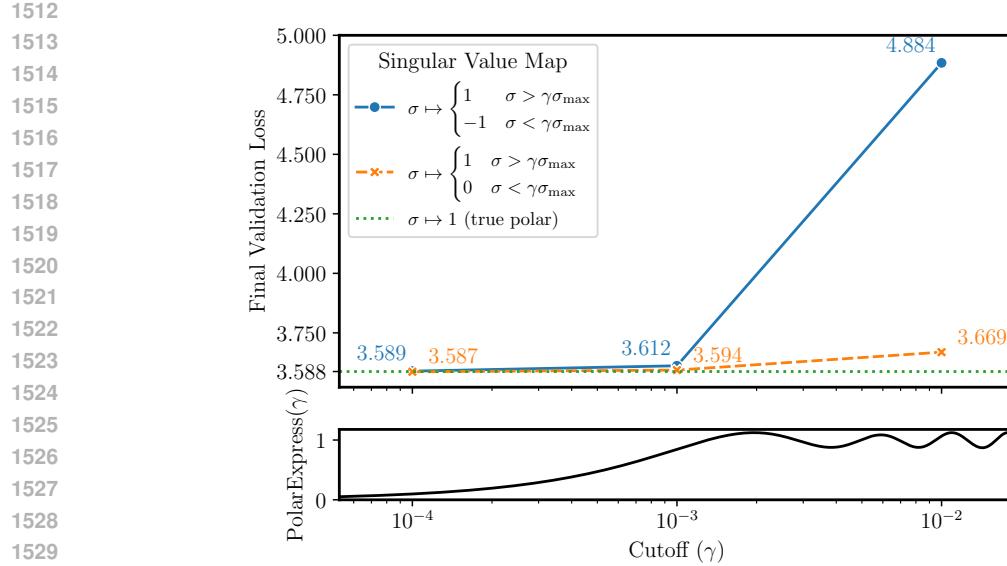


Figure 9: Impact of small singular directions of momentum matrix on optimization quality. We compare three variations of the Muon update rule. Exact Muon (green) processes the momentum  $\mathbf{M} = \mathbf{U}\Sigma\mathbf{V}^\top$  by mapping each singular value to 1:  $\text{polar}(\mathbf{M}) = \mathbf{U}\mathbf{V}^\top$ . Truncated Muon (orange) maps the larger singular values to 1 and the smaller singular values to 0. Reverse Muon (blue) maps the larger ones to 1 and the smaller ones to  $-1$ . Computations are performed in bfloat32. All runs train GPT-2 Small on 1 billion tokens of FineWeb data with learning rate 0.05 and weight decay 0.1. When the cutoff that defines ‘‘large’’ and ‘‘small’’ singular values is  $\gamma \approx 10^{-3}$ , all three methods perform well, showing that the small singular directions do not matter. Bottom panel shows the polynomial defined by composing five iterations of Polar Express. Five iterations is just enough for singular values  $\geq 10^{-3}$  to nearly converge.

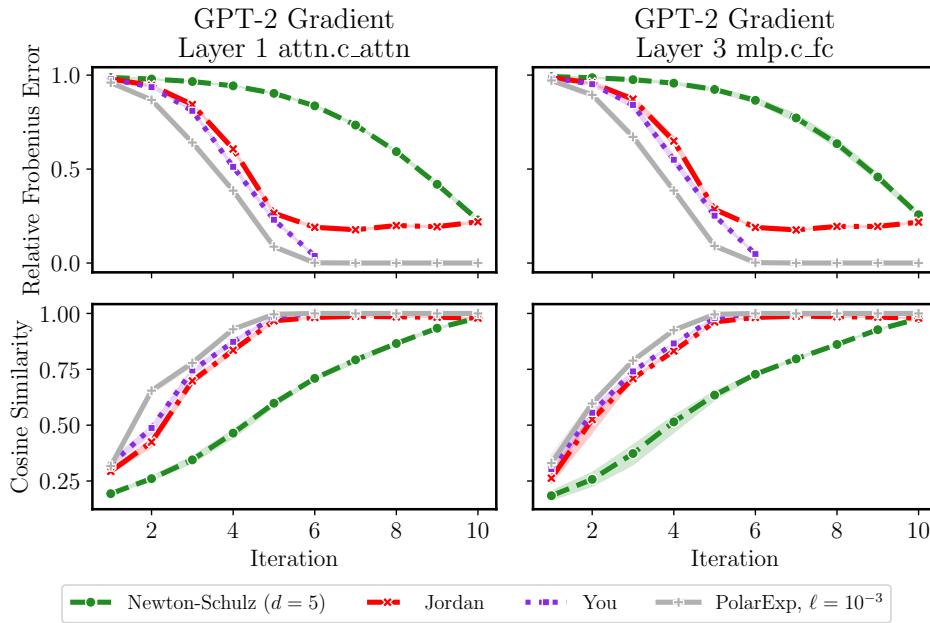
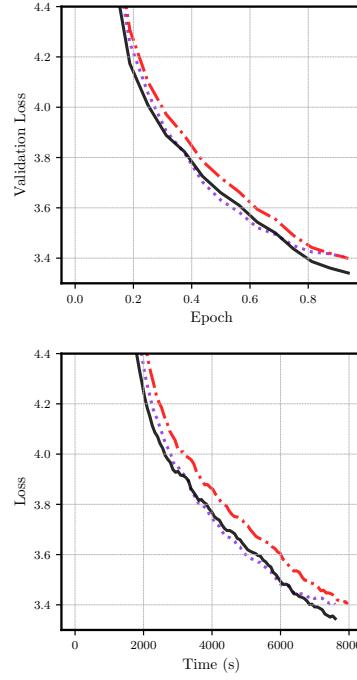
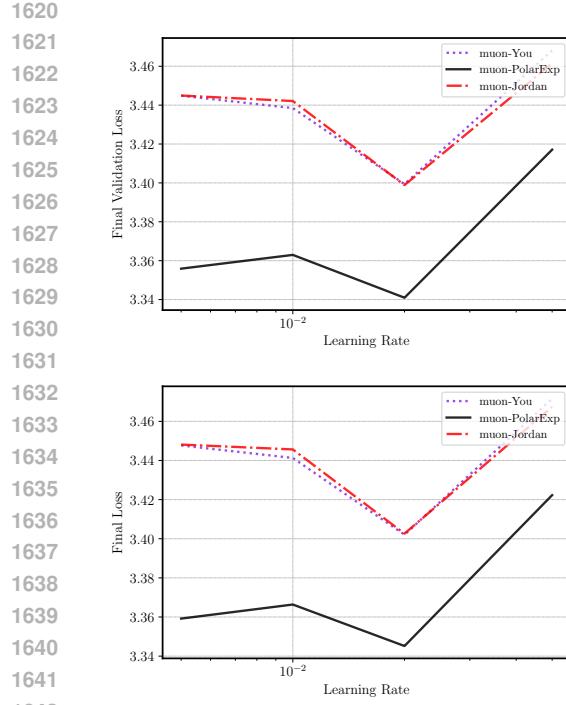
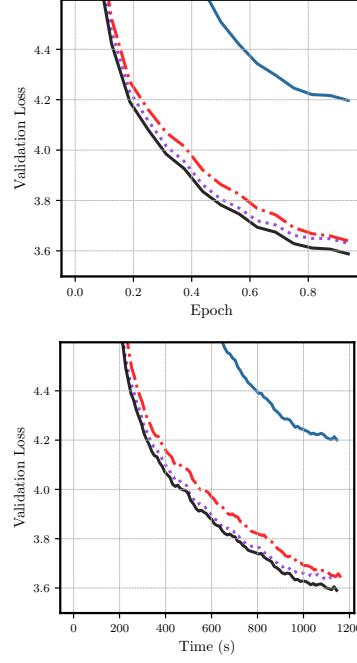
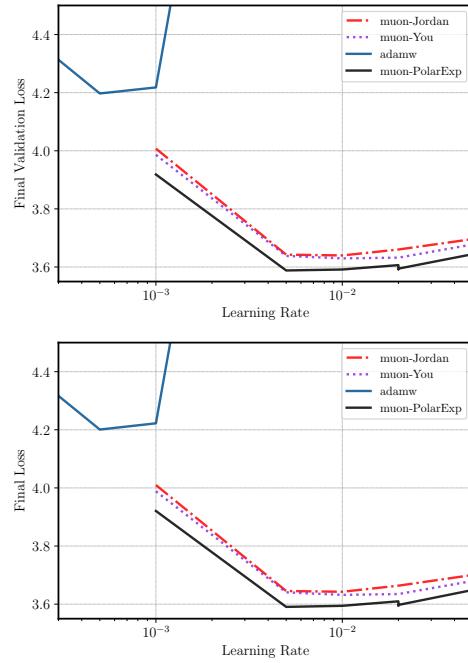


