Extending Lagrangian and Hamiltonian Neural Networks with Differentiable Contact Models

Abstract

The incorporation of appropriate inductive bias plays a critical role in learning dynamics from data. A growing body of work has been exploring ways to enforce energy conservation in the learned dynamics by encoding Lagrangian or Hamiltonian dynamics into the neural network architecture. These existing approaches are based on differential equations, which do not allow discontinuity in the states and thereby limit the class of systems one can learn. However, in reality, most physical systems, such as legged robots and robotic manipulators, involve contacts and collisions, which introduce discontinuities in the states. In this paper, we introduce a differentiable contact model, which can capture contact mechanics: frictionless/frictional, as well as elastic/inelastic. This model can also accommodate inequality constraints, such as limits on the joint angles. The proposed contact model extends the scope of Lagrangian and Hamiltonian neural networks by allowing simultaneous learning of contact and system properties. We demonstrate this framework on a series of challenging 2D and 3D physical systems with different coefficients of restitution and friction. The learned dynamics can be used as a differentiable physics simulator for downstream gradient-based optimization tasks, such as planning and control.

1 Introduction

A large class of real-world physical systems evolves in a piecewise-continuous manner. For example, while playing tennis, tennis balls collide with the ground and rackets with high elasticity but follow smooth trajectories governed in between those collisions. The ability to walk/run depends heavily on the contacts between the legs and the ground. Unfavorable contact properties can significantly hinder this ability; for example, lack of friction makes it very difficult to walk on an icy road or sand. Robotic manipulators and grippers also rely on contacts and collisions to perform their assigned tasks. These examples highlight the importance of contacts and collisions, which can be found everywhere.

Encoding energy conservation into the computation graph of a neural network constitutes an effective way to improve its data-efficiency and generalization performance in inferring the dynamics of a physical system from its trajectory data [1–12]. However, most of these models assume that the trajectories of the system are smooth and governed by ordinary differential equations (ODE), in particular the Euler-Lagrange equation or Hamiltonian dynamics. Hence, they cannot deal with systems that involve contacts and collisions. On the other hand, another line of work, such as Neural Physics Engine (NPE) [13] and Interaction Networks (IN) [14], use relational priors to learn motions of physical systems. Although these models are not ODE-based and cannot learn the underlying continuous dynamics governing the smooth portion of the trajectories, they can accommodate collisions and contacts. Yang et al. [15] directly learn physical constraints and contact with an iterative neural projection (INP) algorithm. A more recent work by Hochlehnert et al. [16] learns frictionless contact while encoding a discrete form of the Euler-Lagrange equation. However,

these models have only been tested on frictionless 2D contacts. Learning of properties associated with frictional contacts and 3D contacts remains a relatively underexplored topic in the literature.

By introducing a contact model, which can handle frictional contacts both with or without elasticity, this work incorporates contacts and collisions into the energy-conserving neural networks and extends their scope. The proposed contact model solves convex optimization problems to calculate the jump in velocity during contact. To make our contact model differentiable, we build upon the recent progress on differentiating through convex optimization problems [17]. We demonstrate the performance of our contact model in learning coefficients of restitution and friction associated with a variety of 2D and 3D contacts. In addition, we also demonstrate the usefulness of this differentiable contact model along with differentiable dynamics in downstream planning tasks.

1.1 Related Work

Lagrangian/Hamiltonian-inspired Neural Networks: In the last few years, an increasing volume of work has proposed neural network models to learn the underlying dynamics from data while enforcing energy conservation. This line of works leverage Lagrangian dynamics [1, 3, 5, 7–9] or Hamiltonian dynamics [2, 4, 6, 10] to incorporate the physics prior of energy conservation into deep learning. Recently, Finzi et al. [11] show that using Cartesian coordinates and enforcing explicit constraints improve learning in both Lagrangian and Hamiltonian settings. To learn the underlying dynamics governed by an ODE, many of these prior works have used Neural ODE [18] which can learn an ODE from observed trajectories. However, real systems often exhibit non-smooth trajectories caused by sudden/abrupt changes in the velocity due to contacts and collisions. Although Neural ODE based approaches have recently been extended for learning dynamics with jump discontinuities [19–21], they cannot accommodate the physical constraints (e.g., maximum dissipation principle, non-negative normal force) associated with contacts as well as assumes knowledge about the time-instances when jump discontinuities occur. Among the energy-conserving neural networks, only [10] attempted to address contacts and collisions; to capture the elastic collision of a billiard ball, it manually reverses the momentum of the ball orthogonal to the contact surface. However, this specialized technique cannot be applied to frictional or inelastic contacts, or objects which can rotate.

Contact Model: Our contact model shares similarity with the contact model of MuJoCo [22, 23]. However, there are three differences: (1) our contact model handles elastic contacts while MuJoCo only focuses on inelastic contacts; (2) MuJoCo solves the convex optimization problem with a generalization of the Projected Gauss-Seidel method while we leverage the open-source `scs` solver [24] to solve the optimization problem; (3) the dynamics in MuJoCo is described using generalized coordinates, while we use Cartesian coordinates, since it has been shown in previous work [11], the use of Cartesian coordinates improves the learning of system properties. Another category of contact models solve contact impulses by solving a linear complementarity problem (LCP) [25]. Recently, a number of works [26–30] has proposed differentiable LCPs for downstream planning and control tasks. However, their performance on learning the contact properties has yet to be tested.

