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ABSTRACT

This paper presents a unified framework for simulating and visualizing dynamic fracture phenomena in extreme mechanical collisions using multi-view image inputs. While existing methods primarily address elastic deformations at contact surfaces, they fail to capture the complex physics of extreme collisions, often producing non-physical artifacts and material adhesion at fracture interfaces. Our approach integrates two key innovations: (1) an enhanced Collision Material Point Method (Collision-MPM) with momentum-conserving interface forces derived from normalized mass distributions, which effectively eliminates unphysical adhesion in fractured solids; and (2) a fracture-aware 3D Gaussian continuum representation that enables physically plausible rendering without post-processing. The framework operates through three main stages: First, performing implicit reconstruction of collision objects from multi-view images while sampling both surface and internal particles and simultaneously learning surface particle Gaussian properties via splatting; Second, high-fidelity collision resolution using our improved Collision-MPM formulation; Third, dynamic fracture tracking with Gaussian attribute optimization for fracture surfaces rendering. Through comprehensive testing, our framework demonstrates significant improvements over existing methods in handling diverse scenarios, including homogeneous materials, heterogeneous composites, and complex multi-body collisions. The results confirm superior physical accuracy, while maintaining computational efficiency for rendering.

1 INTRODUCTION

Dynamic fracture simulation stands at the intersection of computational physics and computer graphics, enabling realistic modeling of material failure across diverse applications. While physics-based methods like the Material Point Method (MPM) have advanced significantly since their introduction (Stomakhin et al., 2013), critical gaps remain in handling extreme mechanical collisions and achieving seamless simulation-to-rendering pipelines.

Recent advances in explicit scene representation, particularly 3D Gaussian splatting (Kerbl et al., 2023), have revolutionized real-time rendering capabilities. Building upon this foundation, several studies have successfully integrated physical simulation with Gaussian representations, including (Borycki et al., 2024; Cai et al., 2024; Feng et al., 2024; Xie et al., 2024; Zhang et al., 2025). Among these, (Xie et al., 2024) established a significant milestone by coupling 3D Gaussians with MPM simulations, demonstrating remarkable adaptability across various material types. Parallel developments include (Cai et al., 2024)'s physics parameter estimation through Gaussian differentiability and (Borycki et al., 2024)'s GASP framework for point-wise physical attribute embedding. While these approaches have made substantial progress in bridging the simulation-rendering gap, their applicability remains constrained to moderate mechanical conditions and specific material classes. A critical limitation emerges when addressing extreme mechanical collisions, such as the high-energy fragmentation observed in brittle materials (Wolper et al., 2019). Current methodologies face two fundamental challenges in these scenarios: (1) unphysical adhesion artifacts in MPM simulations, and (2) inadequate fracture surface representation for rendering. These challenges become particularly apparent in high-energy impact scenarios like brittle material fragmentation or multi-body collisions.

To address these limitations, we introduce Fracture-GS, a unified framework for simulating and visualizing dynamic fracture phenomena in extreme mechanical collisions using multi-view image inputs. First, the Signed Distance Function (SDF) of the object is constructed from multi-view im-



Figure 1: Two extreme mechanical collision scenarios are simulated using our proposed framework: (1) single-object impacts against wind, and (2) multi-object collision interactions with complex fracture dynamics. Notably, the flowerpot object comprises heterogeneous materials (with distinct properties for leaves, stems, and the pot itself), and the table object consists of legs and tabletop, while other teapot is modeled as homogeneous materials.

ages to implicitly represent the volumetric geometry of the object, followed by sampling both surface and internal particles within the SDF-constrained domain to ensure spatial coherence; Meanwhile, surface particles learn Gaussian attributes using isotropic kernels. Subsequently, we proposed an enhanced collision-MPM, which is used to perform extreme collisions between multiple objects with dynamic fracture. It can effectively alleviate the non-physical adhesion phenomenon caused by MLS-MPM. Then, all fracture particles are tracked through a hardening-aware tracking criterion defined by (Wolper et al., 2019). Finally, based on the tracking fracture particles, we can efficiently regenerate their Gaussian attributes through the proposed fracture particles Gaussian optimization strategy, enabling high-quality rendering, as shown in Figure 1.

In summary, our key contributions are:

- A unified physics-rendering framework that combines our enhanced Collision-MPM with fracture-aware 3D Gaussian representations, enabling high-fidelity simulation and visualization of extreme mechanical collisions across diverse materials.
- An improved Collision-MPM formulation that introduces momentum-conserving interface forces derived from normalized mass distributions, effectively eliminating the non-physical adhesion artifacts prevalent in conventional MPM simulations.
- A dynamic fracture tracking and rendering pipeline that identifies fracture surfaces through hardening parameter and reconstructs Gaussian attributes via minimal-volume enclosing ellipsoid (MVEE) optimization.

2 PRELIMINARY AND RELATED WORK

2.1 3D GAUSSIAN SPLATTING AND DYNAMIC SCENE RECONSTRUCTION

3D Gaussian Splatting (3DGS) is a fast neural rendering method that primarily represents volumetric scenes using a collection of anisotropic 3D Gaussian kernels. Each Gaussian kernel is defined by a set of differentiable parameters $\{x_p, \sigma_p, A_p, c_p\}$, where x_p denotes the spatial position, σ_p represents opacity, A_p is the covariance matrix, and c_p is the view-dependent color function. The covariance matrix A_p can be further decomposed into scale s_p and rotation r_p components, which control the spatial distribution and orientation of the Gaussian kernel, respectively. The color function c_p achieves view-dependent characteristics through spherical harmonics, capturing complex lighting and material effects. The rendering process involves projecting (splatting) the 3D Gaussian kernels onto a 2D image plane, incorporating viewpoint transformation, opacity blending, and depth sorting. The final color of the i -th pixel is computed using the following formula:

$$C_i = \sum_{k \in I} \sigma_k c_k (d_i) \prod_{j=1}^{k-1} (1 - \sigma_j) \quad (1)$$

108 Here, I stands for the set of Gaussian kernels, σ_k and σ_j represent the opacity of the k -th and j -th
 109 Gaussian kernels, respectively. d_i means the viewing direction from the camera to the i -th pixel, and
 110 $c_k(d_i)$ indicates the color of the Gaussian kernels in the viewing direction d_i . Since 3D Gaussian
 111 explicitly represents the scene, its learning and rendering speeds are generally faster than those of
 112 3D reconstruction methods based on NeRF. As a result, a wide range of applications based on 3D
 113 Gaussian have emerged.

114 Dynamic 3D scene reconstruction has long been a challenging problem, aiming to reconstruct dy-
 115 namic scenes from various representations such as videos and images. The introduction of Neural
 116 Radiance Fields (NeRF) has significantly advanced this field, leading to a series of dynamic scene
 117 reconstruction methods that build upon NeRF’s framework. These methods focus on addressing
 118 challenges such as non-ideal input conditions, including sparse views and motion blur, to enhance
 119 reconstruction quality. Notable works in this domain include (Pumarola et al., 2021), which handles
 120 dynamic scenes by incorporating temporal information, and (Fang et al., 2022), which improves ef-
 121 ficiency through time-aware neural voxels. More recently, (Song et al., 2023) proposed a streamable
 122 dynamic scene representation that decomposes neural radiance fields for efficient reconstruction and
 123 rendering.