Figure 10: Convergence of degree-5 polynomial methods, considering only singular values larger than  $\sigma_{\max}/10^3$ . Test matrices are gradients of two layers of a randomly-initialized GPT-2 model on a batch of language modeling data. Polar Express converges in just five or six iterations and outperforms other methods.

1566 H.2 TRAINING GPT-2  
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15691570 **Additional Metrics** We report additional results from the experiment of Section 4.2. In addition to  
1571 showing validation loss vs. learning rate and training step, we also report *training* loss vs. learning  
1572 rate and training *time*. The results are shown in Figures 11a and 11b. The upper rows of each  
1573 subfigure are identical to Figure 1 and Figure 4, and are repeated here for ease of comparison.  
1574  
15751576 **Weight Decay** As described in Section 4.3, we reran our GPT-2 training runs with weight decay  
1577 of 0.1. This change had little effect on the results, as shown in Figure 12.  
1578  
15791580 **Number of Training Tokens** We also reran some of our GPT-2 training runs using 10 billion  
1581 tokens of training data instead of 1 billion. As described in Section 4.3, 10 billion tokens roughly  
1582 matches the Chinchilla scaling rule for GPT-2-Large and exceeds it for GPT-2-Small. Results are  
1583 shown in Figure 13. Note that the top row of Figure 13a is identical to Figure 6. *Polar Express*  
1584 still outperforms the baselines across all conditions, but the gap shrinks as the training loss con-  
1585 verges.  
15861587 H.3 IMAGE CLASSIFICATION  
15881589 We conducted experiments on the CIFAR-10 and CIFAR-100 image classification benchmarks  
1590 (Krizhevsky, 2009) using ResNet-20 and ResNet-110 architectures with batch normalization (He  
1591 et al., 2016). We used a range of learning rates in the range  $10^{-6}$  to 1 with a constant learning-rate  
1592 schedule, a batch size of 128, and 50 epochs of training data. We used three different random seeds  
1593 for each hyperparameter setting to assess stability and variability. As a baseline, we also included  
1594 AdamW and SGD with momentum (Kingma & Ba, 2015). Results are given in Figures 14 and 15.  
1595 For these experiments we see that all the Muon variants performed well, matching or exceeding  
1596 the training loss and validation accuracy of AdamW and sgd-m while also being more stable with  
1597 respect to the choice of learning rate. However, we do not see a marked difference between the  
1598 varieties of Muon. Indeed, even Newton-Schulz (degree = 5) performs equally well in this context,  
1599 despite being significantly less accurate than *Polar Express*, *Jordan* or *You*.  
16001601 Next we train a Vision Transformer (patch size 4, embedding dimension 512, depth 6, 8 heads, MLP  
1602 dimension 512, dropout 0.1) on CIFAR-10 for 200 epochs with batch size 512 using a constant  
1603 learning rate schedule. Results are shown in Figure 16. Muon with *Polar Express* achieved  
1604 the best training and validation loss (closely followed by *Jordan*'s and *You*'s methods). However,  
1605 improved loss did not entirely translate to better accuracy: both Muon and Newton-Schulz and  
1606 Adam performed well in terms of validation accuracy. Overall, these experiments do not show a  
1607 consistent advantage for *Polar Express*. Further work may be beneficial to fully realize the  
1608 potential benefits of Muon and to further tune *Polar Express* for these settings.  
16091610 I INITIALIZATION FOR MATRICES WITH LARGE SPECTRAL GAPS  
16111612 In Section 3, we constructed a sequence of polynomials that is adapted to the range of the singular  
1613 values  $[\ell, u]$ . Assuming nothing else about the input, these polynomials are optimal since they pro-  
1614 vide a good approximation to 1 across the entire interval. However, in many applications, the spec-  
1615 trum has large gaps; that is, there are several large outlying singular values that are well-separated  
1616 from the rest. For these matrices, it is not necessary for the polynomial to be accurate on the entire  
1617 interval  $[\ell, u]$ , only on the range of the small singular values plus a few other isolated points. In this  
1618 section, we take advantage of this structure to accelerate our method by preprocessing the matrix to  
1619 eliminate the largest singular values.1620 The first step is to find small intervals containing each of these large singular values. To find lower  
1621 bounds, we use subspace iteration, which is a generalization of the power method that approximates  
1622 multiple singular values simultaneously. Fix  $k$ , the number of singular values we wish to eliminate.  
1623

