A composable machine-learning approach for steady-state simulations on high-resolution grids

Abstract

In this paper we show that our Machine Learning (ML) approach, CoMLSim (**Co**mposable **M**achine **L**earning **Sim**ulator), can simulate PDEs on highly-resolved grids with higher accuracy and generalization to out-of-distribution source terms and geometries than traditional ML baselines. Our unique approach combines key principles of traditional PDE solvers with local-learning and low-dimensional manifold techniques to iteratively simulate PDEs on large computational domains. The proposed approach is validated on more than 5 steady-state PDEs across different PDE conditions on highly-resolved grids and comparisons are made with the commercial solver, Ansys Fluent as well 4 other state-of-the-art ML methods. The numerical experiments show that our approach outperforms ML baselines in terms of 1) accuracy across quantitative metrics and 2) generalization to out-of-distribution conditions as well as mesh resolutions. Additionally, we provide results of conducting a large number of ablations experiments to highlight components of our approach that strongly influence the results. We conclude that our local-learning and iterative-inferencing approach reduces the challenge of generalization that most ML models face.

1 Introduction

Engineering simulations utilize solutions of partial differential equations (PDEs) to model various complex physical processes like weather prediction, combustion in a car engine, thermal cooling on electronic chips etc., where the complexity in physics is driven by several factors such as Reynolds number, heat source, geometry, etc. Traditional PDE solvers use discretization techniques to approximate PDEs on discrete computational domains and combine them with linear and non-linear equation solvers to compute PDE solutions. The PDE solutions are dependent on several conditions imposed on the PDE, such as geometry and boundary conditions of the computational domain as well as source terms such as heat generation, buoyancy etc. Moreover, many applications in engineering require highly-resolved computational domains to accurately approximate the PDEs solutions as well as to capture the intricacies of geometry and source terms imposed on the PDEs. As a result, traditional PDE solvers, although accurate and generalizable across various PDE conditions, can be computationally slow, especially in the presence of complicated physics and large computational domains. In our work, we specifically aim to increase the speed of engineering simulations using ML techniques on complex PDE conditions used for enterprise simulations.

The idea of using Machine Learning (ML) with PDEs is not unique and has been explored for several decades [1, 2]. ML approaches are computationally fast, but they fall short in terms of accuracy, generalization to a wide range and out-of-distribution PDE conditions and in their ability to scale to highly-resolved computational grids, when compared to traditional PDE solvers. A few shortcomings of the current methods are outlined below and subsequently are motivations for this work.

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Figure 1: CoMLSim performance vs Ansys Fluent for a high Reynolds number, steady-state turbulent flow over a car (never seen during training) at 90 mph. The vectors and contours predictions match reasonably with MAE of 0.0145. High flow gradients that affect the wake pattern are resolved accurately and comparable to Fluent.

1. ML approaches use static inferencing to predict PDE solutions as a function of the PDE conditions. In many cases, solutions and conditions have high-dimensional and sparse representations, which are challenging to generalize with such black-boxed static inferencing.

2. Most engineering applications require highly-resolved computational grids to capture detailed solution features, such as hot spots on an electronic chip surface. This adversely impacts training in terms of GPU memory, computational time and data requirements.

3. Most ML approaches fail to use powerful information from traditional PDE solvers related to solver methods, numerical discretization etc.

In this work, we introduce a novel ML approach, **Co**mposable **M**achine **L**earning **Sim**ulator (CoML-Sim) to work with high-resolution grids without compromising accuracy. To achieve that, our method decomposes highly-resolved computational grids into smaller local subdomains with Cartesian grids. PDE solutions and conditions on each subdomain are represented by lower-dimensional latent vectors determined using autoencoders. The latent vectors of solutions corresponding to conditions are predicted in each subdomain and stitched together to maintain local consistency, analogous to flux conservation in traditional solvers. In this paper, we describe our method (shown in fig. 2) for building CoMLSim for different PDEs, showcase different experiments to compare accuracy and generalizability with different ML methods and provide an extensive ablation study to understand the impact of CoMLSim's different components on the results.