124 Another approach to scene representation is based on 3D Gaussian Splats (3DGS), which explicitly
 125 represents 3D scenes using a set of Gaussian kernels and achieves fast rendering through Gaussian
 126 splatting. Due to its high efficiency and interpretability, several dynamic scene reconstruction meth-
 127 ods based on 3DGS have been proposed, such as (Lin et al., 2024; Huang et al., 2024; Wu et al.,
 128 2024; Yang et al., 2024; Sun et al., 2024; Zhang et al., 2024; Dahmani et al., 2024). Among these,
 129 (Lin et al., 2024) incorporates DDDM (Deformable Dynamic Model) into the optimization pro-
 130 cess of Gaussians, eliminating the need to reconstruct Gaussian for each frame and directly guiding
 131 Gaussian deformation using DDDM. On the other hand, (Huang et al., 2024) learns control points
 132 for Gaussian and uses a small number of control points to govern the motion of the entire Gaussian
 133 set, achieving efficient reconstruction of dynamic motion processes. (Dahmani et al., 2024) divides
 134 the dynamic sequence into different windows based on the motion number, and train dynamic Gaus-
 135 sian models for different windows, together with different canonical spaces and deformation fields.

136 2.2 MATERIAL POINT METHOD AND PHYSICS-BASED GAUSSIAN APPROACHES

137 The Material Point Method (MPM) is a numerical approach based on a hybrid Eulerian-Lagrangian
 138 framework, used to solve governing equations in continuum mechanics and facilitate bidirectional
 139 information transfer between particles and grids. Let $p \in \{P_a, P_b\}$ denote the classification criterion
 140 for distinguishing particle subsets. Following established methodologies, the process adheres to the
 141 MLS-MPM framework (Hu et al., 2018), which is divided into the following three stages:

- 142 • **Particle-to-Grid (P2G) Stage:** The particle mass m_p and velocity v_p are interpolated to
 143 neighboring grid nodes through a weighted projection scheme governed by basis functions:

$$144 \quad m_i^n = \sum_p w_{ip}^n m_p, \quad (2)$$

$$145 \quad m_i^n v_i^n = \sum_p w_{ip}^n m_p (v_p^n + C_p^n(x_i - x_p^n)), \quad (3)$$

146 where w_{ip}^n represents the interpolation kernel (e.g., quadratic B-spline) evaluated at particle
 147 position x_p for grid node i , P is the set of active particles, and m_i and v_i are the aggregated
 148 mass and velocity at grid node i , respectively.

- 149 • **Grid Update Stage:** Grid velocities are advanced by solving the discrete momentum con-
 150 servation equations through an explicit forward Euler integration scheme:

$$151 \quad m_i^n (v_i^{n+1} - v_i^n) = -\Delta t \cdot f_i^* + \Delta t \cdot f_i^{ext}, \quad (4)$$

$$152 \quad f_i^* = \sum_p \frac{4}{\Delta x^2} V_p w_{ip}^n \sigma_p^n (x_i^n - x_p^n), \quad (5)$$

153 where Δx denotes the grid size, f_i^* is the grid force calculated from the particle volume
 154 V_p , Cauchy stress σ_p^n , and positions x_i^n and x_p^n , and f_i^{ext} is the external force (typically
 155 gravity).

162 • **Grid-to-Particle (G2P) Stage:** The updated grid velocities v_i^{n+1} are mapped back to La-
 163 grangian particles to update their kinematic states for the next time step $n+1$. This transfer
 164 is achieved through interpolation:
 165

$$166 \quad v_p^{n+1} = \sum_i w_{ip}^n v_i^{n+1}, \quad (6)$$

167

$$168 \quad C_p^{n+1} = \frac{4}{\Delta x^2} \sum_i w_{ip}^n v_i^{n+1} (x_i^n - x_p^n)^T. \quad (7)$$

170

171 After obtaining the updated particle velocities, the particle positions are advanced through an explicit
 172 time integration scheme:
 173

$$x_p^{n+1} = x_p^n + \Delta t v_p^{n+1}. \quad (8)$$

174 The Material Point Method (MPM) has emerged as a powerful tool for simulating complex physical
 175 phenomena, combining the advantages of both Lagrangian and Eulerian approaches. Recent ad-
 176 vancements in MPM simulations have significantly expanded its applications and performance opti-
 177 mization. (Hu et al., 2018) introduces a variant of the MPM based on Moving Least Squares (MLS),
 178 referred to as MLS-MPM, for simulating complex physical phenomena involving displacement dis-
 179 continuities and bidirectional rigid-body coupling. By incorporating the Compatible Particle-In-Cell
 180 (CPIC) algorithm, this method enables the handling of discontinuities in material points, infinitely
 181 thin boundaries, and bidirectional coupling with rigid bodies. As a result, it is capable of simulating
 182 material cutting, dynamic open boundaries, and interactions between rigid and deformable bodies.
 183 (Wolper et al., 2019) integrates a phase field into the MPM to develop a crack-tracking approach
 184 known as PFF-MPM. Additionally, it proposes an incompressible plastic flow rule that maintains
 185 constant volume during plastic stress projection. Meanwhile, (Fang et al., 2020) proposed a novel
 186 framework for fluid-solid coupling using IQ-MPM. This method combines a "ghost matrix" operator
 187 splitting scheme with weak-form governing equations to achieve stable and efficient coupling un-
 188 der the CFL time step constraint. It supports discrete consistency with hybrid Lagrangian-Eulerian
 189 solvers and uses an interface quadrature (IQ) technique to handle free-slip boundaries, avoiding the
 190 "stickiness" issues in traditional MPM implementations. Moreover, the effectiveness of employing
 191 GPUs to enhance the computational efficiency of MPM implementations has been well-documented
 192 in several studies, such as (Gao et al., 2018; Hu et al., 2019; Qiu et al., 2023). However, these meth-
 193 ods all require post-processing of the simulation results to achieve high-quality rendering effects.
 194

195 The integration of Gaussian-based techniques into physical simulations leverages the rendering ef-
 196 ficiency of Gaussian, thereby eliminating the need for post-processing with dedicated rendering
 197 engines such as Houdini after the completion of the physical simulation, and directly yielding re-
 198 sults with Gaussian splatting. Specifically, (Xie et al., 2024) incorporates Gaussian kernels into
 199 the dynamic simulation process of MPM, enabling continuum mechanics simulations based on 3D
 200 Gaussian kernels and achieving real-time rendering in simple scene simulations. (Zhang et al., 2025)
 201 uses diffusion on images to obtain prior motion videos of objects, simulates the motion process based
 202 on MPM, leverages the differentiability of MPM to learn the material field in specified regions, and
 203 completes the forward simulation process to generate dynamic videos. (Cai et al., 2024) aims to
 204 guide the learning of physical properties of objects using 3D Gaussian splats (3DGS). It first recon-
 205 structs the static Gaussians in the initial state and learns the deformation models of Gaussians based
 206 on dynamic inputs, optimizing the initial velocities and physical parameters (Young's modulus and
 207 Poisson's ratio) of the static Gaussians through differentiable MPM. (Tan et al., 2024) generates
 208 high-quality, physics-based videos from a single image. However, none of these methods address
 209 the extreme cases where objects undergo fragmentation, which is the primary focus of this work.