Differentiable Simulation: The recent past has also witnessed a growing interest in differentiable physics simulation that can be used in many downstream tasks (e.g., parameter estimation, planning, and control) [31–36]. Jiang et al. [37] use an LCP formulation to learn contact impulses for perfectly inelastic contacts. DiffTaichi [33] focuses on material point method and only provides intuitively simple contact mechanisms. The support for partially elastic frictional contact has yet to be provided. Geilinger et al. [34] differentiates through the dynamics solver analytically. NeuralSim [36] formulates a nonlinear complementarity problem to learn contact impulses and then solves it using Projected Gauss-Siedel. GradSim [38] uses a relaxed Coulomb model to learn contacts from video sequences. However, these prior differentiable simulation models do not focus on the energy aspect of the system and their performance on the prediction of system energy has not been evaluated.

1.2 Contribution

The main contribution of this work is three-fold. First, by introducing a differentiable contact model, we extend the usefulness of Lagrangian/Hamiltonian-inspired deep learning methods from collision-free systems to more realistic systems with contact and collisions. Second, we demonstrate the simultaneous learning of system and contact properties in a variety of physical systems by integrating the contact model with Constrained Lagrangian/Hamiltonian neural networks (CLNN/CHNN). We

show that the learned contact properties, i.e., coefficients of restitution and friction, are interpretable and match the ground truth with high accuracy. Finally, the learned contact model with CLNN/CHNN can be used to solve downstream gradient-based optimization tasks.

2 Preliminaries

2.1 Rigid body dynamics without contacts

Consider a rigid body system whose configuration at time t is described by a set of coordinates $\mathbf{x}(t) := (x_1(t), x_2(t), ..., x_D(t))$. Then the time evolution of this system can be expressed as the following second-order ODE,

$$\ddot{\mathbf{x}} = \mathbf{h}(\mathbf{x}, \dot{\mathbf{x}}; \mathbf{p}_s), \tag{1}$$

where $\mathbf{p}_s$ denote system properties, which may include inertia of objects $\mathbf{M}(\mathbf{x})$ and potential energy $V(\mathbf{x})$. The vector-valued function $\mathbf{h}$ can be derived from the laws of physics, e.g, Lagrangian/Hamiltonian dynamics. By introducing $\mathbf{v} := \dot{\mathbf{x}}$, Eqn. (1) can be written as the following first-order ODE

$$\begin{pmatrix} \dot{\mathbf{x}} \\ \dot{\mathbf{v}} \end{pmatrix} = \begin{pmatrix} \mathbf{v} \\ \mathbf{h}(\mathbf{x}, \mathbf{v}; \mathbf{p}_s) \end{pmatrix} = \mathbf{g}(\mathbf{x}, \mathbf{v}; \mathbf{p}_s). \tag{2}$$

There are two popular choices for the coordinates $\mathbf{x}$ – the generalized coordinates and the Cartesian coordinates. The generalized coordinates are usually chosen as a set of independent coordinates which implicitly enforces holonomic constraints (equality constraints, see Appendix B for details) in the system. The Cartesian coordinates are in general not independent of each other, so that the holonomic constraints in the system must be enforced explicitly in the dynamics (2). Although $\mathbf{g}$ is usually derived with generalized coordinates, this work uses Cartesian coordinates and explicit constraints [11] to demonstrate the results. The proposed contact model is independent of the set of coordinates and $\mathbf{g}$. We provide the expression of $\mathbf{g}$ used in this work in Appendix B.

2.2 Rigid body dynamics with contacts

In robotics tasks, the above assumption of no collision and contact no longer holds. For example, legged robots move around through repeated collisions/contacts between the robot legs and the ground, and robot arms grasp objects by making frictional contact with them. The difficulty of modeling these phenomena is that they essentially make the dynamics discontinuous. For example, when a ball hits the ground, its velocity changes from pointing downward to pointing upward in an infinitesimally small period of time, which can be modeled as an instantaneous change in velocity. In general, contacts, collisions, and joint limits can all be modeled in this way. Algorithm 1 shows the general procedure of modeling rigid body with contacts, where a jump in velocity is calculated by a contact model whenever there exist active contacts.

Algorithm 1: Rigid Body Dynamics with Contact

Input :

$t_0, t_1, ..., t_N$	Sequence of time points
$(\mathbf{x}_0, \mathbf{v}_0)$	Initial condition at t_0
$\mathbf{p}_s = (\mathbf{M}(\mathbf{x}), V(\mathbf{x}))$	System properties
$\mathbf{p}_c = (\boldsymbol{\mu}, \mathbf{e}_P)$	Contact properties
$\mathbf{g}(\mathbf{x}, \mathbf{v}; \mathbf{p}_s)$	First-order system dynamics

Initialize output trajectories $\mathcal{T} = \{(\mathbf{x}_0, \mathbf{v}_0)\}$.
for $i = 0 \to N - 1$ **do**
  $(\mathbf{x}_{i+1}, \mathbf{v}_{i+1}) \leftarrow$
    $\texttt{ODESolve}(\mathbf{g}, (\mathbf{x}_i, \mathbf{v}_i), t_i, t_{i+1})$
  Get active contacts $\mathbf{c}_a$ (collision detection)
  if *exists active contacts* **then**
    $\Delta\mathbf{v} \leftarrow$
      $\texttt{ContactModel}(\mathbf{x}_{i+1}, \mathbf{v}_{i+1}, \mathbf{c}_a, \mathbf{p}_s, \mathbf{p}_c)$
    $\mathbf{v}_{i+1} \leftarrow \mathbf{v}_{i+1} + \Delta\mathbf{v}$
  $\mathcal{T} \leftarrow \mathcal{T} \cup \{(\mathbf{x}_{i+1}, \mathbf{v}_{i+1})\}$