Significant contributions of this work:

1. The CoMLSim approach combines traditional PDE solver strategies such as domain discretization, flux conservation, solution methodologies etc. with ML techniques to accurately model numerical simulations.

2. Our approach operates on local subdomains and solves PDEs in a low-dimensional space. This enables generalizing to out-of-distribution PDE conditions and scaling to higher-dimensional grids.

Figure 2: CoMLSim solution algorithm

3. The iterative inferencing algorithm is unsupervised and allows for coupling with traditional PDE solvers.

Related Works: The use of ML for solving PDEs has gained tremendous traction in the past few years. Much of the research has focused on improving neural network architectures and optimization techniques to enable generalizable and accurate learning of PDE solutions. More recently, there has been a lot of focus on learning PDE operators with neural networks (NNs) [3, 4, 5, 6, 7, 8, 9]. The neural operators are trained on high-fidelity solutions generated by traditional PDE solvers on a computational grid of specific resolution and do not require any knowledge of the PDE. However, accuracy and out-of-distribution generalizability deteriorates as the resolution of computational grid increases (for example, 2048^2 in 2-D or 256^3 in 3-D). Furthermore, these methods are limited by an upper cap on GPU memory.

A different research direction focuses on training neural networks with physics constrained optimization [10, 11]. These method use Automatic Differentiation (AD) [12] to compute PDE derivatives. The physics-based approaches have been extended to solve complicated PDEs representing complex physics [13, 14, 15, 16, 17, 18, 19, 20, 21, 22, 23, 24, 25, 26, 27, 28, 29, 30]. More recently, alternate approaches that use discretization techniques using higher order derivatives and specialize numerical schemes to compute derivatives have shown to provide better regularization for faster convergence [31, 32, 33, 34]. However, the use of optimization techniques to solve PDEs, although accurate, has proved to be extremely slow as compared to traditional solvers and hence, non-scalable to high resolution meshes.

Domain decomposition techniques have been successfully used in the context of ML applied to PDEs [35]. These techniques enable local learning from smaller restricted domains and have proved to accelerate learning of neural networks. In [36, 35], the authors proposed a hybrid method combining ML with domain decomposition to reduce the computational cost of finite element solvers. Similarly, domain decomposition has been extensively used in the context of physics informed neural networks to reduce training costs by enabling distributed training on multiple GPUs [37, 38]. Lu et al. [39] and Wang et al. [40] learn on localized domains but infer on larger computational domains using a stitching process. Bar-Sinai et al. [41] and Kochkov et al. [42] learn coefficients of numerical discretization schemes from high fidelity data, which is sub-sampled on coarse grids. Beatson et al. [43] learn surrogate models for smaller components to allow for cheaper simulations. Greenfeld et al. [44] learn local prolongation operators from discretization matrices, to improve the rate of convergence of the multigrid linear solvers. Additionally, other ML methods improve the learning of ML models by compressing PDE solutions on to lower-dimensional manifolds. This has shown to improve accuracy and generalization capability of neural networks [45, 46, 47, 48, 49, 50]. In this work, domain decomposition is combined with latent space learning is employed to accurately represent solutions on local subdomains and allow scalability to high-resolution computational grids. We use a Cartesian grid to discretize each subdomain. However, Graph Neural Networks (GNNs) [51, 52, 53, 54, 55] as well as FEM use other discretizations such as triangle or polygonal meshes. These ideas can be used to handle unstructured meshes in subdomains in future extensions of this work.