207 3 METHOD

209 In this section, we introduce Fracture-GS, a unified framework for simulating and visualizing dy-
 210 namic fracture phenomena in extreme mechanical collisions using multi-view image inputs. As
 211 illustrated in Figure 2, we first reconstruct the geometry of colliding objects using multi-view im-
 212 ages combined with existing implicit 3D reconstruction algorithms (Xiao et al., 2024), followed by
 213 sampling both surface and internal particles within the SDF-constrained domain to ensure spatial
 214 coherence; Subsequently, surface particles Gaussian kernels are trained using 3D Gaussian Splat-
 215 tting from input images. Then, we proposed an enhanced collision-MPM, which is used to perform
 216 extreme collisions between multiple objects with dynamic fracture. It can effectively alleviate the

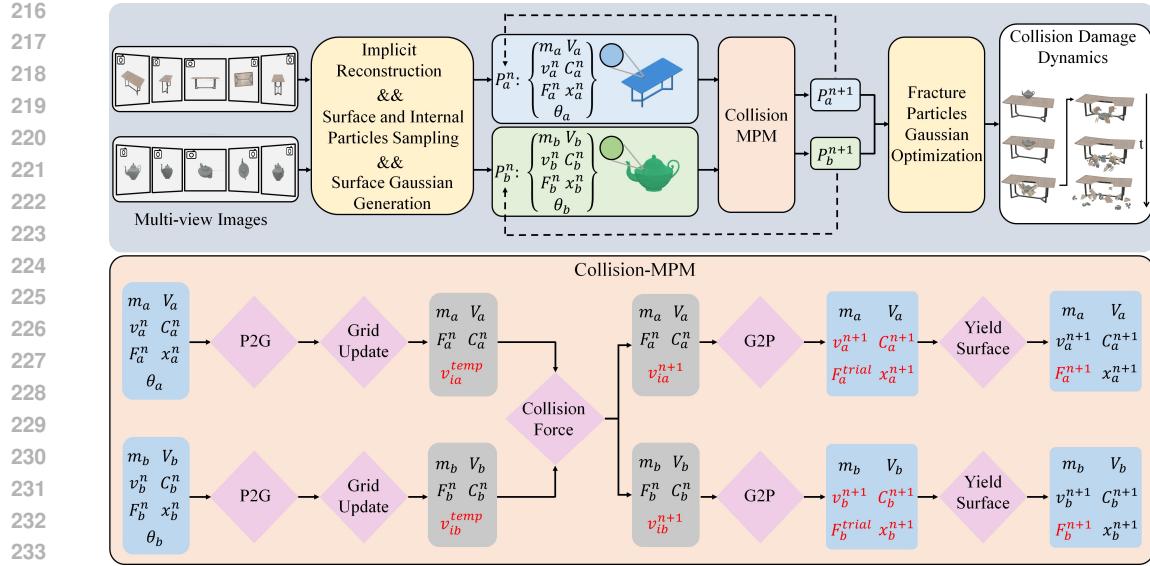


Figure 2: Pipeline. The object is implicitly reconstructed from multi-view images, followed by sampling both surface and internal particles. Surface particles learn Gaussian attributes using isotropic kernels. Next, the sampled particles undergo extreme mechanical collision simulation with dynamic fracture using our enhanced **Collision-MPM**. Finally, fracture particles are tracked and their Gaussian attributes are optimized through our proposed **Fracture Particle Gaussian Optimization** strategy, enabling high-quality rendering of the simulation results. For **Collision-MPM**, the key parameters are highlighted in red. The yield surface determines whether a particle enters the plastic region, triggering a return mapping procedure to project stress back to the yield surface and update the particle’s deformation gradient. Parameters in the blue bottom plate are computed in the Lagrangian coordinate system, while those in the gray bottom plate are computed in the Eulerian coordinate system.

non-physical adhesion phenomenon caused by MLS-MPM. Finally, to enhance the visual realism of the mechanical simulation, all fracture particles are tracked through a hardening-aware tracking criterion defined by (Wolper et al., 2019), based on the tracking fracture particles (as shown in Figure 3 (right)), we can efficiently regenerate their Gaussian attributes through the proposed fracture particles Gaussian optimization strategy, enabling high-quality rendering.

The all particle attributes of the colliding objects $P_a : \{m_a, V_a, C_a, F_a, v_a, x_a, \theta_a\}$ and $P_b : \{m_b, V_b, C_b, F_b, v_b, x_b, \theta_b\}$ include mass (m), volume (V), deformation gradient (F), velocity gradient (C), velocity (v), position (x), and elastoplastic parameters (θ), where θ comprises Young’s modulus (E), Poisson’s ratio (γ), hardening tracking parameters (α), cohesion coefficient (β), and hardening factor (ξ).

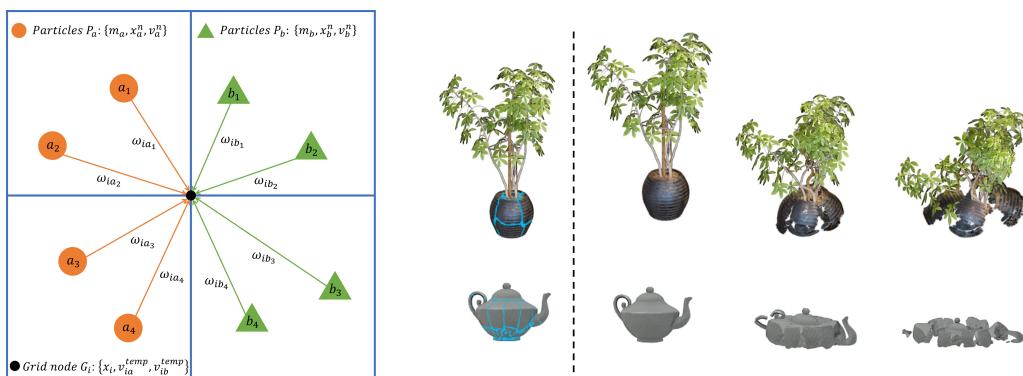


Figure 3: (left) Illustration of the mass distributions of particles P_a and P_b at grid node G_i . (right) Tracking fracture particles through hardening factor α .

270 3.1 COLLISION-MPM
271

272 Although the Material Point Method (MPM) itself can solve the problem of collisions between
273 objects, its performance has limitations. For instance, in fluid-solid coupling simulations, fluid
274 particles may unnaturally adhere to solid surfaces—a phenomenon similarly observed in multi-body
275 fracture simulations, where fragmented solids exhibit unphysical adhesion to neighboring objects.
276 To mitigate these artifacts, we propose a Collision-MPM framework, integrating the collision force
277 (Yan et al., 2018) into the MPM framework. This method, originally designed to prevent fluid-solid
278 interpenetration, ensures collision realism across material interfaces.

279 In our **Collision-MPM** framework, the particle information of two objects is independently trans-
280 ferred to the grid. First, after completing the P2G and Grid Update processes, the mass distributions
281 \hat{n}_{ia} and \hat{n}_{ib} of object particles P_a and P_b on grid node G_i are calculated. Simultaneously, the grid
282 velocities v_i^{n+1} are replaced with v_i^{temp} . Subsequently, the interface direction tendencies n_{ia} and
283 n_{ib} at grid node G_i for object particles P_a and P_b are determined:

$$\hat{n}_{ia} = \frac{\sum_{p_a} m_a \nabla \omega_{ia}(x_i - x_a^n)}{\left\| \sum_{p_a} m_a \nabla \omega_{ia}(x_i - x_a^n) \right\|} \quad (9)$$

$$\hat{n}_{ib} = \frac{\sum_{p_b} m_b \nabla \omega_{ib}(x_i - x_b^n)}{\left\| \sum_{p_b} m_b \nabla \omega_{ib}(x_i - x_b^n) \right\|} \quad (10)$$

290 As shown in Figure 3 (left), \hat{n}_{ia} and \hat{n}_{ib} illustrate the mass distributions of particles P_a and P_b on
291 grid node G_i . We utilize these distributions to compute the interface direction tendency, which can
292 also be interpreted as the tendency of the contact surface’s normal direction:
293

$$n_{ia} = -n_{ib} = \frac{\hat{n}_{ia} - \hat{n}_{ib}}{\|\hat{n}_{ia} - \hat{n}_{ib}\|} \quad (11)$$