From a simulation perspective, with known system properties, contact properties (coefficients of friction and restitution), and vector field $\mathbf{g}$, the trajectory of the system can be simulated by Algorithm 1. From a learning perspective, we frame the problem as learning unknown system and contact properties from a given set of trajectories given a model prior of vector field $\mathbf{g}$. In this case, we can parametrize the unknown system and contact properties $(\mathbf{p}_s, \mathbf{p}_c)$ by neural networks and learnable parameters, predict trajectories by Algorithm 1 and minimize the difference between the predicted and actual trajectories by backpropagation. This training scheme requires all operations in the forward pass (Algorithm 1) to be differentiable. There are two key parts in the forward pass – the ODE solver module and the contact model. Operations in the ODE solver are in general differentiable, and Neural

140 ODE [18, 39] provides a framework of backpropagating through ODE solvers with constant memory
141 cost. In this work, we provide a differentiable contact model so that we can extend these previous
142 works to learn rigid body dynamics with contacts.

3 A differentiable contact model

144 In this section, we introduce a differentiable contact model for learning rigid body dynamics with
145 contacts. The proposed contact model solves post-contact velocities by solving contact impulses in
146 two phases [40] – a *compression phase*, starting from the first contact of objects till the maximum
147 compression, and a *restitution phase*, starting right after the compression phase till the separation of
148 objects. We start by presenting the constraints imposed by frictional contacts.

3.1 Contact constraints

150 This work focuses on two types of contact – frictional contact and limit constraint. For any conceptual
151 frictional contact i in a 3D contact space, the contact impulse $\mathbf{f}_i \in \mathbb{R}^3$ must lie in the friction cone,

$$\mu_i f_{i,n} \geq \sqrt{f_{i,t_1}^2 + f_{i,t_2}^2}, \quad \forall i, \tag{3}$$

152 where $\mu_i \geq 0$ is the coefficient of friction for conceptual contact i. In addition, the normal impulses
153 must be non-negative, since objects can only push but not pull others:

$$f_{i,n} \geq 0, \quad \forall i. \tag{4}$$

154 For any limit constraint such as limit in joint angle or distance, the contact space is essentially one
155 dimensional and the constraints on contact impulses are only $f_{i,n} \geq 0$. This is equivalent to setting
156 up a 3D contact space like that in frictional contact and letting $\mu_i = 0$ in Eqn. (3).

3.2 Contact model in compression phase

158 The idea behind solving contact impulses during the compression phase is the maximum dissipation
159 principle [41], which states that the compression impulses should maximize the rate of energy
160 dissipation. Equivalently, the compression impulses are those that minimizes the kinetic energy at the
161 end of the compression phase. This can be described by an optimization problem [15, 41, 22, 23].
162 We choose the following form, which is similar to the one used in Mujoco [22, 23],

$$\underset{\mathbf{f}_C^c}{\text{Minimize}} \; \frac{1}{2}(\mathbf{f}_C^c)^T \mathbf{A} \mathbf{f}_C^c + (\mathbf{f}_C^c)^T \mathbf{v}_C^{c-} \tag{5}$$

$$\text{subject to } (3), (4).$$

163 where $\mathbf{f}_C^c$ denotes the impulses in compression phase, $\mathbf{A}$ is the inverse inertia in the contact space, and
164 $\mathbf{v}_C^{c-}$ represents the velocity in the contact space before the compression phase. A concise derivation
165 of (5) from the first principle is provided in Appendix C.

3.3 Contact model in restitution phase

167 Similarly, we can set up an optimization problem to solve for the contact impulses in the restitution
168 phase $\mathbf{f}_C^r$. We assume the restitution follows Poisson's hypothesis, where the normal components
169 in $\mathbf{f}_C^r$ equals those in $\mathbf{f}_C^c$ scaled by the coefficient of restitution e_P. We adopt Poisson's hypothesis
170 instead of the popular Newton's hypothesis used in prior works [14, 26, 33, 10], because the latter
171 might result in unrealistic energy increase in certain systems [42]. Please see Appendix D for a
172 discussion. We set up the following constraint:

$$f_{i,n}^r \geq e_{P,i} \cdot f_{i,n}^c, \quad \forall i. \tag{6}$$

173 We have inequality instead of equality here since we would like to compensate for existing penetration
174 in the simulation. Since we simulate the rigid body system in discrete time steps, almost every time
175 when a collision is detected, penetration has already occurred among the objects involved in that
176 collision. Consider the case where the collision is perfectly inelastic, i.e., COR $e_P = 0$, then the
177 true normal impulse during restitution phase would be zero, which fails to fix existing unrealistic

penetration. By setting up the constraint as in Eqn. (6), a larger normal impulse is allowed to fix existing penetration. The optimization problem in the restitution phase is

$$\underset{\mathbf{f}_C^r}{\text{Minimize}} \; \frac{1}{2}(\mathbf{f}_C^r)^T \mathbf{A}\mathbf{f}_C^r + (\mathbf{f}_C^r)^T(\mathbf{v}_C^{c+} - \mathbf{v}_C^*) \tag{7}$$

subject to $(3), (4), (6)$.

with $\mathbf{v}_C^{c+}$ as the velocity in contact space after the compression phase and $\mathbf{v}_C^*$ as the target velocity, used for fixing penetration. If there's no penetration, $\mathbf{v}_C^* = \mathbf{0}$. A detailed discussion of penetration compensation can be found in Appendix G. From the principle of maximum dissipation, the equality in Eqn. (6) would hold for the solution when COR is large and penetration is small, thus respecting Poisson's hypothesis when penetration can be fixed naturally.