2 CoMLSim approach

2.1 Similarities with traditional PDE solvers

Consider a set of coupled steady-state PDEs with n solution variables. For the sake of notation simplicity, we take $n = 2$, such that $u(x, y, z)$ and $v(x, y, z)$ are defined on a computational domain Ω with boundary conditions specified on the boundary of the computational domain, Ω_b. The coupled PDEs are defined as follows:

$$L_1(u, v) - F_1 = 0; L_2(u, v) - F_2 = 0 \tag{1}$$

where, L_1, L_2 denote PDE operators, F_1, F_2 represent PDE source terms and $u = u_b, v = v_b$ for $\Omega = \Omega_b$. The PDE operators can vary for different PDEs. For example, in a non-linear PDE such as the unsteady, incompressible Navier-Stokes equation the operator, $L = \vec{a}.\vec{\nabla} - \vec{\nabla}.\vec{\nabla}$

Traditional FVM/FDM based PDE solvers solve PDEs in Eq. 1 by computing solutions variables, u, v, and their linear/non-linear derivatives on a discretized computational domain. Iterative solution

algorithms are used to conserve fluxes between neighboring computational elements and determine consistent PDE solutions over the entire domain at convergence. The CoMLSim approach is designed to perform similar operations but at the level of subdomains (n^3 elements) using ML techniques.

2.2 Solution algorithm for steady-state PDEs

The solution algorithm of CoMLSim approach is shown in Fig. 2 and described in detail in Alg. 1. Similar to traditional solvers, the CoMLSim approach discretizes the computational domain into smaller subdomains contain n^p computational elements, where p refers to spatial dimensionality and n is predetermined. Each subdomain has a constant physical size and represents both PDE solutions and conditions. For example, a subdomain cutting across the cylinder in Fig. 2 represents the geometry and boundary conditions of the part-cylinder as well as the corresponding solution. In a steady-state problem, the solutions on the domain are initialized either with uniform solutions or they are generated from coarse-grid PDE solvers. The initial solution $\vec{p} = (u(\vec{x}), v(\vec{x}))$ is shown with randomly oriented flow vectors in Fig. 2. Pretrained encoders $e_{\vec{p}}, e_g, e_b, e_s$ are used to encode the initial solutions as well as user-specified PDE conditions into lower-dimensional latent vectors, $\eta_{\vec{p}}$ corresponding to PDE solutions and η_g, η_b, η_s corresponding to geometry, boundary conditions and source terms, respectively. Solution and condition latent vectors on groups of neighboring subdomains are concatenated together evaluated using a pre-trained flux conservation autoencoder (Θ) to get a new set of solution latent vectors ($\eta'_{\vec{p}}$) on each subdomain, which are more locally consistent than the original vectors. The new solution latent vectors are iteratively passed through Θ to improve their local consistency and the iteration is stopped when the L_2 norm of change in solution latent vectors meets a specified tolerance, otherwise the iteration continues with the updated latent vectors. The latent vectors of the PDE conditions are not updated and help in steering the solution latent vectors to an equilibrium state that is decoded to PDE solutions using pretrained decoders ($g_{\vec{p}}$) on the computational domain. The converged solution in Fig. 2 is represented with flow vectors that are locally consistent with neighboring subdomains.The iterative procedure used in the CoMLSim approach can be implemented using several linear or non-linear equation solvers, such as Fixed point iterations [56], Gauss Seidel, Newton's method etc., that are used in commercial PDE solvers. Out of those we have explored Point Jacobi and Gauss Seidel, which are described below. The use of autoencoders for compressed solution and condition representation in tandem with iterative inferencing solution algorithm is inspired from [50, 31] but the main difference in this work is that the solution procedure is carried out on local subdomains as opposed to entire computational domains to align with principles used in traditional PDE solvers.