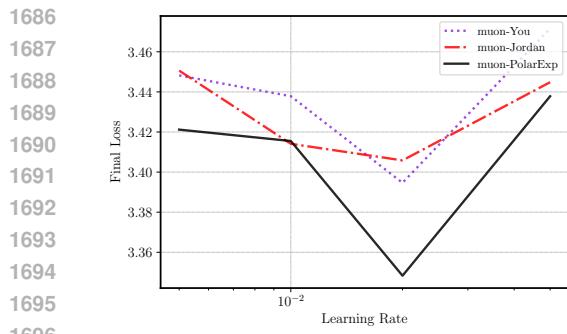
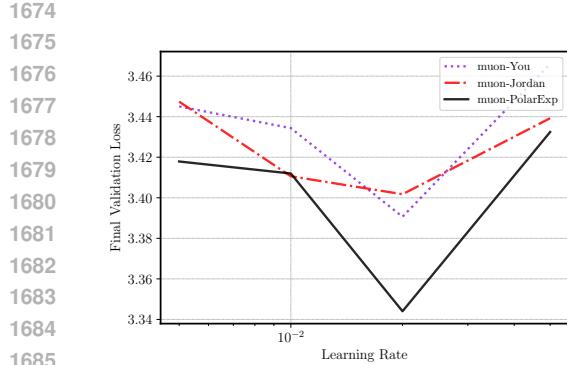


(a) GPT-2-Large (774M params). Best final validation losses were muon-You (lr = 0.02): 3.399, muon-Jordan (lr = 0.02): 3.398 and muon-PolarExp (lr = 0.02): 3.340.

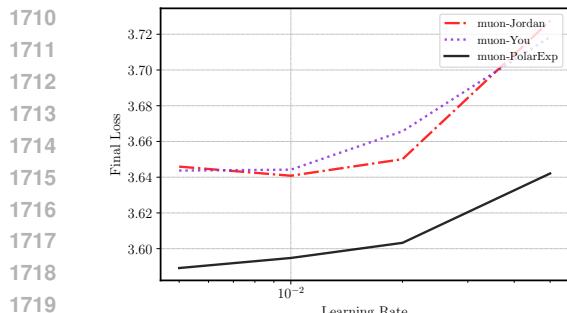
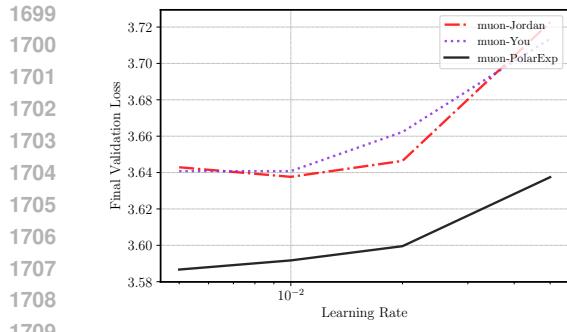


(b) GPT-2-Small (124M params). Best final validation losses were adamw (lr = 0.001): 4.197, muon-Jordan (lr = 0.01): 3.639, muon-You (lr = 0.01): 3.629 and muon-PolarExp (lr = 0.005): 3.588.

Figure 11: Training GPT-2 on 1 billion tokens of FineWeb data (Aroca-Ouellette et al., 2023) without weight decay. The label muon-<method> denotes Muon with 5 iterations of <method> to compute polar( $\mathcal{M}$ ). Top left: final validation loss vs. learning rate. Bottom left: final training loss vs. learning rate. Top right: validation loss vs. number of iterations for best learning rate. Bottom right: training loss vs. time for best learning rate.

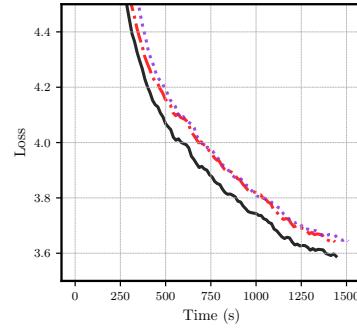
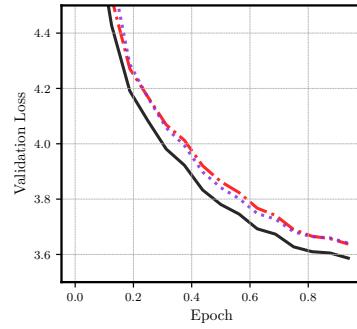
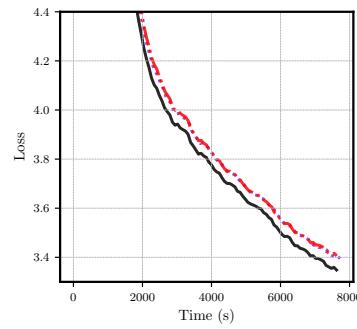
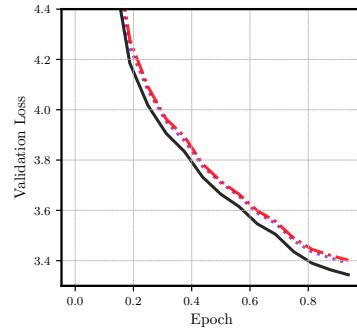