3.4 Differentiability

Solving optimization problems (5) and (7) for contact impulses allow us to calculate instantaneous velocity change to perform simulation (Algorithm 1). Moreover, we would like to back-propagate through the contact model to learn unknown properties. In fact, our proposed contact model is differentiable, thanks to recent progress on differentiable convex optimization layers. Both problems (5) and (7) are convex optimization problems with convex quadratic objectives (we show that $\mathbf{A}$ is positive semi-definite in Appendix E) as well as linear constraints and second-order cone constraints. We can then express our problems using disciplined parametrized programming (DPP) and set up these two problems as differentiable layers using CvxpyLayers [17]. Thus, our model is differentiable and can be used in dynamics and parameter learning as well as downstream tasks.

4 Dynamics and parameter learning experiments

4.1 Simulated systems

To evaluate the proposed contact model, we simulate five different systems with contacts (Fig. 1) and propose eight dynamics and parameter learning tasks based on these systems with different contact properties (Table 1). Previous work has studied the *bouncing point masses* (Fig. 1(a)) which is often referred to as billiards or bouncing balls [10, 26, 33]. To make this task more challenging, we let the size of each object be different. We also propose the *bouncing disks* (Fig. 1(b)) where each disk can rotate. The 2-pendulum colliding with the ground has been used to study and analyze contact models for more than three decades [43]. We make this task more challenging by studying a *3-Pendulum colliding with the ground* (Fig. 1(c)). The gyroscope is a 3D system that exhibits complex dynamics. A *gyroscope colliding with a wall* (Fig. 1(d)) is a system where the normal contact impulse does not point towards the center of mass (c.o.m), and Newton's hypothesis might give an unrealistic result with increased energy after collisions [42]. The *rope* (Fig. 1(e)) has also been studied in previous works [14, 15]. Please refer to Appendix H for further details about these systems and the tasks.

4.2 Dynamics and parameter learning experimental setup

In each task, we jointly learn system and contact properties from trajectory data by extending CLNN/CHNN with the proposed contact model.

Data: For each task, the training set is generated by randomly sampling 800 collision-free initial conditions and then simulating the dynamics for 100 time steps. Since for some systems, there are

(a)   (b)   (c)   (d)   (e)

Figure 1: Simulated systems with contact. From left to right: (a) bouncing point masses, (b) bouncing disks, (c) chained pendulums with ground, (d) gyroscope with a wall, and (e) rope. The black lines in bouncing disks show the orientations of disks.

Table 1: Benchmark tasks. The columns D, E, max(C) denote dimension of the dynamics, number of equality constraints, and the maximum number of contacts that could be simultaneously active , respectively.

Name	System	D	E	max(C)	Space	Same e_P, μ for all contacts	Conserve energy
BP5-e	Bouncing point masses	10	0	8	2D	Y	Y
BP5	Bouncing point masses	10	0	8	2D	N	N
CP3-e	Chained pendulums w/ ground	6	3	1	2D	Y	Y
CP3	Chained pendulums w/ ground	6	3	1	2D	Y	N
BD5	Bouncing disks	30	15	8	2D	N	N
Rope	Rope	400	0	~ 399	2D	Y	N
Gyro-e	Gyroscope w/ a wall	12	7	1	3D	Y	Y
Gyro	Gyroscope w/ a wall	12	7	1	3D	Y	N

very few data points in a trajectory that involves collision, we select a small chunk containing 5 consecutive time steps from each simulated trajectory such that the final training set contains 800 trajectories of length 5, where around half of the trajectories contain collisions and the other half are collision-free. We also make sure that the initial state of these selected chunks is collision-free. The evaluation and test set are generated in a similar way with 100 trajectories, respectively.

Architecture and training details: In the experiments, we assume the system properties, i.e., object inertia and potential energy, as well as contact properties, i.e., coefficients of friction and restitution, are unknown and need to be learned from trajectory data. The system properties are parametrized as in CLNN and CHNN [11]. As for contact properties, all coefficient of friction are non-negative, so they are parametrized by scalar learnable parameters passed through *ReLu* function. As each coefficient of restitution lies in the interval of $[0, 1]$, it is parametrized by a learnable parameter passed through a *hard sigmoid* function. The predicted trajectories are generated by running Algorithm 1 with parametrized system and contact properties. We use RK4 as the ODE solver in Neural ODE. We compute the L_1-norm of the difference between predicted and true trajectories, and use it as the loss function for training. The gradients are computed by differentiating through Algorithm 1, and learnable parameters are updated using the AdamW optimizer [44, 45] with a learning rate of 0.001.

Models: We implement two slightly different versions of the contact model. The first version, referred to as CM, set up optimization problems exactly as stated in (5) and (7). The second version, referred to as CMr, adds a diagonal positive regularization matrix $\mathbf{R} = \epsilon\mathbf{I}$ to $\mathbf{A}$ in (5) and (7), such that $(\mathbf{A} + \mathbf{R})$ is always positive definite, which ensures a unique global minimum in each problem. These two versions are combined with CLNN and CHNN to set up the following four neural network models: (*i*) CM-CD-CLNN, (*ii*) CM-CD-CHNN, (*iii*) CMr-CD-CLNN, and (*iv*) CMr-CD-CHNN. The "CD" in model names emphasizes that we assume that a collision detection module is given.