Algorithm 1: Solution methodology of CoMLSim approach

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1 Domain Decomposition: Computational domain $\Omega \rightarrow$ Subdomains Ω_c
2 Initialize solution on all Ω_c: $\vec{p}(\vec{x}) = 0.0$ for all $\vec{x} \in \Omega_c$
3 Encode solutions on Ω_c: $\eta_{\vec{p}} = e_u(\vec{p}(\Omega_c))$
4 Encode conditions on Ω_c:
5 $\eta_g = e_g(g(\Omega_c)), \eta_b = e_b(b(\Omega_c)), \eta_s = e_s(s(\Omega_c))$
6 Set convergence tolerance: $\epsilon_t = 1e^{-8}$
7 **while** $\epsilon > \epsilon_t$ **do**
8   **for** $\Omega_c \in \Omega$ **do**
9   Gather neighbors of Ω_c: $\Omega_{nb} = [\Omega_c, \Omega_{left}, \Omega_{right}, ...]$
10         $\eta'_{\vec{p}} = \Theta(\eta_{\vec{p}}^{nb}, \eta_b^{nb}, \eta_g^{nb}, \eta_s^{nb})$
11   Compute L_2 norm: $\epsilon = ||\eta_{\vec{p}} - \eta'_{\vec{p}}||_2^2$
12   Update: $\eta_{\vec{p}} \leftarrow \eta'_{\vec{p}}$ for all $\Omega_c \in \Omega$
13 Decode PDE solution on all Ω_c: $\vec{p} = g_{\vec{p}}(\eta_{\vec{p}\Omega_c})$

Stability: The iterative inferencing approach proposed in this work is similar to traditional approaches used in solving system of linear equations. Our algorithm uses fixed point iterations to solve Eq. 2

$$\eta'_{\vec{p}} = \Theta(\eta_{\vec{p}}^{nb}, \eta_b^{nb}, \eta_g^{nb}, \eta_s^{nb}) \tag{2}$$

where, Θ corresponds to the weights of flux conservation network, $\eta_{()}$ refers to the latent vectors and nb refers to neighboring subdomains defined below in Eq. 3.

$$\eta_{\vec{p}}^{nb} = \left[\eta_{\vec{p}}^s, \eta_{\vec{p}}^{left}, \eta_{\vec{p}}^{right}, \eta_{\vec{p}}^{top}, \eta_{\vec{p}}^{bottom}\right] \tag{3}$$

150 Similar to linear systems, the stability of our iterative inferencing approach is governed by the
151 condition number [57] defined in Eq. 4

$$||\frac{\partial \Theta(\eta_{\vec{p}}, \eta_b, \eta_g, \eta_s)}{\partial \eta_{\vec{p}}}|| < 1 \tag{4}$$

152 **Solution algorithms: Point Jacobi vs Gauss Seidel:** Since the algorithm loops over all the sub-
153 domains in a specified order, while updating the solution encoding of a subdomain c, there are
154 subdomains in the neighborhood that already have an updated solution encoding. The Gauss Seidel
155 method uses these updated solution encodings ($\eta'_{\vec{p}}$) on neighboring subdomains to update the solution
156 encoding on subdomain, c. On the other hand, Point Jacobi method does not make use of this
157 and hence can be easily vectorized for significantly faster computation. The implementation has a
158 remarkable similarity with how traditional solvers use these methods [58] and leaves the door open to
159 use other non-linear optimizers with physics-based constraints.

2.3 Neural network components in CoMLSim

161 Our algorithm employs two types of autoencoders, CNN autoencoders to establish a lower-
162 dimensional representation of PDE solutions and conditions from local subdomains grids and FCNN
163 Autoencoders for flux conservation, where the goal is to learn a reduced representation of solution
164 and condition latent vectors in a local neighborhood. Let us consider an example of solving the
165 Laplace equation, $\nabla^2 u = 0$ in 2D for arbitrarily shaped computational domains, Ω_g. In this example,
166 the CoMLSim algorithm will require 3 autoencoders, described in Eq. 5, to learn lower-dimensional
167 representations on local subdomains.

$$\left.\begin{array}{c} u' = S(u, \eta_u) \\ g' = G(g, \eta_g) \\ \eta'_{nb} = \Theta(\eta_{nb}, \zeta) \end{array}\right\} \tag{5}$$