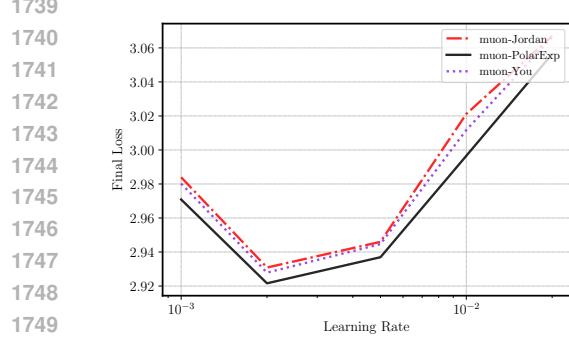
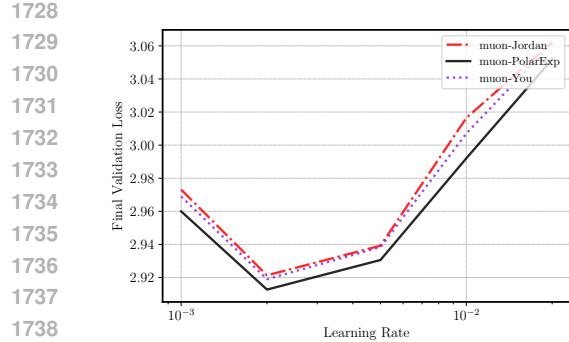
1697 (a) GPT-2-Large (774M params). Best final validation losses were muon-You (lr = 0.02): 3.390,  
1698 muon-Jordan (lr = 0.02): 3.401 and muon-PolarExp (lr = 0.02): 3.344.



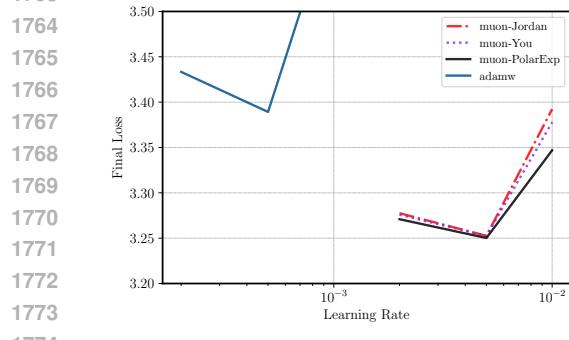
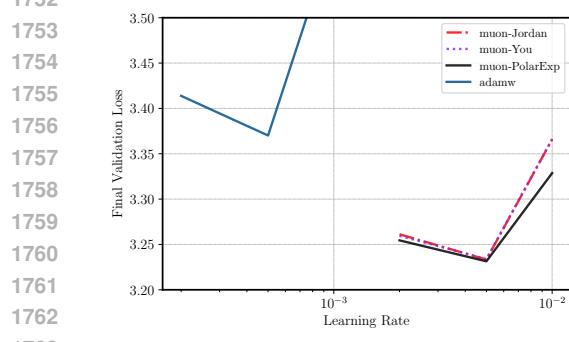
1721 (b) GPT-2-Small (124M params). Best final validation losses were muon-Jordan (lr = 0.01): 3.638,  
1722 muon-You (lr = 0.005): 3.641 and muon-PolarExp (lr = 0.005): 3.587.

1723 Figure 12: Training GPT-2 on 1 billion tokens of FineWeb data (Aroca-Ouellette et al., 2023) with  
1724 weight decay 0.1. The label muon-<method> denotes Muon with 5 iterations of <method> to  
1725 compute  $\text{polar}(\mathcal{M})$ . Top left: final validation loss vs. learning rate. Bottom left: final training loss  
1726 vs. learning rate. Top right: validation loss vs. number of iterations for best learning rate. Bottom  
1727 right: training loss vs. time for best learning rate.





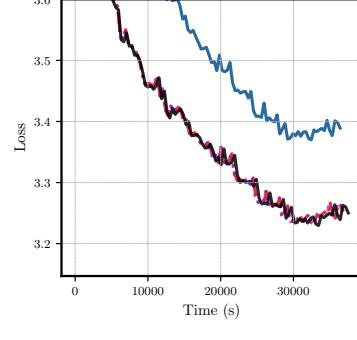
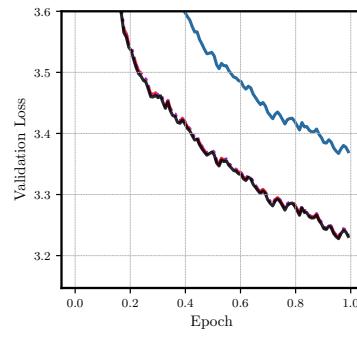
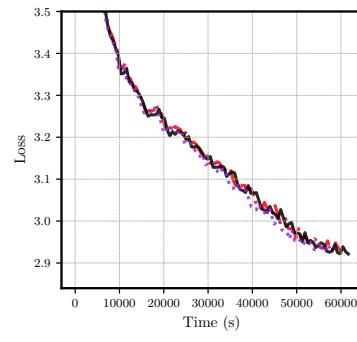
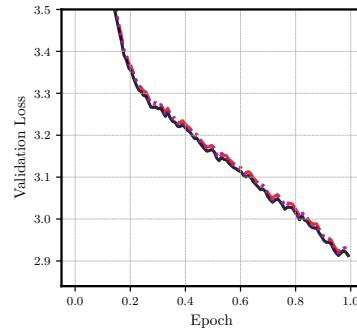
(a) GPT-2-Large (774M params) with weight decay 0.1. Best final validation losses were muon-Jordan (lr = 0.002): 2.921, muon-You (lr = 0.002): 2.919 and muon-PolarExp (lr = 0.002): 2.913.