294 The computation of collision forces is conditionally activated based on relative velocity analysis at
295 material interfaces. Specifically, collision forces are generated at grid node G_i only if the inequality
296 $(v_{ia}^{temp} - v_{ib}^{temp}) \cdot n_{ia} > 0$ is satisfied, where v_{ia}^{temp} and v_{ib}^{temp} denote the velocities of P_a and P_b
297 particles on grid node G_i after the grid update. The collision forces are derived from the momentum
298 conservation principle during the collision process:
299

$$f_i^c = \frac{p_{ia}^{temp} m_{ia}^n - p_{ib}^{temp} m_{ib}^n}{(m_{ia}^n + m_{ib}^n) \Delta t} \quad (12)$$

$$f_{ia}^c = -f_{ib}^c = \mu(f_i^c \cdot n_{ib}) n_{ib} \quad (13)$$

300 Here, f_i^c represents the collision force on grid node i , m_{ia}^n and m_{ib}^n denote the mass contributions of
301 P_a and P_b particles on grid node G_i , respectively. $p_{ia}^{temp} = m_{ia}^n v_{ia}^{temp}$ and $p_{ib}^{temp} = m_{ib}^n v_{ib}^{temp}$ rep-
302 resent the momenta of P_a and P_b particles on grid node G_i , respectively. μ is a constant controlling
303 the magnitude of the collision force. $f_{ia, collision}$ and $f_{ib, collision}$ denote the collision forces acting
304 on P_a and P_b particles on grid node G_i . After computing the collision forces, they are integrated
305 into the grid velocity update:
306

$$v_{ia}^{n+1} = v_{ia}^{temp} + \frac{f_{ia}^c}{m_{ia}^n} \Delta t, \quad v_{ib}^{n+1} = v_{ib}^{temp} + \frac{f_{ib}^c}{m_{ib}^n} \Delta t \quad (14)$$

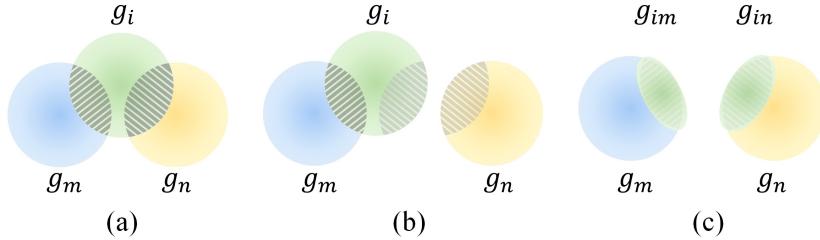
307 where v_{ia}^{n+1} and v_{ib}^{n+1} represent the updated velocities of P_a and P_b particles at grid node G_i ,
308 respectively.
309

310 3.2 FRACTURE PARTICLES GAUSSIAN OPTIMIZATION (FPGO).
311312 3.2.1 CONTINUUM MECHANICS AND CONSTITUTIVE MODEL.
313

314 To simulate the dynamic behavior of an elastic-plastic object, it is essential to solve the conservation
315 equations for momentum and mass:
316

$$\rho \frac{Dv}{Dt} = \nabla \cdot \sigma + f, \quad \frac{D\rho}{Dt} + \rho \nabla \cdot v = 0 \quad (15)$$

324 Here, ρ denotes density, v represents the velocity field, and f is an external force. The Cauchy
 325 stress tensor, denoted by σ , is given by: $\sigma = \frac{1}{\det(F)} \frac{\partial \psi(F^E)}{\partial F} F^{E^T}$, where $\psi(F)$ is the strain energy
 326 density function (or constitutive model), which describes the relationship between stress and strain
 327 in a material. The total deformation gradient, F , is decomposed into elastic and plastic components:
 328 $F = F^E F^P$, enabling the simulation of plastic deformation. In this work, we use the NACC model
 329 in (Wolper et al., 2019). It extends CCC (Coherent Cam Clay) model with non-associated flow
 330 rules to better simulate plastic deformation while maintaining volume during plastic projection, and
 331 it introduced four plastic parameters α, β, ξ, M to control the simulation effect of the plastic model.
 332
 333
 334



343 **Figure 4: Illustration of Fracture Particles Gaussian Optimization**

344 3.2.2 FRACTURE TRACKING AND RENDERING

345 The enhanced Collision-MPM framework successfully simulates extreme mechanical collisions
 346 with dynamic fracture phenomena. While one might intuitively consider directly rendering using
 347 the initially learned Gaussian attributes of each surface particles, this naive approach leads to signif-
 348 icant non-physical artifacts at fracture interfaces (highlighted by red boxes in Figure 5).

349 The underlying mechanism of these artifacts can be explained as follows: As shown in Figure 4, con-
 350 sider three Gaussian particles $\{g_i, g_m, g_n\}$ in initial configuration. When fracture occurs at particle
 351 g_i , causing displacement relative to g_m (Figure 4(b)), the increased interparticle distance reduces or
 352 eliminates the Gaussian overlap region between g_i and g_m , thereby disrupting the continuity of the
 353 rendering field and generating visual artifacts (As demonstrated in Figure 4(a) and (b)).

354 To address this challenge, we propose a novel fracture-aware Gaussian attribute optimization strat-
 355 egy consisting of four key components. First, leveraging the hardening parameter α from the NACC
 356 constitutive model (Wolper et al., 2019), we dynamically identify fractured particles at each timestep
 357 (visualized as green particles in Figure 4). Taking particle g_i as an example (Figure 4), our tracking
 358 begins when its hardening parameter exceeds the hardening parameter threshold α .

359 For each identified fractured particle g_i , we perform neighborhood analysis within a radius d_c to
 360 locate adjacent intact particles $\{g_m, g_n\}$. This adaptive search range ensures proper coverage of
 361 potential interaction zones while maintaining computational efficiency.

362 The core optimization involves Gaussian cloning and attribute reconstruction. We first compute the
 363 minimal-volume enclosing ellipsoid (MVEE) for the overlap regions between g_i and its neighbors
 364 g_m, g_n , generating two new Gaussian particles g_{im} and g_{in} to replace the original g_i . The attribute
 365 assignment follows two principles: (1) optical properties including opacity α and color c are directly
 366 inherited from g_i to maintain visual consistency; (2) spatial parameters are recomputed through:

$$367 \{ \mu_{new}, \Sigma_{new} \} = \text{MVEE}(\Omega_{cross}(g_i, g_j)) \quad (16)$$

368 where Ω_{cross} denotes the original overlap region (Implementation details are provided in Appendix
 369 A.4.).

370 During final rendering, we implement an occlusion-aware sampling scheme: if a pixel’s rendering
 371 path contains multiple optimized particles (g_{im}, g_{in}), only the nearest particle contributes to shading.
 372 This prevents overcounting while preserving physical correctness. As demonstrated in Fig 1, our
 373 approach generates physically plausible transitional particles that maintain both visual continuity
 374 across fracture surfaces and mechanical accuracy in collision regions.

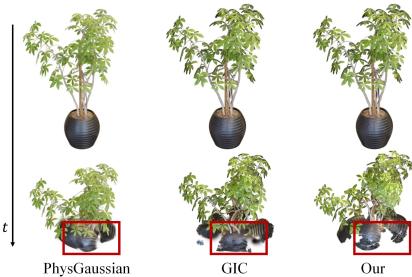
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4 EXPERIMENTS

379

4.1 EXPERIMENTAL DATA AND PHYSICAL PARAMETER SETTINGS

380 To comprehensively validate the effectiveness of our experiments, we selected three objects for col-
 381 lision (More experimental results are provided in the supplementary material): Ficus, Teapot and
 382 Table. Among these, the Ficus plant and the table are heterogeneous material objects. Specifically,
 383 the Ficus plant consists of leaves, branches, and a ceramic pot, each made of different materials,
 384 while the table has a tabletop and legs constructed from distinct materials. The teapot is homoge-
 385 neous in material composition. The physical parameters for all objects are detailed in Appendix
 386 A.1. For specific implementation details, please refer to A.2. Additional experimental results are
 387 provided in A.5 and Appendix A.6.