168 where, u, η_u, S refers to the solution, its latent vector and the weights of the autoencoder, respectively,
169 g, η_g, G refers to a representation of geometry, such as Signed Distance Fields (SDF) [59], its latent
170 vector and the weights of the autoencoder, Θ refers to the weights of flux conservation autoencoder and
171 η_{nb} represents a set of concatenated latent vectors, η_u and η_g on a group of neighboring subdomains.
172 All the autoencoders are trained with samples of PDE solutions generated for the same use case. In
173 the case of coupled PDEs, a single autoencoder is trained for all solution variables. Additional details
174 related to the autoencoder networks are provided in the supplementary materials.

175 **Why Autoencoders?:** Solutions to classical PDEs such as the Laplace equation can be represented
176 by homogeneous solutions as follows:

$$\phi(x, y) = a_0 + a_1 x + a_2 y + a_3(x^2 - y^2) + a_4(2xy) + ... \tag{6}$$

177 where, $\vec{A} = a_0, a_1, a_2, ..., a_n$ are constant coefficients that can be used to reconstruct the PDE solution
178 on any local subdomain. $\vec{A}$ can be considered as a compressed encoding of the Laplace solutions.
179 Since, it is not possible to explicitly derive such compressed encodings for other high dimensional and
180 non-linear PDEs, the CoMLSim approach relies on autoencoders to compute them. It is known that
181 non-linear autoencoders with good compression ratios can learn powerful non-linear generalizations
182 [60, 61, 62]. Autoencoders also have great denoising abilities, which improve robustness and stability,
183 when used in iterative settings [50]. In their paper, Park and Lee [63] demonstrate that the latent
184 manifold established by a trained autoencoder is stable for varying intensities of Gaussian noise. This
185 quality of autoencoders is useful for the convergence of the iterative inferencing approach.

3 Experiments

187 In this section as well as in the supplementary materials, we consider a number of use cases with
188 varying degrees of difficulty resulting from the PDE formulation as well as source terms, geometry
189 and boundary conditions. The PDEs have applications in fluid mechanics, structural mechanics and
190 semiconductor simulations. The source code demonstrating an experiment and the datasets used in
191 our experiments and analysis can be made available upon acceptance.

192 **Details of experiments:** Here we provide some details about the experiments. Additional details and
193 experiments on Laplace and Darcy equation may be found in the supplementary materials.

1. **2D Poisson equation:** The Poisson's equation, shown in Eq. 7, is very popular in engineering simulations, for example, chip temperature prediction, pressure equation in fluids etc.

$$\nabla^2 u = f \tag{7}$$

where, u is the solution variable and f is the source term. This PDE is solved on a 1024x1024 grid to resolve the high-frequency features of the source term. The source, f, is sampled from a Gaussian mixture model shown in Eq. 8.

$$\sum_{j=0}^{1024} \sum_{j=0}^{1024} f_{i,j} = \sum_{j=0}^{1024} \sum_{j=0}^{1024} \sum_{k=0}^{30} A_k \exp\left(-\frac{x - \mu_{x,k}}{\sigma_{x,k}} - \frac{y - \mu_{y,k}}{\sigma_{y,k}} \right) \tag{8}$$

where, x, y correspond to the grid coordinates. A_k randomly assumes either 0 or 1 to vary the number of active Gaussians in the model. μ_x, μ_y and σ_x, σ_y are the mean and standard deviations of Gaussians in x and y directions, respectively. The means and standard deviations vary randomly between 0 to 1 and 0.001 to 0.01, respectively. The smaller magnitude of standard deviation results in hot spots that require highly-resolved grids. In this case, 256 training and 100 testing solutions are generated using Ansys Fluent.