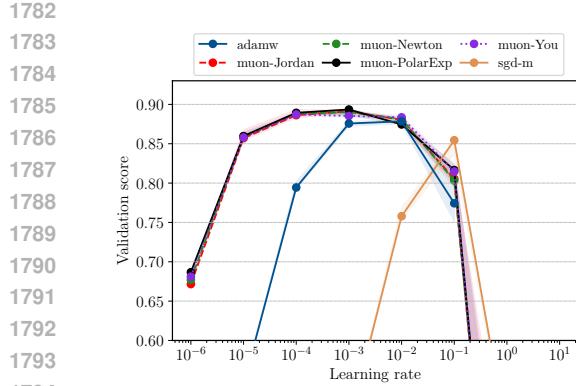


(b) GPT-2-Small (124M params) without weight decay. Best final validation losses were adamw (lr = 0.0005): 3.370, muon-Jordan (lr = 0.005): 3.233, muon-You (lr = 0.005): 3.234 and muon-PolarExp (lr = 0.005): 3.231.

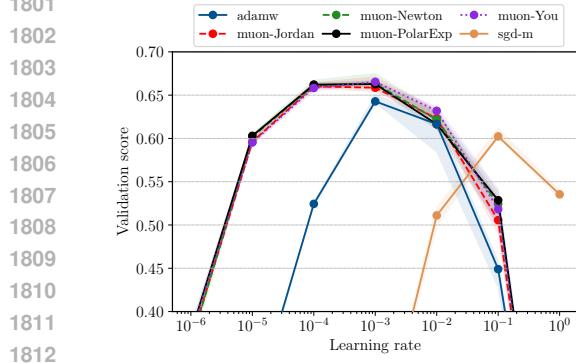
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Figure 13: Training GPT-2 on 10 billion tokens of FineWeb data (Aroca-Ouellette et al., 2023). The label muon-<method> denotes Muon with 5 iterations of <method> to compute polar( $M$ ). Top left: final validation loss vs. learning rate. Bottom left: final training loss vs. learning rate. Top right: validation loss vs. number of iterations for best learning rate. Bottom right: training loss vs. time for best learning rate.

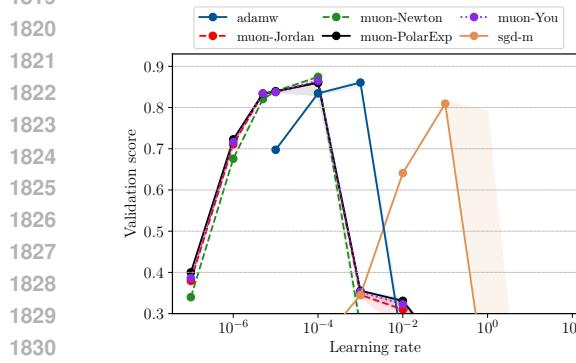




1795 Figure 14: CIFAR10 with a RESNET20. Shaded regions show range over three random seeds.  
1796 The best validation accuracy for each method was sgd-m (lr = 0.1): 0.855 Adamw (lr = 0.01):  
1797 0.878 muon-You (lr = 0.001): 0.887, muon-Newton (lr = 0.001): 0.890, muon-Jordan (lr =  
1798 0.001): 0.891, muon-PolarExp (lr = 0.001): 0.893.  
1799  
1800



1813 Figure 15: CIFAR100 with RESNET110. Shaded regions show range over three random  
1814 seeds. The best validation accuracy for each method was sgd-m (lr = 0.1): 0.602, Adamw  
1815 (lr = 0.01): 0.643, muon-Jordan (lr = 0.001): 0.660, muon-Newton (lr = 0.001): 0.663,  
1816 muon-PolarExp (lr = 0.001): 0.663, muon-You (lr = 0.001): 0.665,  
1817  
1818  
1819



1832 Figure 16: CIFAR10 with a ViT. Shaded regions show range over three random seeds. The best  
1833 validation accuracy for each method was sgd-m (lr =  $10^{-1}$ ): 0.809, muon-PolarExp (lr =  
1834  $10^{-5}$ ): 0.860, Adamw (lr =  $10^{-3}$ ): 0.861, muon-Jordan (lr =  $10^{-5}$ ): 0.861, muon-You (lr =  
1835  $10^{-5}$ ): 0.865, muon-Newton (lr =  $10^{-4}$ ): 0.874.

Letting  $\sigma_1 \geq \dots \geq \sigma_n$  denote the singular values of  $M$ , subspace iteration produces estimates  $\tilde{\sigma}_1 \geq \dots \geq \tilde{\sigma}_k$  satisfying  $\sigma_i \geq \tilde{\sigma}_i$  for all  $i \in 1, \dots, k$ .<sup>7</sup> To find upper bounds on each  $\sigma_i$ , we can use the fact that  $\|M\|_F^2 = \sum_{j=1}^n \sigma_j^2$  as follows:

$$\sigma_i^2 = \|M\|_F^2 - \sum_{\substack{j=1 \\ j \neq i}}^n \sigma_j^2 \leq \|M\|_F^2 - \sum_{\substack{j=1 \\ j \neq i}}^k \sigma_j^2 \leq \|M\|_F^2 - \sum_{\substack{j=1 \\ j \neq i}}^k \tilde{\sigma}_j^2 \quad (24)$$

That is, for each  $i \in [n]$ ,

$$\sigma_i \in \left[ \tilde{\sigma}_i, \sqrt{\|M\|_F^2 - \sum_{\substack{j=1 \\ j \neq i}}^k \tilde{\sigma}_j^2} \right]$$

Setting  $i = k + 1$ , the above also provides an upper bound for the tail of the spectrum,  $\sigma_{k+1}, \dots, \sigma_n$ .