390 Figure 5: Integrated comparison showing both qualitative visualizations (left) and quantitative met-
 391 400 rics (right), with our method achieving the best results.

401

4.2 COMPARISON WITH STATE-OF-THE-ART METHODS

402 We compare our method with two state-of-the-art Gaussian splatting based simulation frameworks.
 403 First, **PhysGaussian** (Xie et al., 2024) is a physics-integrated Gaussian framework that simulates
 404 and renders mechanical behaviors under external forces from multi-view inputs. Second, **GIC** (Cai
 405 et al., 2024), originally designed for material property estimation from videos, is adapted for com-
 406 parison by utilizing only its forward simulation component with given material parameters, analogous
 407 to PhysGaussian. All methods employ identical initialization conditions to ensure fair comparison:
 408 the same static 3D Gaussian reconstruction pipeline following GIC’s methodology; the same NACC
 409 constitutive model was adopted for physical simulation, but the above two comparison methods did
 410 not include FPGO and used the MLS-MPM; and identical initial conditions and material parameters.
 411 We use PSNR (Hore & Ziou, 2010), LPIPS (Zhang et al., 2018), and FID (Heusel et al., 2017) as
 412 primary metrics to evaluate reconstruction quality. Due to the absence of ground truth for dynamic
 413 fracture sequences, we employ a self-referencing evaluation scheme using established image quality
 414 metrics. The specific implementation details of our self-referencing metric calculation are provided
 415 in Appendix A.3.

416 **User Study.** We also conducted a human evaluation to assess simulation fidelity, following meth-
 417 ods from prior work (Liu et al., 2025; Wei et al., 2024). Ten participants with varying experience
 418 in simulation and vision rated to **Fracture Simulation Fidelity (FSF)**, checking if it was realistic
 419 and as expected. Rendered videos of simulations were presented in random order, with participants
 420 rating each on a five-point scale (1 = poor, 5 = excellent). Mean scores appear in Figure 5 (right).

421 **Results.** Quantitative and qualitative results in Figure 1 and 5. Our method significantly outper-
 422 forms other methods in GS-based simulation. Competing methods often produce artifacts due to
 423 inadequate or neglected handling of fracture surfaces and collision adhesion, which degrades sim-
 424 ulation quality. Although our approach excels in simulation, the object’s FID is high due to reliance
 425 on training view interpolation for Gaussian restoration and cannot perform Gaussian re-
 426 construction on hidden areas. Future work will explore 3D AI-based texture generative inpainting
 427 to improve this.

428 **Effect of Collision-MPM (C-MPM) and Fracture Particles Gaussian Optimization (FPGO).**
 429 To validate C-MPM effectiveness, we conduct ablation studies comparing against the conven-
 430 tional MLS-MPM approach. As demonstrated in Figure 6, the right column reveals significant

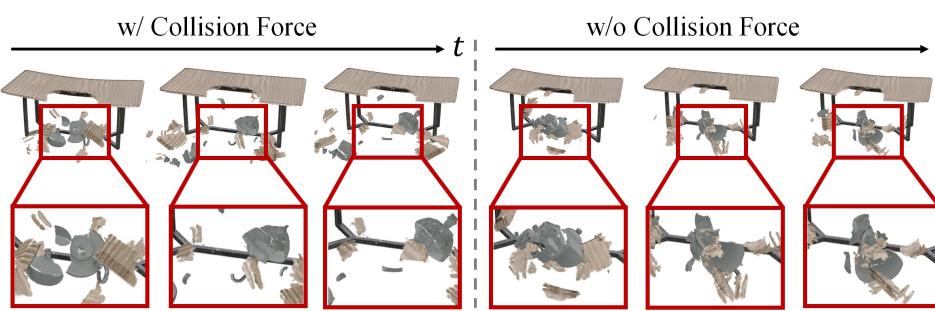


Figure 6: **Collision-MPM** effectively resolves the non-physical adhesion artifacts in multi-body collisions. As demonstrated in the red box regions, wood fragments from the fractured table exhibit natural separation behavior rather than adhering unnaturally to the teapot surface.

non-physical adhesion artifacts (highlighted in red boxes) in MLS-MPM simulations, whereas our method successfully eliminates these unrealistic phenomena through momentum-conserving interface forces derived from normalized mass distributions. Regarding the FPGO module, quantitative experiments in Figure 5 demonstrate its critical advantages. Both baseline methods - PhyGaussian and GIC - directly render using initially learned Gaussian attributes, resulting in various visual artifacts at fracture interfaces. In contrast, our dynamic attribute optimization strategy significantly enhances visual realism by reconstructing Gaussian properties through MVEE-based interpolation.

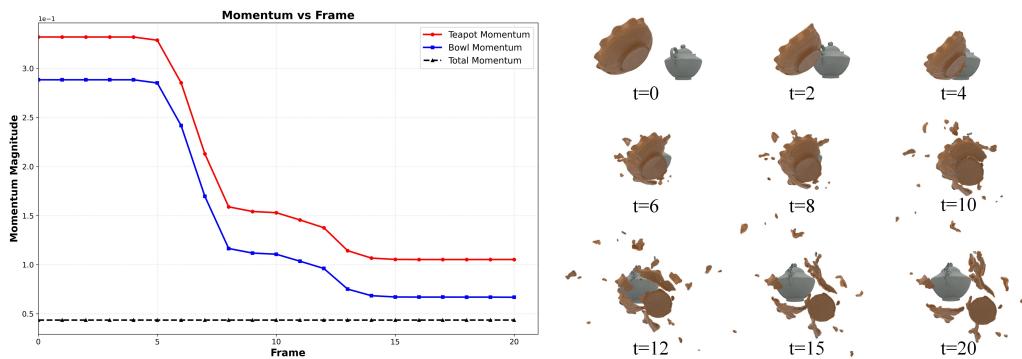


Figure 7: Momentum conservation during a collision between two objects. The total momentum (black curve) remains constant throughout the simulation, demonstrating strict adherence to the conservation law, while the individual momenta of the objects exchange during the impact.

Momentum Conservation Validation. To rigorously validate that our simulation framework strictly adheres to the law of momentum conservation, we designed a controlled experiment. This experiment involves two objects—a teapot and a bowl—propelled towards each other with initial velocities in a environment free from external influences such as gravity and friction. The figure 7 tracks the system’s evolution over time ($t = 0$ to $t = 20$). It visually demonstrates that while individual momenta change during the collision interval ($t \approx 5$ to $t \approx 15$), the “Total Momentum” curve remains constant in both magnitude and direction throughout the entire sequence. This constant total momentum can be clearly demonstrated from the figure that our simulation method accurately maintains the total linear momentum of the system.

Energy Stability Analysis. To quantitatively validate the concern regarding non-physical energy growth and numerical instability, we conduct a thorough energy evolution analysis throughout a representative simulation involving the collision and fracture of a teapot on a table surface. As illustrated in Figure 8, we track the kinetic, elastic, and gravitational potential energy components for both the teapot and the table individually, as well as for the combined system over the first 50 frames. The results demonstrate that the total energy of the system remains strictly bounded and does not exhibit any anomalous increase. Energy is transferred in a physically consistent manner: kinetic energy converts into elastic deformation energy upon impact, and part of it is dissipated through

fracture processes, while gravitational potential energy varies accordingly with object height. The smooth transitions and the absence of energy blow-up confirm that our method inherently prevents non-physical energy accumulation. This energy behavior serves as an effective unit test, validating that our simulation not only captures complex dynamic and fracture phenomena but also maintains numerical stability under severe contact and deformation conditions.