2. **2D non-linear coupled Poisson equation:** The coupled non-linear Poisson's equation is shown below in Eq. 9. u, v are the solution variables and f is the source term, similar to the description in Eq. 8. This PDE has applications in reactive flow simulations.

$$\nabla^2 u = f - u^2; \qquad \nabla^2 v = \frac{1}{u^2 + \epsilon} - v^2 \tag{9}$$

The data generation is similar to Experiment 2 with a slight difference that the source term is less stiff with a standard deviation that varies between 0.005 to 0.05.

3. **3D Reynolds-Averaged Navier-Stokes external flow:** This use case consists of a 3-D channel flow with resolution 304x64x64 at high Reynolds number over arbitrarily shaped objects. The corresponding PDEs are presented in supplementary materials or may be referred from [64]. The characteristics of flow generated on the downstream has important applications in the design of automobiles and airplanes. In this case, we use 5 primitive 3-D geometries namely, cylinder, cuboid, trapezoid, airfoil wing and wedge and their random combinations and rotational augmentations to create 150 training geometries and 50 testing geometries, which are solved with Ansys Fluent to generate the data. Out-of-distribution testing is carried out on simplistic automobile geometries as shown in Figure 1.

4. **3D chip cooling with Natural convection:** In this case, we extend the complexity of Experiments 1, 2 and 3 to a industrial 3-D case of chip cooling with natural convection. In this problem, the computational grid (128^3 resolution) consists of a chip that is subjected to the powermap specified by Eq. 8. The heating of the chip results in generation of velocity in the fluid and at steady-state, a balance is achieved. We use Ansys Fluent to generate 300 training and 100 testing solutions for velocity, pressure, temperature for arbitrary sources.

4   Results and Ablation studies

In this section, we provide a variety of results for the experiments outlined in Section 3. Additionally, we conduct thorough ablation studies to understand the different components of the CoMLSim approach in more detail. We use the 2-D Linear Poisson's equation experiment for these studies unless otherwise specified. Details regarding the CoMLSim set up for each experiment, baseline network architectures as well as additional results corresponding to contour and line plots and comparisons across other metrics are provided in supplementary materials.

4.1   Comparison with Ansys Fluent and other ML baselines

The CoMLSim is compared with other ML baselines namely, UNet [65], FNO [8] and DeepONet [7] for the experiments outlined in Section 3. We compare the mean absolute error with respect to Ansys Fluent and averaged over the unseen testing cases and all solution variables. In the case of 3-D chip cooling we report the L_∞ norm of temperature as this metric is more suited to this industrial application. It may be observed from Table 1 that the CoMLSim approach performs better than other

Figure 3: (A): Error vs No. of Gaussians for linear Poisson's; (B): Error vs No. of Gaussians for non-linear Poisson's; (C): Error vs grid resolution for linear Poisson's; (D): MAE and convergence iterations vs flux conservation bottleneck size

ML baselines. All the ML methods perform better for the non-linear Poisson's case as compared to the linear Poisson's because the source term is less stiff but the CoMLSim approach does a better job in modeling the non-linearity arising due solution coupling.

Table 1: Comparison with baselines

Experiment	Metric	CoMLSim	UNet	FNO	DeepONet	FCNN
2-D Linear Poisson's	L_1	0.011	0.132	0.031	0.061	0.267
2-D non linear Poisson's	L_1	0.0053	0.0877	0.0278	0.527	0.172
3-D NS external flow	L_1	0.012	0.0625	0.038	0.81	0.125
3-D chip cooling	L_∞	15.2	95.21	60.836	45.27	192.7

4.2 Assessment of generalizability

In this section, we assess the generalizability of our method specifically with respect to out-of-distribution PDE conditions and scaling to higher resolution grids. The results are compared with Ansys Fluent as well as other ML baselines discussed in section 4.1.