Baselines: We also set up three baselines. In the first baseline, MLP-CD-CLNN, we calculate the instantaneous velocity change from a multi-layer perceptron (MLP) instead of the proposed contact model. Our second baseline, IN-CP-CLNN, calculates velocity change from an interaction network (IN) [14] without requiring a collision detection module, since IN has the ability to learn collisions and contact. IN requires true system and contact properties as input. Here the "CP" in the model name emphasizes true contact properties are given and the system properties learned by CLNN are fed into IN. Our last baseline, IN-CP-SP, is the original interaction network which has shown strong ability in predicting 2D rigid body trajectories without equality constraints, but haven't been tested on systems with equality constraints or 3D systems. The name emphasizes that true system and contact properties are known and no collision detection module is needed. Similar to IN [14], we transform each trajectory into multiple one-step pairs to train these baseline models. We also attempted to use the LCP formulation of contact model [26] as a baseline. However, the implementation of gradient computation of the LCP function in [26] results in NaN in our examples. Please refer to Appendix J for additional details. As the forward computation of LCP works as expected, we use LCP-generated training data to test the robustness of our model.

4.3 Dynamics and parameter learning results

Our implementation relies on publically available codebases including Pytorch [46], CHNN [11], Symplectic ODE-Net [4] and Neural ODE [18]. We handle training using Pytorch Lightning [47] for the purpose of reproducibility. All data generation and training are done using a Intel Xeon CPU.

Figure 2: Trajectory relative error (log scale) with 95% confidence interval error bars. Each error is averaged over 100 test trajectories of length 5.

Figure 3: Relative error (log scale) along long test trajectories (50 times steps). Each curve is averaged over 100 test trajectories. Vertical dashed lines show the trajectory length during training.

Prediction: We report the average relative L_1 error over the test trajectories of 7 models on 8 tasks in Fig. 2. In all tasks, our models beat baseline models. The performance difference between CLNN and CHNN is minor since their architectures are similar. In most tasks, CM outperforms CMr. In the BP5-E task, CM beats CMr by 2 orders of magnitude. IN does not perform well even in BP5 tasks since our training set (3.2k one-step pairs) is much smaller than the dataset (1M one-step pairs) used in the IN paper. We also report average relative L_1 errors along test trajectories of 50 time steps in Fig. 3 in order to show each model's ability in long term prediction. We observe that our contact model CM and CMr outperform baselines in all tasks.

Interpretable mass ratio: Without direct supervision on mass, deep learning algorithms are unlikely to recover the true mass, as pointed out in [4]. However, one can still inspect the ratio of learned mass values to see how well this physical property is learned. The mass ratio plays an important role in determining the motion of objects when they interact with each other, e.g., during collisions. In our BP5 task, CM-CD-CLNN learns the mass ratio $[m_2/m_1, m_3/m_1, m_4/m_1, m_5/m_1] = [2.0001, 6.0036, 8.0014, 10.0024]$, which is very close to the true ratio $[2, 6, 8, 10]$. In fact, our framework is able to accurately learn mass ratios across tasks (please see Appendix I for details).

Interpretable contact properties: Table 2 shows the learned contact properties by our 4 models in 6 tasks where the contact properties are the same for all contacts. For all tasks, CM can learn contact properties accurately which explains its good performance in prediction. CMr is an approximate model and does not infer contact properties as accurately as CM. The interpretability of the learned

Table 2: Learned contact properties from our 4 models on 6 tasks that has unique contact properties for all contacts. Bold numbers are the best learned contact properties in each task across 4 models.

	BP5-e μ	BP5-e e_P	CP3-e μ	CP3-e e_P	CP3 μ	CP3 e_P	Rope μ	Rope e_P	Gyro-e μ	Gyro-e e_P	Gyro μ	Gyro e_P
Ground Truth	0.000	1.000	0.000	1.000	0.500	0.000	0.000	0.000	0.000	1.000	0.100	0.800
CM-CD-CLNN	**0.000**	**1.000**	**0.000**	**1.000**	0.500	0.005	0.026	0.000	**0.000**	**1.000**	**0.100**	**0.800**
CM-CD-CHNN	**0.000**	**1.000**	**0.000**	**1.000**	**0.500**	**0.004**	**0.017**	**0.000**	**0.000**	**1.000**	**0.100**	**0.800**
CMr-CD-CLNN	**0.000**	**1.000**	0.002	1.000	0.500	0.023	0.037	0.011	0.002	1.000	0.099	0.892
CMr-CD-CHNN	**0.000**	**1.000**	0.002	1.000	0.500	0.023	0.046	0.019	0.002	1.000	0.099	0.893

Figure 4: Energy of the predicted trajectories of all 7 models on a sampled test initial condition from BP5-e, CP3-e and Gyro-e tasks. The true energy in each task is represented by the horizontal black line in the middle, which is constant along the trajecotry.

contact properties along with mass ratio explains the performance of our framework and shows that the proposed contact model indeed extends Lagrangian and Hamiltonian neural networks.

Energy: The prior of Lagrangian/Hamiltonian dynamics conserve energy along each collision-free trajectory, which is one of the reason that Lagrangian/Hamiltonian-based neural network models perform better in prediction and generalization [1, 2, 4, 7]. Fig. 4 illustrates how the total energy changes over time for the predicted trajectories of 7 models on 3 tasks that conserve energy since the contacts are elastic and frictionless (i.e., $e_P = 1, \mu = 0$). Models using CM perform the best in conserving energy in all three tasks since CM learns contact properties perfectly (Table 2). This demonstrates that the proposed contact model can uncover the energy conserving aspect even though energy conservation has not been not enforced explicitly. Models using CMr lose energy each time collision happens in CP3-e and Gyro-e, since they learn positive coefficients of friction in these tasks (Table 2). The baseline models perform the worst in terms of energy conservation.

Sample efficiency: We use the BP5 task to demonstrate the sample efficiency of this proposed framework. We vary the training sample size from 25 to 800 trajectories and report the validation loss (L_1-norm) of CM-CD-CLNN, MLP-CD-CLNN, and Interaction Network. Figure 5 shows that our framework works well even with limited training data.