The second step is to find an odd polynomial that well-approximates the constant function on each of these intervals and on the tail simultaneously. For simplicity, we treat only the  $k = 1$  case here. Assume that  $M$  is normalized to  $\|M\|_F = 1$  and let  $z = \tilde{\sigma}_1$  be the lower bound produced by subspace iteration (which reduces to the power method in this case). Then (24) gives  $\sigma_1 \in [z, 1]$  and  $\sigma_2, \dots, \sigma_n \leq \sqrt{1 - z^2}$ . Assume that these intervals do not overlap, that is,  $\sqrt{1 - z^2} \leq z \iff z \geq 1/\sqrt{2}$ . Then we construct the unique odd cubic polynomial  $p(x) = ax + bx^3$  that satisfies  $p(\sqrt{1 - z^2}) = 1$  and  $p(z) = 1$  by setting

$$a = \frac{z^2(z + \sqrt{1 - z^2}) - \sqrt{1 - z^2}}{z\sqrt{1 - z^2}(2z^2 - 1)} \quad b = \frac{\sqrt{1 - z^2} - z}{z\sqrt{1 - z^2}(2z^2 - 1)} \quad (25)$$

Because  $p(0) = 0$  and  $p$  has at most one local extremum on  $\mathbb{R}_{\geq 0}$ , these conditions immediately guarantee that  $p$  is concave-increasing on  $[0, \sqrt{1 - z^2}]$ , so it must lie above the line  $x \mapsto x/\sqrt{1 - z^2}$ . Furthermore,  $p$  is decreasing on  $[\sigma_1, 1]$ , so it maps  $\sigma_1 \in [z, 1]$  to  $[p(1), 1]$ . By minimizing  $p(1)$  over all valid  $z$  (that is, over the interval  $z \in [1/\sqrt{2}, 1]$ ), one can further show that  $p(1) > 1/\sqrt{2}$ , so  $\sigma_1$  cannot be decreased very much by applying  $p$ . Thus, the largest singular value of  $p(M)$  is still at most 1, while the smaller singular values have increased by a potentially large factor of  $1/\sqrt{1 - z^2}$ . When there is a large outlying singular value,  $z$  is close to 1 and this initialization scheme makes much more progress than a standard iteration of PolarExpress would have.

In Figure 17, we demonstrate the benefit of using the  $p$  given by (25) on a synthetic matrix whose spectrum follows a power law decay. That is,  $\sigma_j(M) = j^{-5}$ , so this matrix has a large outlying singular value  $\sigma_1 \gg \sigma_2$ . Applying (25) costs almost as much as performing an iteration of a degree-5 polynomial method, so for fair comparison, we count it as an additional iteration in this plot. For both Newton-Schulz and Polar Express, performing the extra spectrum-aware initialization step described in this section leads to significant speedups in convergence.

## J FAST POLYNOMIAL ITERATION FOR RECTANGULAR MATRICES

In this section, we describe a simple method for applying an iterative polynomial method to a rectangular matrix. For matrices with a large aspect ratio, this method yields significant computational savings. We emphasize that this method is applicable to *any* computation of the form  $(p_T \circ \dots \circ p_1)(X)$ , where each  $p_t$  is an odd polynomial. Thus, it can be used to apply Newton-Schulz or Jordan’s polynomials in addition to our own.

As a preliminary, we first describe the baseline approach. Let  $X \in \mathbb{R}^{m \times n}$  with  $m \geq n$ , where  $\alpha := m/n \geq 1$  is called the aspect ratio. Any odd polynomial  $p$  of degree  $d = 2q + 1$  can be represented as  $p(x) = xh(x^2)$ , where  $h$  is a polynomial of degree  $q$ . Thus,  $p(X) = Xh(X^\top X)$ . Furthermore,  $h$  can be written in a factored form called Horner’s rule to reduce the number of multiplications. For instance, if  $h(y) = a + by + cy^2 + dy^3$ , Horner’s rule gives  $h(y) = a + y(b + y(c + dy))$ . For a

<sup>7</sup>Let  $Q_0 \in \mathbb{R}^{n \times k}$  be a random matrix with orthonormal columns and define  $Q_{t+1}, R_{t+1} = \text{qr}(M^\top M Q_t)$ , where  $\text{qr}$  is the QR decomposition. Subspace iteration outputs the singular values  $\tilde{\sigma}_1, \dots, \tilde{\sigma}_k$  of  $M Q_T, \tilde{\sigma}_1, \dots, \tilde{\sigma}_k$ . By the Cauchy interlacing theorem,  $\tilde{\sigma}_k \leq \sigma_k$ .

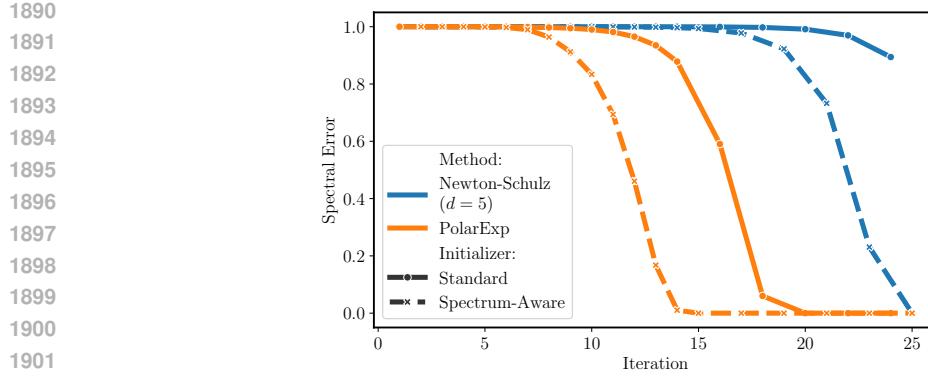


Figure 17: Benefits of the spectrum-aware initialization scheme of [Appendix I](#). Using this scheme improves convergence of both Newton-Schulz and Polar Express on a synthetic  $32 \times 32$  matrix with  $\sigma_j(\mathbf{M}) = j^{-5}$ . Note that we count the spectrum-aware initialization as an additional iteration.

matrix,  $h(\mathbf{Y}) = a\mathbf{I} + \mathbf{Y}(b\mathbf{I} + \mathbf{Y}(c\mathbf{I} + d\mathbf{Y}))$ . Thus for  $\mathbf{Y} \in \mathbb{R}^{n \times n}$ , computing  $h(\mathbf{Y})$  costs about  $(\deg(h) - 1) \cdot n^3$  operations, and computing  $p(\mathbf{X}) = \mathbf{X}h(\mathbf{X}^\top \mathbf{X})$  costs  $2mn^2 + (\frac{d-1}{2} - 1) \cdot n^3 = (\frac{d-3}{2} + 2\alpha) \cdot n^3$  operations. This process could be repeated for each iteration  $p_1, \dots, p_T$ . Notice that if we instead computed  $h(\mathbf{X}\mathbf{X}^\top)\mathbf{X}$ , the result would be the same but the cost would be higher.