Out-of-distribution source term: The source term for linear and non-linear Poisson equations are sampled from a Gaussian mixture model described in Eq. 8, where the number of Gaussian is randomly chosen between 0 and 30. In this experiment, we evaluate the generalizability of our approach for source terms with exactly 30, 40, 50 and 60 number of Gaussians, corresponding to out-of-distribution for the training data distributions. For each case, the results averaged over 10 testing samples are compared with Ansys Fluent. It may be observed from 3A and B, that the accuracy of all ML approaches decreases as you move further away from the source term distribution. However, the CoMLSim approach performs significantly better than other ML baselines and the accuracy is reasonable even for the case with 60 Gaussian mixture model, which is substantially different from the training distribution.

Figure 4: Example of out-of-distribution source (60 Gaussians) comparison for non-linear Poisson's

Increasing mesh resolution: Next, we evaluate the performance of CoMLSim for higher grid resolutions and compare it with other ML baselines. In this experiment, we compute solutions for the 100 source term conditions in the test set but on 4 different mesh resolutions, 1024^2, 2048^2, 4096^2, 8192^2, respectively. The mean and standard deviation of Gaussians in the testing set are proportionally increased with the mesh resolution to ensure that the problem definition does not change. It is important to note that none of the networks are retrained for higher resolution grids. It maybe observed from table 3 that the CoMLSim continues to scale to higher resolution with similar accuracy as the original mesh resolution. The slight drop in accuracy due to increase in mesh resolution is attributed to accumulation of error caused by the flux conservation networks resulting from the increase in number of subdomains. Other ML baselines such as UNet and FNO scale up to a grid resolution of 2048^2 but cannot evaluate beyond that due to GPU memory constraints. DeepONet has a point-wise inference and can also scale to higher-resolution grids but its accuracy is lower than CoMLSim.

Out-of-distribution geometries: Next, we evaluate the performance of CoMLSim on 2 unseen car models presented at different 3-D angles of rotation to the external flow. It should be noted that these geometries are more complicated than the primitive objects considered for training our approach 3. The mean absolute error with respect to Ansys Fluent across all solution variables are 0.0145, 0.06924, 0.04175, 0.91 and 0.13725 for CoMLSim, UNet, FNO, DeepONet and FCNN respectively. It may be observed that CoMLSim significantly outperforms these baselines. Additional results and analysis are provided in supplementary materials.

4.3 Analysis of Subdomain size

In this experiment, we analyze the effect of subdomain resolution on accuracy and computational cost of the CoMLSim approach. We train 3 instances at resolutions of 32^2, 64^2 and 128^2. The compression ratio in autoencoders is kept the same for the different subdomain resolutions. The mean absolute error on testing set is 0.015, 0.011 and 0.029, respectively. Additionally, the iterations required to convergence are 420, 162 and 89, respectively. The convergence history is shown in Fig. 5A. The 128^2 subdomain resolution has a lower accuracy because it is challenging to train accurate autoencoders as bigger subdomains capture a large amount of information. On the other hand, the computational cost is the lowest for the 128^2 resolution because the number of subdomains in the entire computational domain are significantly less and hence, the solution algorithm converges faster.

4.4 Flux conservation bottleneck layer size

The flux conservation autoencoder as it is the primary workhorse of the CoMLSim solution algorithm. As a result, we evaluate the effect of the bottleneck layer size on the accuracy and computational speed. Each instance of CoMLSim with different bottleneck size is tested on 100 unseen test cases. We verify that all models satisfy the stability criterion specified in Eq. 3. It may be observed from Figure 3D that as the compression ratio of the flux conservation autoencoder decreases, it begins to overfit and the testing error as well as the number of convergence iterations and computational time significantly increase. On the other hand, if the compression ratio is too large the testing error increases because the autoencoders underfit. In alignment with the collective intuition about autoencoders, there exists an optimum bottleneck size compression ratio where the best testing error is obtained for small computational times.

4.5 Robustness and stability

A long standing challenge in the field of numerical simulation is to guarantee the stability and convergence of non-linear PDE solvers. However, we believe that the denoising capability of autoencoders [66, 60, 67, 68, 69] used in our iterative solution algorithm presents a unique benefit, irrespective of the choice of initial conditions. Here, we empirically demonstrate the robustness and stability of our approach.