Figure 5: Trajectory absolute errors of BP5 task.

Scalability: In Table 3, we report the average wall clock time in each iteration (forward pass and backward pass) during training of three sizes of ropes. The time scales approximately linearly with the numbers of coordinates (D) and contacts (C).

Table 3: Average wall clock time in each iteration of ropes. The last column shows increases in time

D	max(C)	time (s)	
100	~ 99	0.869	1x
200	~ 199	1.563	1.7x
400	~ 399	3.225	3.7x

Robustness: We evaluate the robustness of our framework by training our model using data generated by LCP formulation and noisy data. (See Appendix K for details.) When trained on LCP data, our framework can learn accurate contact properties in 2D tasks. For the 3D Gyro tasks, the learned contact properties are not as accurate (e.g., learned COR of 0.822 instead of 0.800). This is expected since the LCP formulation relaxes the 3D friction cone into a (linear) polyhedral cone and the direction of friction impulses would deviate from those given by our contact model, which is based on the second-order friction cone. In addition, we observe that the performance of our model does not suffer from noisy data since we incorporate strong physics prior into deep learning. For CMr, we also perform an ablation test to investigate the influence of the amount of regularization. By setting the regularizer as $\mathbf{R} = \epsilon\mathbf{I}$, we observe that a smaller ϵ (e.g. 0.001) result in more accurate learned contact properties, since a smaller ϵ approximate (5) and (7) better. However, making ϵ a learnable parameter does not improve accuracy. Please see Appendix K for more details.

5 Downstream tasks

Since our framework is differentiable, we can use it as a differentiable physics simulator to solve downstream tasks after we have learned the system and contact properties. Here we demonstrate this capability by considering three gradient-based trajectory planning tasks and using CM-CD-CLNN.

Billiards: We study the same billiard task as in DiffTaichi [33]. The goal is to find the initial position and velocity of the white ball such that blue ball hit the black target in 1024 time steps. In order to test

| (a) | (b) | (c) |

Figure 6: Billiards. (a), (b): The solid white ball and arrow shows the initial position and velocity optimized by CM-CD-CLNN and DiffTaichi, respectively, while dashed white ball and arrow shows those of the white ball before optimization. (c): Loss as a function of training iterations.

our framework's ability to solve downstream task and make comparison with DiffTaichi, we assume the parameters such as mass and contact properties are known, the same assumption in DiffTaichi. Fig 6(a) and 6(b) shows the solution found by our proposed model and DiffTaichi, respectively. This task does not have a unique solution since one can place the white ball closer to the billiards with a relatively small initial velocity (e.g. DiffTaichi solution) or place the ball farther away from the billiards with a relatively large initial velocity (e.g. our solution). Since in this example DiffTaichi implements a simpler contact and dynamics model, it has better convergence in learning, as shown in Fig 6(c), where loss is the distance between the black target and the blue ball after 1024 time steps.

Throwing: We present two throwing tasks as shown in Fig. 7 and similar tasks has also been studied in [34, 35]. In the "hit" task (Fig. 7(a)), the initial position of the disk is fixed, the goal is to find a desired initial velocity so that the disk reach the target (black circle) after exactly one bounce off the ground. In the "vertical" task (Fig. 7(b)), the initial position and the translational velocity are fixed, so that first half of the center of mass (c.o.m) trajectory (dashed blue) is fixed. The goal is to find a desired initial angular velocity such that the second half of the c.o.m trajectory is as close to a vertical line (dashed black) as possible. In this task we need to learn a counterclockwise spin such that when the disk bounces off the ground, there are enough friction to stop the horizontal motion.

| (a) | (b) |

Figure 7: Blue hollow circles indicate the initial position of the disk. (a) the "hit" task. The black hollow circle indicate the target position. (b) the "vertical" task.

For these two tasks, we parametrize the initial condition to be learned, simulate the trajectory based on the learned system and contact properties, and minimize the difference between simulated outcome and the goal by gradient descent. We can successfully find the initial conditions to achieve the tasks, evaluated using the true system and contact properties. (Please see the video for additional details).

6 Conclusion

In this work, we have introduced a differentiable contact model, which can capture contact dynamics with different properties. Our contact model extends the applicability of Lagrangian/Hamiltonian-inspired neural networks to enable learning of hybrid dynamics in rigid body systems and offer interpretability about system and contact properties. We show that the proposed framework achieves better prediction with fewer samples and is robust against noisy data or LCP-generated data. Future works will incorporate model-based control and explore interpretable safe control policies for robotics applications. A particular direction could be to develop appropriate energy shaping control policies [48] and integrate them with this proposed learning framework.

Limitations: Our framework assumes a known collision detection module. Although it can be obtained from an idealized touch feedback sensor [16], this information might be unavailable in other scenarios. Future work would explore how to relax this assumption.

Societal impact: We introduce a framework for data-driven dynamics modelling which uses physics-based priors to improve generalization, sample efficiency, and interpretability. Data-driven dynamics modelling, in general, can have a profound effect in learning-based control synthesis, especially in robotics and automation. However, our proposed framework is still a conceptual proposal and has a very low (around 2) Technology Readiness Level (TRL) [49]. We are yet to fully understand its limitations and failure scenarios that can significantly influence its real-world adoption.

References

[1] Michael Lutter, Christian Ritter, and Jan Peters. Deep Lagrangian Networks: Using Physics as Model Prior for Deep Learning. In *International Conference on Learning Representations*, 2019.

[2] Samuel Greydanus, Misko Dzamba, and Jason Yosinski. Hamiltonian neural networks. In *Advances in Neural Information Processing Systems*, volume 32, pages 15379–15389, 2019.