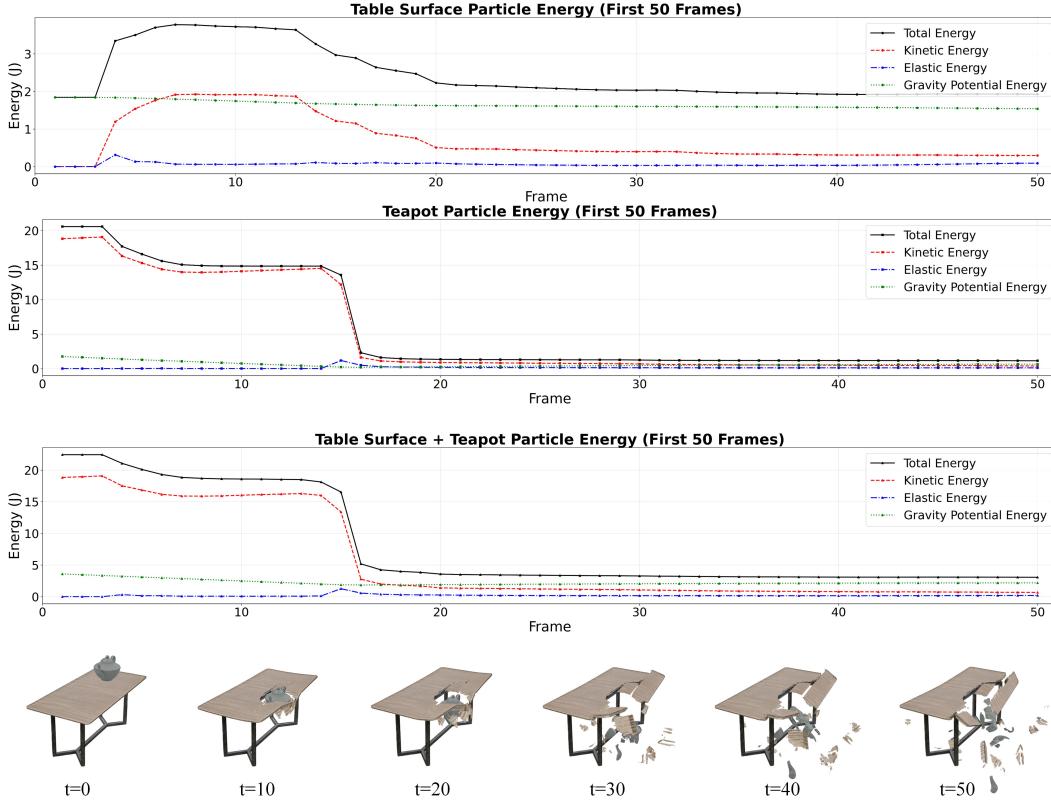


Figure 8: Energy evolution and qualitative visualization during the teapot-table collision and fracture simulation. The top and middle rows plot the kinetic, elastic, and gravitational potential energy components for the table surface, the teapot, and the combined system over the first 50 frames, showing bounded total energy without non-physical growth. The bottom row provides corresponding qualitative visualizations at key frames ($t=0, 10, 20, 30, 40, 50$), depicting the physical progression of the collision and fracture process that correlates with the energy transitions observed in the graphs.

5 CONCLUSION

This paper presents a unified framework for simulating and rendering extreme mechanical collisions with dynamic fracture effects. Our key contributions is the integration of two novel components: (1) an enhanced Collision-MPM formulation that introduces momentum-conserving interface forces to eliminate adhesion artifacts in multi-body collisions, and (2) a fracture-aware Gaussian optimization strategy that maintains visual fidelity during fracture propagation by dynamically reconstructing particle attributes. The framework demonstrates robust performance across various scenarios, including high-velocity impacts and heterogeneous material fractures. Qualitative and quantitative evaluations show significant improvements over existing methods in both physical accuracy and rendering quality.

Current limitations include computational demands that prevent real-time performance for complex scenes, and the need for manual parameter setting. Future directions will focus on two key aspects: (1) GPU optimization and adaptive time-stepping for faster computation, (2) development of learning-based approaches for automatic parameter estimation. These improvements will further enhance the framework’s practicality for industrial applications in virtual prototyping and visual effects production.

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648 **A APPENDIX**649 **A.1 MATERIAL PARAMETERS**

651 Table 1 summarizes the material parameters employed in our collision simulations, including
 652 Young’s modulus (E), Poisson’s ratio (ν), density, and non-associated flow rule parameters
 653 (α, β, ξ, M). These physically-based values were assigned to each object component to validate
 654 our framework’s capability in handling diverse material properties. The selected parameters reflect
 655 realistic material contrasts, enabling quantitative analysis of how mechanical properties—such as
 656 the stiffness variation between flexible leaves and rigid branches— influence fracture propagation
 657 patterns across all experiments.

658 **Table 1: Material Parameters**

660 Scene	661 E (MPa)	662 ν	663 Density	664 NACC (α, β, ξ, M)
662 Bowl	663 5×10^4	664 0.46	665 2	666 0.98, 0.5, 1, 2.36
663 Ficus leaf	664 8×10^4	665 0.39	666 0.6	667 0.94, 2, 3, 2.36
664 Ficus branch	665 1×10^6	666 0.39	667 5	668 0.94, 2, 3, 2.36
665 Ficus pot	666 2×10^4	667 0.39	668 2	669 0.98, 0.5, 2, 2.36
666 Teapot	667 5×10^5	668 0.46	669 5	670 0.98, 0.5, 1, 2.36
667 Table top	668 1.5×10^4	669 0.39	670 1	671 0.99, 0.5, 1, 2.36
668 Table leg	669 1×10^8	670 0.39	671 1000	672 0.94, 2, 3, 2.36

671 **A.2 EXPERIMENTS DETAIL**

672 Here, we provide additional details regarding the experiments. These include collisions between
 673 single object and rigid surface, collisions among multiple objects, as well as experiments on the
 674 interactions between individual objects and rigid surface under varying physical parameters. We
 675 also present an ablation study on the impact of fracture particle tracking.

676 **Input Resolution & Sampling:** The input resolution starts at 1024×1024 . For each collision
 677 object, we sample 200,000 surface points and 100,000 interior points.

678 **Simulation Setup:** Building upon Warp, the simulation is executed on an 18-core Intel Xeon Gold
 679 5220 CPU and an NVIDIA GeForce RTX 3090 GPU, achieving 100-frame sequences for each
 680 collisions scene.

681 **Implementation Details.** Following established practice in GIC (Cai et al., 2024), our simulation
 682 pipeline begins with isotropic Gaussian reconstruction for visual surface representation. For phys-
 683 ical discretization, we employ an SDF-based voxelization strategy: after constructing a volumetric
 684 grid and identifying interior voxels via SDF filtering, we perform uniform random sampling **inside**
 685 **each interior voxel** using a controllable density parameter N_v (samples per voxel). This preprocess-
 686 ing step decouples visual quality from simulation discretization while ensuring physically plausible
 687 material sampling. To evaluate the influence of the particle sampling density N_v on the simulation
 688 results, we set multiple sets of internal sampling parameters ($N_v = 0, 5, 10, 15, 20, 50, 100, 200$),
 689 and the visualization results are shown in Figure 12.

690 **A.3 METRIC COMPUTATION FOR FRACTURE ASSESSMENT**

692 To quantitatively evaluate the visual plausibility of fracture propagation in the absence of ground
 693 truth dynamic sequences, we designed a self-referencing assessment protocol using three estab-
 694 lished image quality metrics: PSNR, LPIPS, and FID. Our evaluation strategy focuses on measuring
 695 how faithfully each method maintains visual continuity during the fracture process. For each frac-
 696 ture event, we identify the last frame before fracture initiation (detected via the hardening parameter
 697 α) as the reference frame I_{ref} , representing the intact object’s appearance. We then analyze the sub-
 698 sequent 5 frames where fracture propagation becomes fully visible, ensuring this window captures
 699 critical fracture dynamics while maintaining visual comparability to the reference. The core prin-
 700 ciple underlying our metric design is that high-quality fracture rendering should appear as a coherent
 701 extension of the original material. Accordingly, we compute PSNR to measure pixel-level con-
 702 sistency with the pre-fracture state, LPIPS to assess perceptual similarity to the intact appearance,
 703 and FID to evaluate distributional similarity to pre-fracture rendering. Methods that produce visual

702 artifacts or unnatural fracture surfaces consequently exhibit large deviations from I_{ref} , resulting in
 703 degraded metric scores that reflect their reduced visual plausibility.