Figure 5: (A): Convergence history for varying subdomain sizes; (B): Convergence history after coupling with traditional PDE solvers; (C, D): Robustness and stability of CoMLSim approach; (E): Point Jacobi vs Gauss Seidel convergence

In scenario A, we randomly sample 25 initial solutions from a
uniform distribution and in scenario B we sample from 6 differ-
ent distributions, namely Uniform, Gumbel, Laplace, Gamma,
Normal, Logistic etc. The mean convergence trajectory is plot-
ted in Fig. 5C and 5D show that the log of L_2 norm of the convergence error falls below an acceptable
tolerance for all cases and demonstrates the stability of our approach.

4.6 CoMLSim solution methods: Point Jacobi vs Gauss Seidel

Next, we compare the performance of the Point Jacobi and Gauss Seidel implementations in the
CoMLSim approach across 3 metrics, mean absolute error, computational cost and iterations to
converge. Both methods have a similar prediction accuracy with a mean absolute error of 0.013 and
0.017, respectively but the Point Jacobi method is computationally cheaper requiring 4.3 seconds
to converge as compared to 72 seconds of Gauss Seidel. It may be observed from Fig. 5E that the
Gauss Seidel method takes significantly fewer convergence iterations, but the Point Jacobi method is
faster because it can be efficiently vectorized on a GPU.

4.7 Coupling with traditional PDE solvers and super-resolution

In this experiment, we analyze the potential of coupling the CoMLSim approach with a traditional
PDE solver such as Ansys Fluent and subsequently demonstrate how our approach inherently
possesses the capability to perform super-resolution. We mainly carry out 5 experiments, where
we start from a coarse grid solution for the 100 test samples of linear Poisson's equation solved by
Ansys Fluent on a 64x64 grid (256x smaller than the fine resolution), add random noise to the coarse
initial solution with varying amplitudes (10%, 25% and 50%), and use that as an initial solution in
the CoMLSim approach. The convergence trajectory in Fig. 5B shows that the case with coarse grid
initialization converges the fastest, followed by initialization with noise and zero initialization. All
solutions converge to the same accuracy.

4.8 Computational speed

We observe that the CoMLSim approach is about 40-50x faster as compared to commercial steady-
state PDE solvers such as Ansys Fluent for the same mesh resolution in all the experiments presented
in this work. In comparison to the ML baselines, our approach is expected to be slightly slower
because it adopts an iterative inferencing approach. But, this is compensated by solution accu-
racy, generalizability and robustness on high-resolution grids. A detailed analysis is provided in
supplementary materials.

5 Conclusion

In this work, we introduce the CoMLSim approach, which is an unsupervised, low-dimensional and
local machine learning approach for solving PDE on highly-resolved grids and generalizing across
a wide range of PDE conditions. Our approach is inspired from strategies employed in traditional
PDE solvers and adopts iterative inferencing. The proposed approach is demonstrated to predict
accurate solutions for a range of PDEs, generalize reasonably across geometries, source terms and
BCs. Moreover, it scales to high resolution domains.

Broader impact, future work & limitations Although the proposed ML-model can generalize to
out-of-distribution geometries, source terms and BCs, but like other ML approaches, extrapolation
to any and all PDE conditions that are significantly different from the original distribution still
remains a challenge. However, this work takes a big step towards laying down the framework on
how truly generalizable ML-based solvers can be developed. In future, we would like to address
these challenges of generalizability and scalability by training autoencoders on random, application
agnostic PDE solutions and enforcing PDE-based constraints in the iterative inferencing procedure.
Future work will also investigate the potential for hybrid solvers and extensions to transient PDEs
and inverse problems. Finally, we will also address extension of the current approach to unstructured
meshes, which is a current limitation.

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A Appendix

Optionally include extra information (complete proofs, additional experiments and plots) in the appendix. This section will often be part of the supplemental material.