A major drawback of this naive approach is that it has a strong dependence on  $\alpha$ , since two rectangular matrix multiplications must be performed in *each* of the  $T$  iterations. When  $m \gg n$ , these two multiplications dominate the cost. In [Algorithm 3](#), we introduce a simple trick that dramatically reduces this cost, using just two rectangular matrix multiplications to compute *all*  $T$  iterations.

---

**Algorithm 3** Fast Polynomial Iteration for Rectangular Matrices

**input:**  $\mathbf{X} \in \mathbb{R}^{m \times n}$  with  $m > 1.5n$ , odd polynomials  $p_1(x) = xh_1(x^2), \dots, p_T(x) = xh_T(x^2)$ .  
**output:** The matrix  $(p_T \circ \dots \circ p_1)(\mathbf{X})$ .

```

 $\mathbf{Y} = \mathbf{X}^\top \mathbf{X}$   $\triangleright mn^2$ 
 $\text{Let } \mathbf{Q}_0 = \mathbf{I}$ 
 $\text{for } t = 1, 2, \dots, T \text{ do}$ 
 $\quad \mathbf{R}_t = \mathbf{Q}_{t-1}^\top \mathbf{Y} \mathbf{Q}_{t-1}$   $\triangleright 2n^3$ 
 $\quad \mathbf{Q}_t = \mathbf{Q}_{t-1} h_t(\mathbf{R}_t)$   $\triangleright \text{Horner's rule: } \deg(h_t) \cdot n^3$ 
 $\text{end for}$ 
 $\text{return } \mathbf{X} \mathbf{Q}_T$   $\triangleright mn^2$ 

```

---

To see why this works, define  $q_0(x) = x$ ,

$$q_t(x) = \frac{(p_t \circ \dots \circ p_1)(x)}{x} = \frac{p_t((p_{t-1} \circ \dots \circ p_1)(x))}{x} = \frac{p_t(xq_{t-1}(x))}{x} \quad (26)$$

$$= \frac{xq_{t-1}(x) \cdot h_t((xq_{t-1}(x))^2)}{x} = q_{t-1}(x) \cdot h_t(x^2 \cdot q_{t-1}(x)^2) \quad (27)$$

and  $r_t(x) = x^2 \cdot q_{t-1}(x)^2$ . It is clear by induction that  $\mathbf{R}_t = r_t(\mathbf{X})$ ,  $\mathbf{Q}_t = q_t(\mathbf{X})$ , and  $\mathbf{X} \mathbf{Q}_T = (p_T \circ \dots \circ p_1)(\mathbf{X})$ . As promised, this algorithm uses no rectangular multiplications in the for-loop. If each  $p_t$  is degree  $d$ , then the total cost is  $(\frac{d+3}{2}T + 2\alpha) \cdot n^3$ . When  $\alpha > 1.5\frac{T}{T-1}$ , this is smaller than the naive method. We can use this criterion to select either [Algorithm 3](#) or the baseline method at runtime.<sup>8</sup>

[Algorithm 3](#) can introduce numerical errors, especially when working in a low precision format like `bfloat16`. We identify two sources of numerical trouble and propose remedies for each. The first is due to the ill-conditioning of  $\mathbf{X}$ . Let  $\mathbf{X} = \mathbf{U}\Sigma\mathbf{V}^\top$  be the SVD. For large  $T$ ,  $(p_T \circ$

<sup>8</sup>Notice that  $\mathbf{Q}_T \rightarrow \mathbf{Y}^{-1/2}$ . This shows that the Polar Express polynomials also give a method of computing the inverse square root of a PSD matrix.

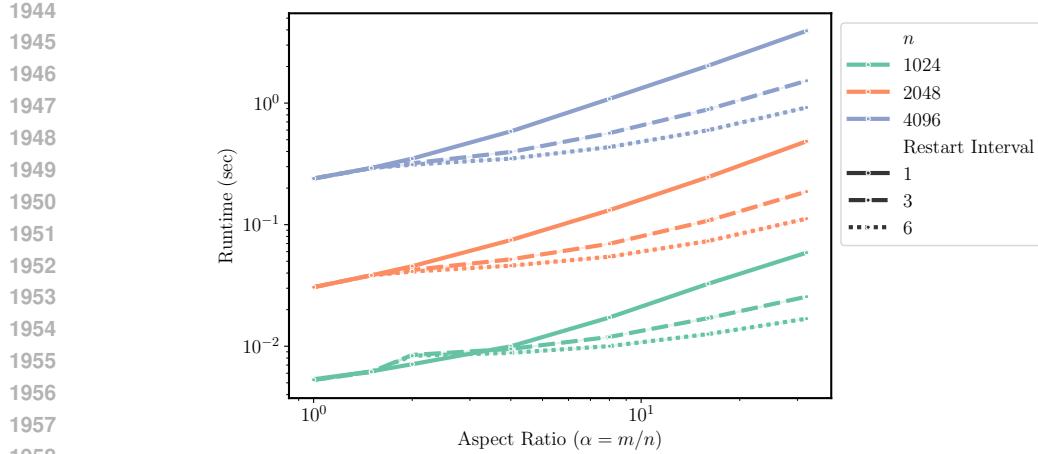


Figure 18: Effects of using [Algorithm 3](#) on runtime on a GPU. We run  $T = 6$  iterations of a degree-5 polynomial method on matrices with various dimensions  $n$  and aspect ratios  $\alpha$ . Restart interval = 6 is [Algorithm 3](#), restart interval = 1 is equivalent to the baseline (that is, not using [Algorithm 3](#)), and restart interval = 3 is an intermediate method that calls [Algorithm 3](#) once to do the first three iterations and again to do the last three iterations for greater stability. When  $\alpha \gg 1$ , increasing the restart interval *significantly* reduces the runtime.