704 It should be noted, however, that this represents a pragmatic compromise in the absence of more
 705 principled evaluation methodologies. We regard the development of dedicated metrics for physics-
 706 based rendering as an important direction for future work.

707

A.4 MINIMAL-VOLUME ENCLOSING ELLIPSOID (MVEE)

708

A.4.1 PROBLEM DEFINITION

709 Given two spheres in 3D space with centers $\mathbf{C}_1, \mathbf{C}_2$ and radii r_1, r_2 , we aim to compute the **minimal-**
 710 **volume enclosing ellipsoid (MVEE)** of their intersection region.

711

Step 1: Intersection Conditions and Geometric Parameters

712

713 **Intersection Criteria.** The spheres intersect if:

$$714 |r_1 - r_2| < d < r_1 + r_2, \quad \text{where } d = \|\mathbf{C}_1 - \mathbf{C}_2\|$$

715 If $d \geq r_1 + r_2$, the spheres are disjoint; if $d \leq |r_1 - r_2|$, one sphere is entirely contained within the
 716 other.

717 **Key Geometric Properties.** The intersection region is a *lens-shaped* volume bounded by two
 718 spherical caps. Its properties include:

- 719 • **Symmetry axis:** The line connecting \mathbf{C}_1 and \mathbf{C}_2 (unit vector $\mathbf{u} = \frac{\mathbf{C}_2 - \mathbf{C}_1}{d}$).
- 720 • **Maximal width:** Perpendicular to \mathbf{u} , determined by the radius of the circle of intersection.

721

Step 2: Analytical Estimation of the Enclosing Ellipsoid

722 **Ellipsoid Center.** The center \mathbf{c} of the MVEE is approximated as a weighted midpoint along the
 723 symmetry axis:

$$724 \mathbf{c} = \mathbf{C}_1 + \left(\frac{r_1}{r_1 + r_2} \right) (\mathbf{C}_2 - \mathbf{C}_1)$$

725 This heuristic prioritizes the larger sphere's influence.

726 **Ellipsoid Axes.** The ellipsoid has three principal axes:

727 1. **Major axis (aligned with \mathbf{u}):**

$$728 \quad a = r_1 + r_2 - d$$

729 2. **Minor axes (perpendicular to \mathbf{u}):** Lengths $b = c$, given by the radius of the intersection
 730 circle:

$$731 \quad b = c = \sqrt{r_1^2 - \left(\frac{d^2 + r_1^2 - r_2^2}{2d} \right)^2}$$

732 **Orientation Matrix.** The ellipsoid's rotation matrix \mathbf{R} is constructed from the orthonormal basis:

$$733 \quad \mathbf{R} = [\mathbf{u} \ \mathbf{v} \ \mathbf{w}], \quad \text{where } \mathbf{v} \perp \mathbf{u}, \mathbf{w} = \mathbf{u} \times \mathbf{v}$$

734

A.5 COMPUTATIONAL EFFICIENCY

735 We analyze the computational overhead of our fracture tracking mechanism by comparing the per-
 736 frame rendering time for three core scenarios in main paper, with and without this feature enabled.
 737 As detailed in Table 2, the incorporation of fracture tracking introduces a moderate and consistent
 738 computational cost, increasing rendering time by approximately 20-80% across the scenes. This
 739 overhead is attributed to the additional steps of dynamically updating the Gaussian splatting model
 740 to reflect new fracture surfaces and collisions in each frame.

741

742 Table 2: Per-frame Gaussian rendering time.

743 Scene	744 w/o fracture tracking	745 w/ fracture tracking
746 Ficus	747 49.19 ms	748 62.78 ms
749 Teapot	750 6.13 ms	751 10.87 ms
752 Teapot & Table	753 27.15 ms	754 29.85 ms

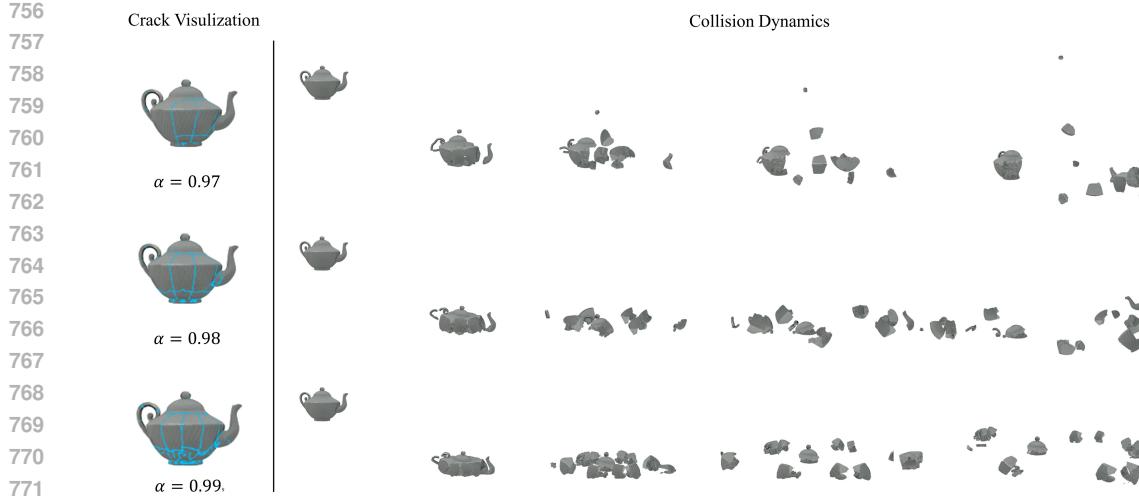


Figure 9: Collisions with three distinct initial hardening factors.

775 A.6 ABLATION STUDIES

776 **Effect of Fracture Particle Tracking and Gaussian Generation.** During the fracture process,
 777 some internal particles become exposed and visible. To achieve higher visual realism, we need to
 778 quickly locate these particles and generate corresponding Gaussian visual attributes. Based on the
 779 hardening tracking parameters (α), we can efficiently track relevant internal particles on the fracture
 780 surface and rapidly generate their Gaussian visual attributes through interpolation with neighboring
 781 external surface particles, enabling high-quality rendering. Additionally, a large number of isolated
 782 particles may be generated during collisions. We utilize the fracture particle tracking mechanism
 783 to exclude these fine particles from Gaussian rendering. This selective exclusion ensures that
 784 only particles meeting specific material integrity criteria participate in the final rendering, thereby
 785 maintaining the physical accuracy while improving computational efficiency.

786 **Single-Object Analysis.** To systematically evaluate our method’s sensitivity to material properties,
 787 we conduct extensive collision experiments on individual objects with varying plasticity parameters,
 788 as shown in Figure 9. Our results demonstrate that increasing the initial hardening factor systemati-
 789 cally enhances object fragmentation and promotes more extensive crack propagation, validating our
 790 method’s ability to capture material-dependent fracture behaviors.

791 **Multi-Object Interactions.** To assess performance in complex scenarios, we execute multiple chal-
 792 lenging multi-body interaction experiments, as shown in Figure 10. These include bowl-teapot col-
 793 lisions, bowl-table impacts, and ficus-table interactions. These experiments demonstrate our frame-
 794 work’s robustness in handling heterogeneous material compositions and complex contact dynamics
 795 across diverse object categories.