[3] Michael Lutter, Kim Listmann, and Jan Peters. Deep Lagrangian Networks for end-to-end learning of energy-based control for under-actuated systems. In *2019 IEEE/RSJ International Conference on Intelligent Robots and Systems (IROS)*, pages 7718–7725, 2019.

[4] Yaofeng Desmond Zhong, Biswadip Dey, and Amit Chakraborty. Symplectic ODE-Net: Learning Hamiltonian Dynamics with Control. In *International Conference on Learning Representations*, 2020.

[5] Manuel A. Roehrl, Thomas A. Runkler, Veronika Brandtstetter, Michel Tokic, and Stefan Obermayer. Modeling System Dynamics with Physics-Informed Neural Networks Based on Lagrangian Mechanics. In *21st IFAC World Congress*, 2020.

[6] Yaofeng Desmond Zhong, Biswadip Dey, and Amit Chakraborty. Dissipative SymODEN: Encoding Hamiltonian Dynamics with Dissipation and Control into Deep Learning. In *ICLR 2020 Workshop on Integration of Deep Neural Models and Differential Equations*, 2020.

[7] Miles Cranmer, Sam Greydanus, Stephan Hoyer, Peter Battaglia, David Spergel, and Shirley Ho. Lagrangian neural networks. In *ICLR 2020 Workshop on Integration of Deep Neural Models and Differential Equations*, 2020.

[8] Yaofeng Desmond Zhong and Naomi Leonard. Unsupervised Learning of Lagrangian Dynamics from Images for Prediction and Control. volume 33, 2020.

[9] Christine Allen-Blanchette, Sushant Veer, Anirudha Majumdar, and Naomi Ehrich Leonard. LagNetViP: A Lagrangian Neural Network for Video Prediction. *arXiv preprint arXiv:2010.12932*, 2020.

[10] Zhengdao Chen, Jianyu Zhang, Martin Arjovsky, and Léon Bottou. Symplectic Recurrent Neural Networks. In *International Conference on Learning Representations*, 2020.

[11] Marc Finzi, Ke Alexander Wang, and Andrew Gordon Wilson. Simplifying Hamiltonian and Lagrangian Neural Networks via Explicit Constraints. volume 33, 2020.

[12] Yaofeng Desmond Zhong, Biswadip Dey, and Amit Chakraborty. Benchmarking Energy-Conserving Neural Networks for Learning Dynamics from Data. *arXiv preprint arXiv:2012.02334*, 2020.

[13] Michael B Chang, Tomer Ullman, Antonio Torralba, and Joshua B Tenenbaum. A compositional object-based approach to learning physical dynamics. In *International Conference on Learning Representations*, 2017.

[14] Peter Battaglia, Razvan Pascanu, Matthew Lai, Danilo Jimenez Rezende, and Koray Kavukcuoglu. Interaction Networks for Learning about Objects, Relations and Physics. In *Advances in Neural Information Processing Systems*, volume 29, pages 4502–4510, 2016.

[15] Shuqi Yang, Xingzhe He, and Bo Zhu. Learning physical constraints with neural projections. In *Advances in Neural Information Processing Systems*, volume 33, 2020.

[16] Andreas Hochlehnert, Alexander Terenin, Steindor Saemundsson, and Marc Deisenroth. Learning Contact Dynamics using Physically Structured Neural Networks. In *Proceedings of The 24th International Conference on Artificial Intelligence and Statistics*, volume 130, pages 2152–2160, 2021.

[17] Akshay Agrawal, Brandon Amos, Shane Barratt, Stephen Boyd, Steven Diamond, and J. Zico Kolter. Differentiable Convex Optimization Layers. In *Advances in Neural Information Processing Systems*, volume 32, pages 9562–9574, 2019.

[18] Ricky TQ Chen, Yulia Rubanova, Jesse Bettencourt, and David K Duvenaud. Neural ordinary differential equations. In *Advances in Neural Information Processing Systems*, volume 31, pages 6571–6583, 2018.

[19] Junteng Jia and Austin R Benson. Neural Jump Stochastic Differential Equations. In *Advances in Neural Information Processing Systems*, volume 32, pages 9847–9858, 2019.

[20] Daehoon Gwak, Gyuhyeon Sim, Michael Poli, Stefano Massaroli, Jaegul Choo, and Edward Choi. Neural ordinary differential equations for intervention modeling. *arXiv preprint arXiv:2010.08304*, 2020.

[21] Calypso Herrera, Florian Krach, and Josef Teichmann. Neural Jump Ordinary Differential Equations. In *International Conference on Learning Representations*, 2021.

[22] Emanuel Todorov. A convex, smooth and invertible contact model for trajectory optimization. In *2011 IEEE International Conference on Robotics and Automation*, pages 1071–1076, 2011.

[23] Emanuel Todorov. Convex and analytically-invertible dynamics with contacts and constraints: Theory and implementation in MuJoCo. In *2014 IEEE International Conference on Robotics and Automation (ICRA)*, pages 6054–6061, 2014.

[24] Brendan O'donoghue, Eric Chu, Neal Parikh, and Stephen Boyd. Conic optimization via operator splitting and homogeneous self-dual embedding. *Journal of Optimization Theory and Applications*, 169(3):1042–1068, 2016.

[25] Mihai Anitescu and Florian A Potra. Formulating dynamic multi-rigid-body contact problems with friction as solvable linear complementarity problems. *Nonlinear Dynamics*, 14(3):231–247, 1997.

[26] Filipe de Avila Belbute-Peres, Kevin Smith, Kelsey Allen, Joshua B Tenenbaum, and J Zico Kolter. End-to-end differentiable physics for learning and control. In *Advances in neural information processing systems*, volume 31, pages 7178–7189, 2018.