$\cdots p_1)(\mathbf{X}) = \mathbf{X} \mathbf{Q}_T \approx \text{polar}(\mathbf{X}) = \mathbf{U} \mathbf{V}^\top$ . Thus,  $\mathbf{Q}_T \approx \mathbf{V}^\top \Sigma^{-1} \mathbf{V}$ . When  $\mathbf{X}$  has very small singular values and the floating point precision is very low, instantiating  $\mathbf{Q}_T$  may be unstable. To mitigate this issue, we use a restarting strategy. Notice that the issue arises only for large  $T$ , for which  $(p_T \circ \cdots \circ p_1)(\epsilon) \approx 1$ . Limiting ourselves to  $T = 3$  iterations improves the conditioning of  $\mathbf{Q}_T$  because  $(p_T \circ \cdots \circ p_1)(\epsilon) \ll 1$ . Thus, to compute  $T > 3$  iterations, we begin with  $\mathbf{X}_0$  and apply [Algorithm 3](#) with the first three polynomials, producing  $\mathbf{X}_3$ . When then apply [Algorithm 3](#) again with the next three polynomials to  $\mathbf{X}_3$ , producing  $\mathbf{X}_6$ , and so on. As  $\mathbf{X}_t$  approaches convergence, its conditioning improves and we may no longer need to restart at all. Note that restarting [Algorithm 3](#) after every iteration is exactly the same as the baseline method.

Second, while the matrix  $\mathbf{Y}$  is positive definite in exact arithmetic, numerical round-off can introduce spurious negative eigenvalues that cause the method to diverge to infinity. To combat this issue, we instead set  $\mathbf{Y} = \mathbf{X}^\top \mathbf{X} + 10^{-3} \mathbf{I}$  during the first application of [Algorithm 3](#). (We also normalize by  $\|\mathbf{X}\|_F + 10^{-3}$  instead of  $\|\mathbf{X}\|_F$ .) In subsequent restarts of [Algorithm 3](#), we set  $\mathbf{Y} = \mathbf{X}^\top \mathbf{X}$  as before. This is akin to slightly increasing each of the singular values of  $\mathbf{X}$ , but it does *not* change the polar factor of  $\mathbf{X}$ . Thus, while the output will be slightly different in the early iterations, the algorithm still converges to the correct answer.

[Figure 18](#) shows that using [Algorithm 3](#) can significantly improve runtime on the GPU when the aspect ratio is large enough. As expected, using [Algorithm 3](#) for many iterations significantly reduces the dependence of the runtime on the aspect ratio. Running six iterations of a degree-5 polynomial method when  $\alpha = 4$  (as with the linear transformations in each MLP block of a transformer) we obtain almost a 2x speedup, and when  $\alpha = 32$ , we obtain a 5x speedup. If we restart every three iterations, the trend is the same but the runtime savings are somewhat smaller.

## J.1 APPLICATION TO MUON

If these problems can be mitigated, the speed afforded by [Algorithm 3](#) suggests an improvement in the way Muon is applied to transformers. In sum, the idea is to replace one large matrix with a small aspect ratio by many smaller matrices with large aspect ratios and apply [Algorithm 3](#) to all of them in parallel. Each multi-head attention layer contains four square weight matrices  $\mathbf{W}_Q, \mathbf{W}_K, \mathbf{W}_V$  and  $\mathbf{W}_O \in \mathbb{R}^{d \times d}$ . The orthogonalization step of Muon is either applied separately to these four matrices or else to  $[\mathbf{W}_Q \mid \mathbf{W}_K \mid \mathbf{W}_V]$  and  $\mathbf{W}_O$ , since typical implementations of multi-head attention store the weights in this concatenated form. However, we believe it is natural to consider each of these four weight matrices to be a concatenation of many smaller linear transformations, each corresponding

1998 to a single attention head. If  $H$  is the number of heads, each of these smaller matrices has size  
 1999  $d \times \frac{d}{H}$ ; that is, they have aspect ratio  $\alpha = H$ . The gradient matrices of  $[\mathbf{W}_Q \mid \mathbf{W}_K \mid \mathbf{W}_V]$  and  $\mathbf{W}_O$   
 2000 can be reshaped into 3-tensors in which each slice is one of these smaller matrices. Since typical  
 2001 transformers like GPT-3 can have as many as 96 heads, this variation of Muon has the potential to  
 2002 reduce the runtime.

2003 We use this idea to train a GPT-Small model on FineWeb1B. We compare four conditions:  
 2004

- 2005 1. The baseline approach used in the rest of this paper
- 2006 2. Splitting up the gradient matrices of  $[\mathbf{W}_Q \mid \mathbf{W}_K \mid \mathbf{W}_V]$  and  $\mathbf{W}_O$  by head and applying  
 2007 Muon to each piece, as described above
- 2008 3. Using [Algorithm 3](#), restarted after three iterations
- 2009 4. Splitting by head *and* using [Algorithm 3](#)

2010 We used `Polar Express` with weight decay of 0.1 for all conditions and swept learning rates  
 2011 0.003, 0.005, 0.01. Otherwise, all hyperparameters were the same as in [Section 4.2](#).  
 2012

2013 Our results showed that these changes had a negligible effect in this setting. They did not affect  
 2014 the optimization quality. Compared to the baseline, splitting by heads actually reduced the final loss  
 2015 slightly from 3.59 to 3.55; using [Algorithm 3](#) increased the loss very slightly, from 3.59 to 3.60 when  
 2016 not splitting by head, and from 3.55 to 3.56 when we did split. However, the runtimes of all 12 runs  
 2017 were nearly identical, showing that at this scale, the FLOP savings of [Algorithm 3](#) is not beneficial.  
 2018 The embedding size of GPT-Small is just 768. These techniques may be more impactful when using  
 2019 a larger model. It may also have more impact outside of deep learning, where `Polar Express`  
 2020 would be run for more than the 5 iterations used in our experiments. We leave exploration of these  
 2021 settings to future work.  
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