796 **Real-world Data Validation.** To further validate the generalizability of our approach, we extend our
 797 evaluation to real-world data from the DTU dataset. As shown in Figure 11 and 12, our pipeline suc-
 798 cessfully performs Gaussian reconstruction and initialization from real images, followed by physi-
 799 cally plausible fracture simulation. These results demonstrate our method’s robustness when applied
 800 to real-world captured data, confirming its practical applicability beyond synthetic environments.

801 **Analysis of sampling point density.** To assess the influence of particle sampling density on
 802 simulation outcomes, we conducted a systematic sweep of internal sampling parameters ($N_v =$
 803 $0, 5, 10, 15, 20, 50, 100, 200$). Using the highest density configuration ($N_v = 200$) as a convergence
 804 reference, we evaluated trajectory deviations across all sampling conditions (Figure 12 and Table 3).
 805 The results reveal a fundamental characteristic of MPM simulations: particle sampling density in-
 806 trinsically influences the resolved mechanical response. As shown in our quantitative analysis and
 807 qualitative visualizations (Figure 12 and Table 3), variations in sampling density naturally lead to
 808 differences in fracture patterns and collision dynamics—this is an inherent property of particle-based
 809 methods where discretization density affects solution convergence.

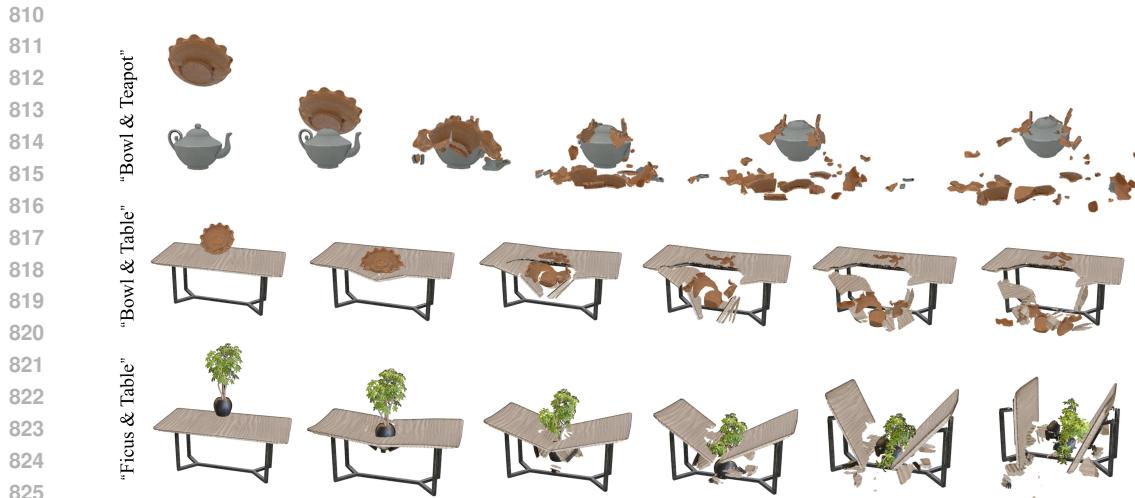


Figure 10: Collisions between different objects: "Bowl & Teapot", "Bowl & Table", "Ficus & Table".



Figure 11: Real-world data validation of our method: results for three different real-world objects.

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Interior Point Density	X-coordinate Error	Y-coordinate Error	Z-coordinate Error
0	0.1998	0.1479	0.0535
5	0.1936	0.1433	0.0474
10	0.1884	0.1360	0.0445
15	0.1772	0.1279	0.0435
20	0.1786	0.1224	0.0412
50	0.1713	0.1191	0.0337
100	0.0935	0.0749	0.0315

Table 3: Coordinate errors at different interior point densities.

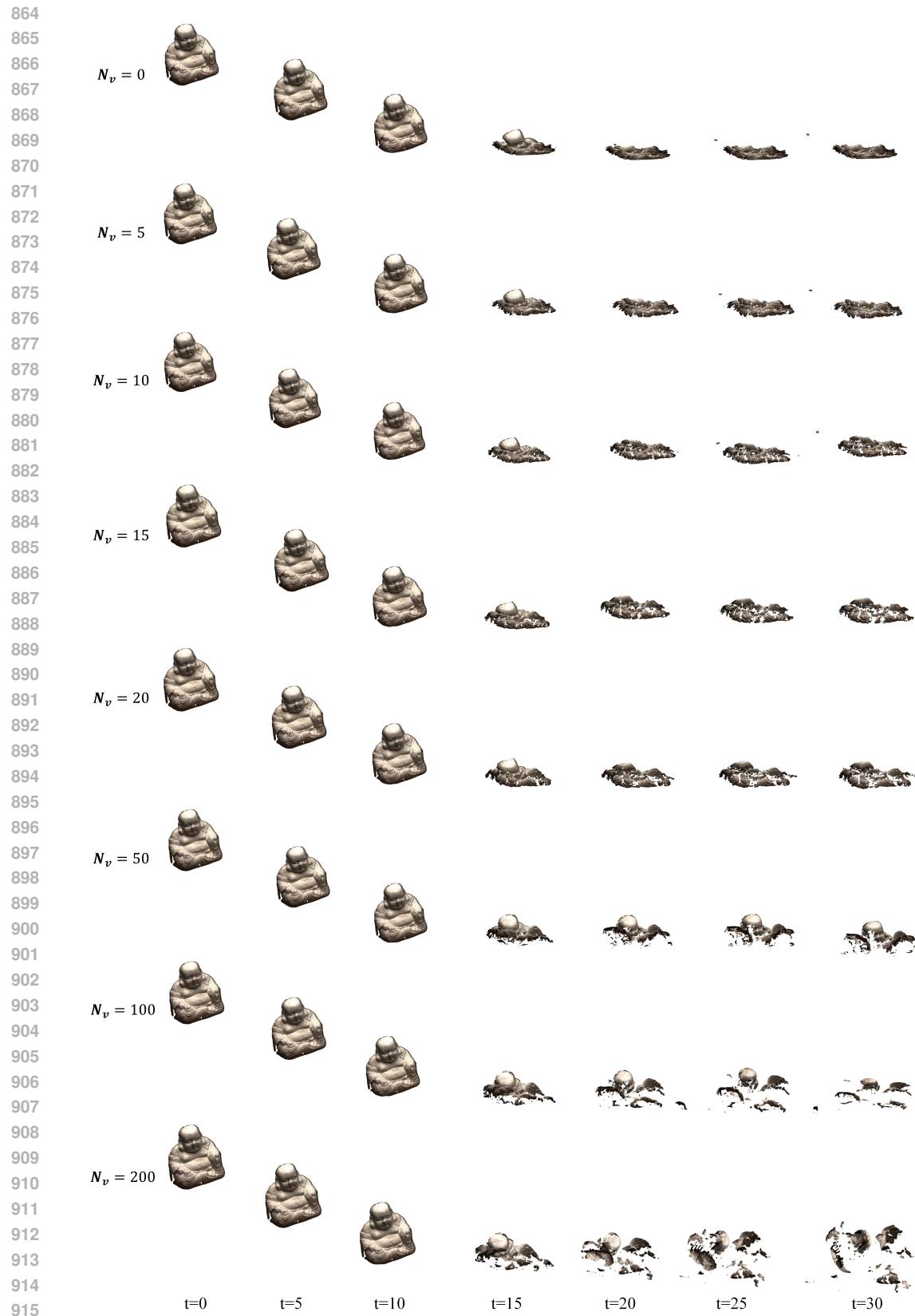


Figure 12: Visual comparison of simulation and rendering results under different internal sampling densities (N_v) across three representative timesteps.