[27] Jonas Degrave, Michiel Hermans, Joni Dambre, and Francis Wyffels. A differentiable physics engine for deep learning in robotics. *Frontiers in Neurorobotics*, 13:6, 2019.

[28] Junbang Liang and Ming C. Lin. Differentiable Physics Simulation. In *ICLR 2020 Workshop on Integration of Deep Neural Models and Differential Equations*, 2020.

[29] Changkyu Song and Abdeslam Boularias. Identifying mechanical models of unknown objects with differentiable physics simulations. In *Proceedings of the 2nd Conference on Learning for Dynamics and Control*, volume 120 of *Proceedings of Machine Learning Research*, pages 749–760. PMLR, 2020.

[30] Changkyu Song and Abdeslam Boularias. Learning to slide unknown objects with differentiable physics simulations. In *Robotics science and systems*, 2020.

[31] Zhenjia Xu, Jiajun Wu, Andy Zeng, Joshua B Tenenbaum, and Shuran Song. DensePhysNet: Learning Dense Physical Object Representations via Multi-step Dynamic Interactions. *arXiv preprint arXiv:1906.03853*, 2019.

[32] Alvaro Sanchez-Gonzalez, Jonathan Godwin, Tobias Pfaff, Rex Ying, Jure Leskovec, and Peter Battaglia. Learning to simulate complex physics with graph networks. In *International Conference on Machine Learning*, pages 8459–8468, 2020.

[33] Yuanming Hu, Luke Anderson, Tzu-Mao Li, Qi Sun, Nathan Carr, Jonathan Ragan-Kelley, and Fredo Durand. DiffTaichi: Differentiable Programming for Physical Simulation. In *International Conference on Learning Representations (ICLR)*, 2020.

[34] Moritz Geilinger, David Hahn, Jonas Zehnder, Moritz Bächer, Bernhard Thomaszewski, and Stelian Coros. ADD: Analytically Differentiable Dynamics for Multi-Body Systems with Frictional Contact. *ACM Transactions on Graphics*, 39(6), November 2020. ISSN 0730-0301.

[35] M. Macklin, K. Erleben, M. Müller, N. Chentanez, S. Jeschke, and T. Y. Kim. Primal/dual descent methods for dynamics. In *Proceedings of the ACM SIGGRAPH/Eurographics Symposium on Computer Animation*, 2020.

[36] Eric Heiden, David Millard, Erwin Coumans, Yizhou Sheng, and Gaurav S. Sukhatme. NeuralSim: Augmenting Differentiable Simulators with Neural Networks. In *IEEE International Conference on Robotics and Automation (ICRA)*, 2021.

[37] Yifeng Jiang, Jiazheng Sun, and C. Karen Liu. Data-Augmented Contact Model for Rigid Body Simulation. *arXiv Preprint*, 1803.04019, 2018.

[38] J. Krishna Murthy, Miles Macklin, Florian Golemo, Vikram Voleti, Linda Petrini, Martin Weiss, Breandan Considine, Jérôme Parent-Lévesque, Kevin Xie, Kenny Erleben, Liam Paull, Florian Shkurti, Derek Nowrouzezahrai, and Sanja Fidler. GradSim: Differentiable simulation for system identification and visuomotor control. In *International Conference on Learning Representations (ICLR)*, 2021.

[39] Stefano Massaroli, Michael Poli, Jinkyoo Park, Atsushi Yamashita, and Hajime Asama. Dissecting Neural ODEs. In *Advances in Neural Information Processing Systems*, volume 33, 2020.

[40] Siméon D Poisson. Mechanics, vol. ii. *Trans. HH Harte, Longman, London*, 1817.

[41] David E Stewart. Rigid-body dynamics with friction and impact. *SIAM Review*, 42(1):3–39, 2000.

[42] Shlomo Djerassi. Collision with friction; Part A: Newton's hypothesis. *Multibody System Dynamics*, 21(1):37, 2009.

[43] Thomas R Kane and David A Levinson. *Dynamics, Theory and Applications*. McGraw Hill, 1985.

[44] Diederik P Kingma and Jimmy Ba. Adam: A method for stochastic optimization. *arXiv preprint arXiv:1412.6980*, 2014.

[45] Ilya Loshchilov and Frank Hutter. Decoupled weight decay regularization. In *International Conference on Learning Representations*, 2019.

[46] Adam Paszke, Sam Gross, Francisco Massa, Adam Lerer, James Bradbury, Gregory Chanan, Trevor Killeen, Zeming Lin, Natalia Gimelshein, Luca Antiga, Alban Desmaison, Andreas Kopf, Edward Yang, Zachary DeVito, Martin Raison, Alykhan Tejani, Sasank Chilamkurthy, Benoit Steiner, Lu Fang, Junjie Bai, and Soumith Chintala. PyTorch: An Imperative Style, High-Performance Deep Learning Library. In *Advances in Neural Information Processing Systems*, volume 32, pages 8024–8035. 2019.

[47] et al. Falcon, W. A. PyTorch Lightning. *GitHub. Note: https://github.com/PyTorchLightning/pytorch-lightning*, 3, 2019.

[48] Jianping Lin, Nikhil Divekar, Ge Lv, and Robert D. Gregg. Energy shaping control with virtual spring and damper for powered exoskeletons. In *2019 IEEE 58th Conference on Decision and Control (CDC)*, pages 3039–3045, 2019. doi: 10.1109/CDC40024.2019.9029624.

[49] Steven Hirshorn and Sharon Jefferies. Final Report of the NASA Technology Readiness Assessment (TRA) Study Team. 2016.